

中性子過剰原子核の存在限界とその新しい原理 ~核力に基づく大規模計算による解析

Naofumi Tsunoda

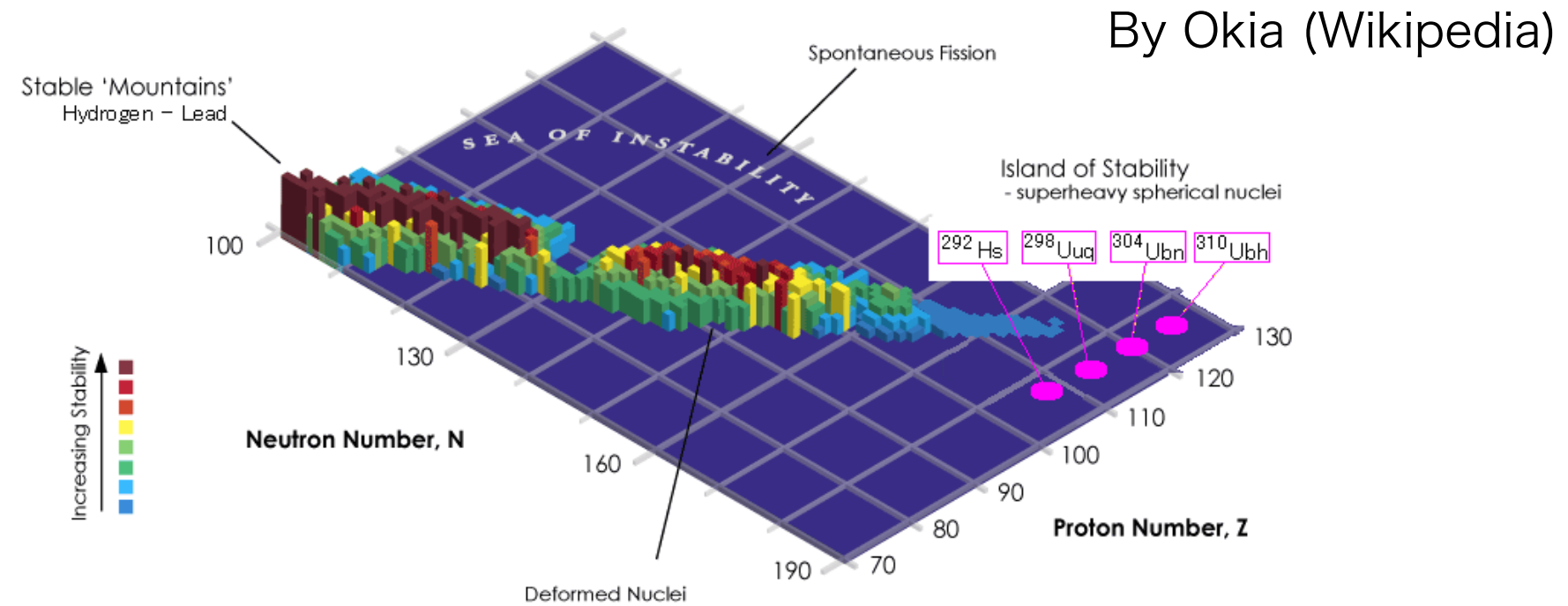
Center for Nuclear Study, the University of Tokyo

素粒子・原子核・宇宙「京からポスト京に向けて」
シンポジウム

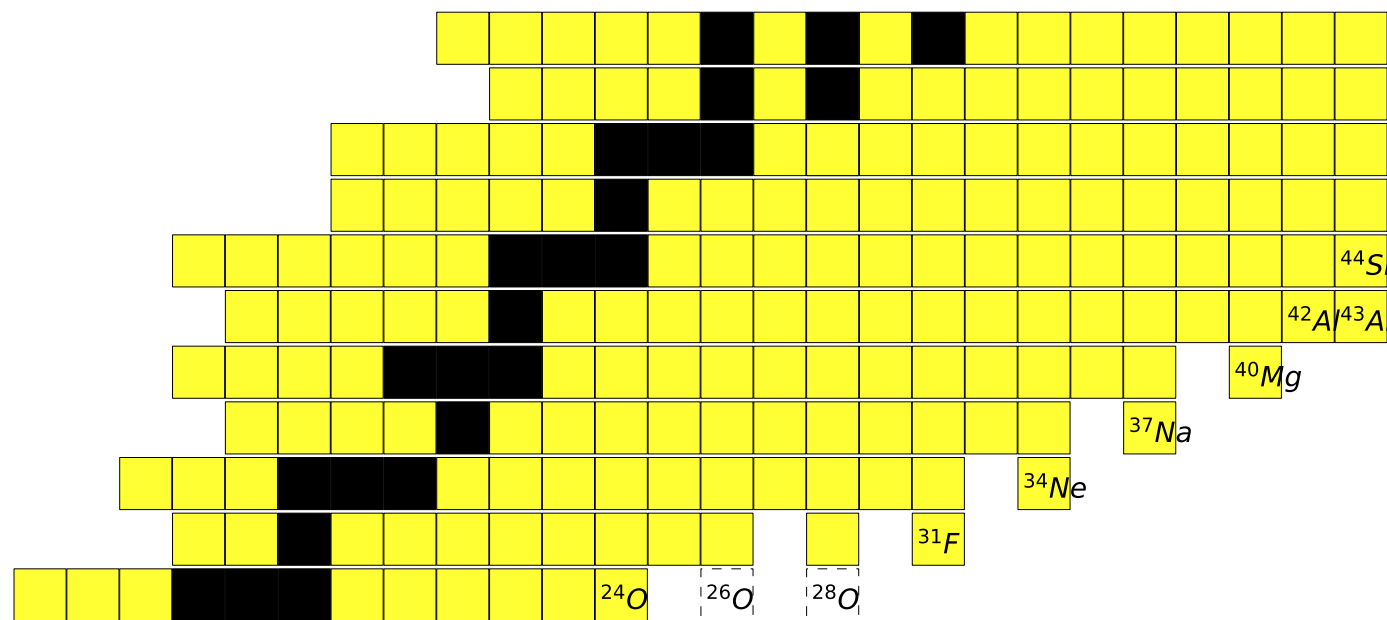
2019/01/09-01/10

This work has been supported by MEXT and JICFuS as a priority issue (Elucidation of the fundamental laws and evolution of the universe) to be tackled by using Post “K” Computer.

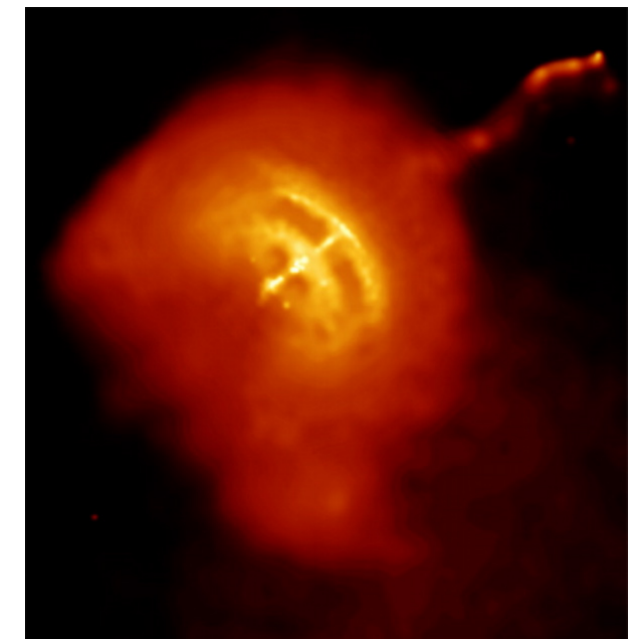
安定の島



原子核のドリップライン



中性子星



By NASA/CXC/PSU/G.Pavlov et al.

BROOKHAVEN
NATIONAL LABORATORY

National Nuclear Data Center

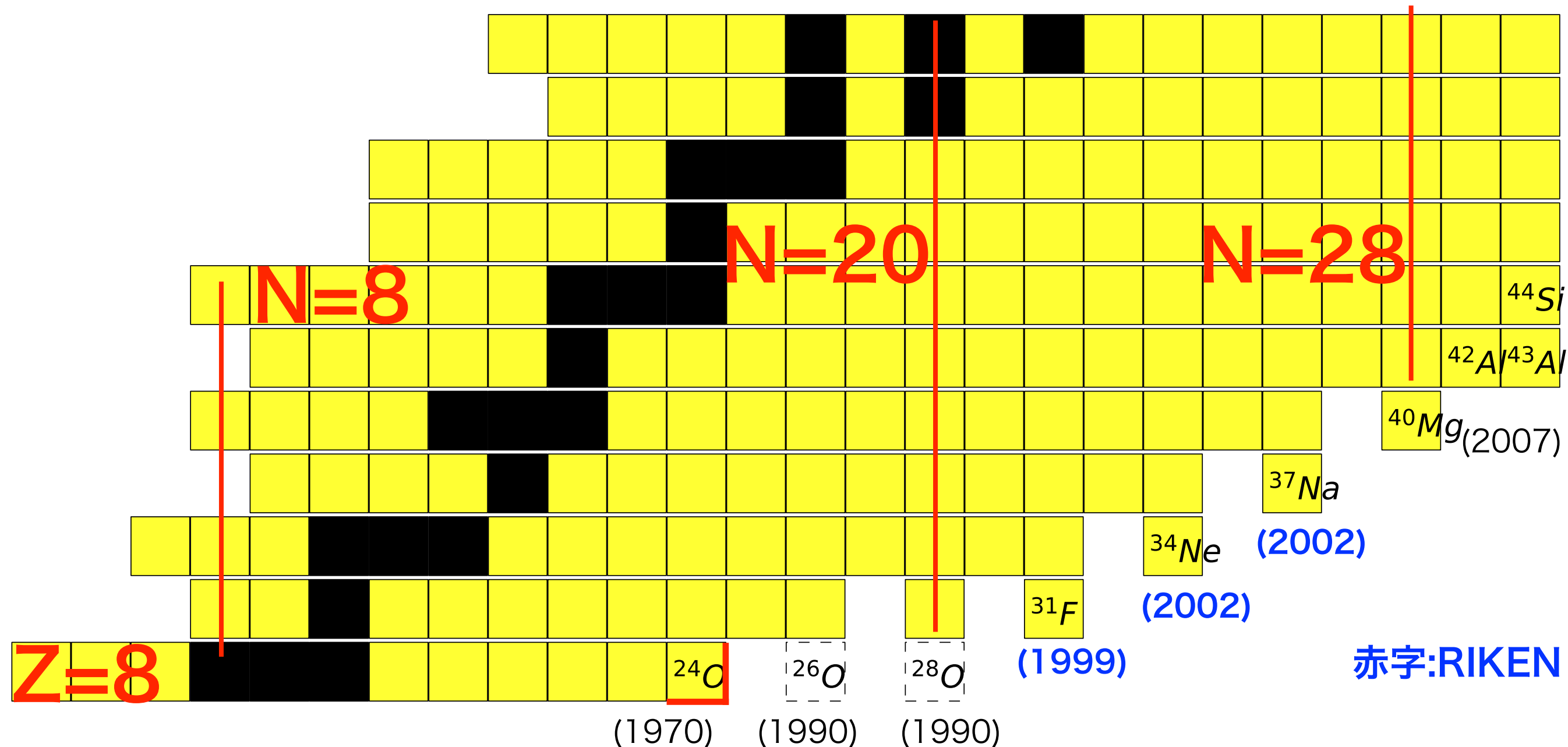
ドリップラインの場所？
原理、メカニズム？

ドリップラインを確定する
ことは、原子核物理学
の重要な使命の一つ!!

National Nuclear Data Center, information extracted from the NuDat 2 database,
<http://www.nndc.bnl.gov/nudat2/>
"Users should feel free to use the information from NuDat 2 (tables
and plots) in their work, reports, presentations, articles and books."

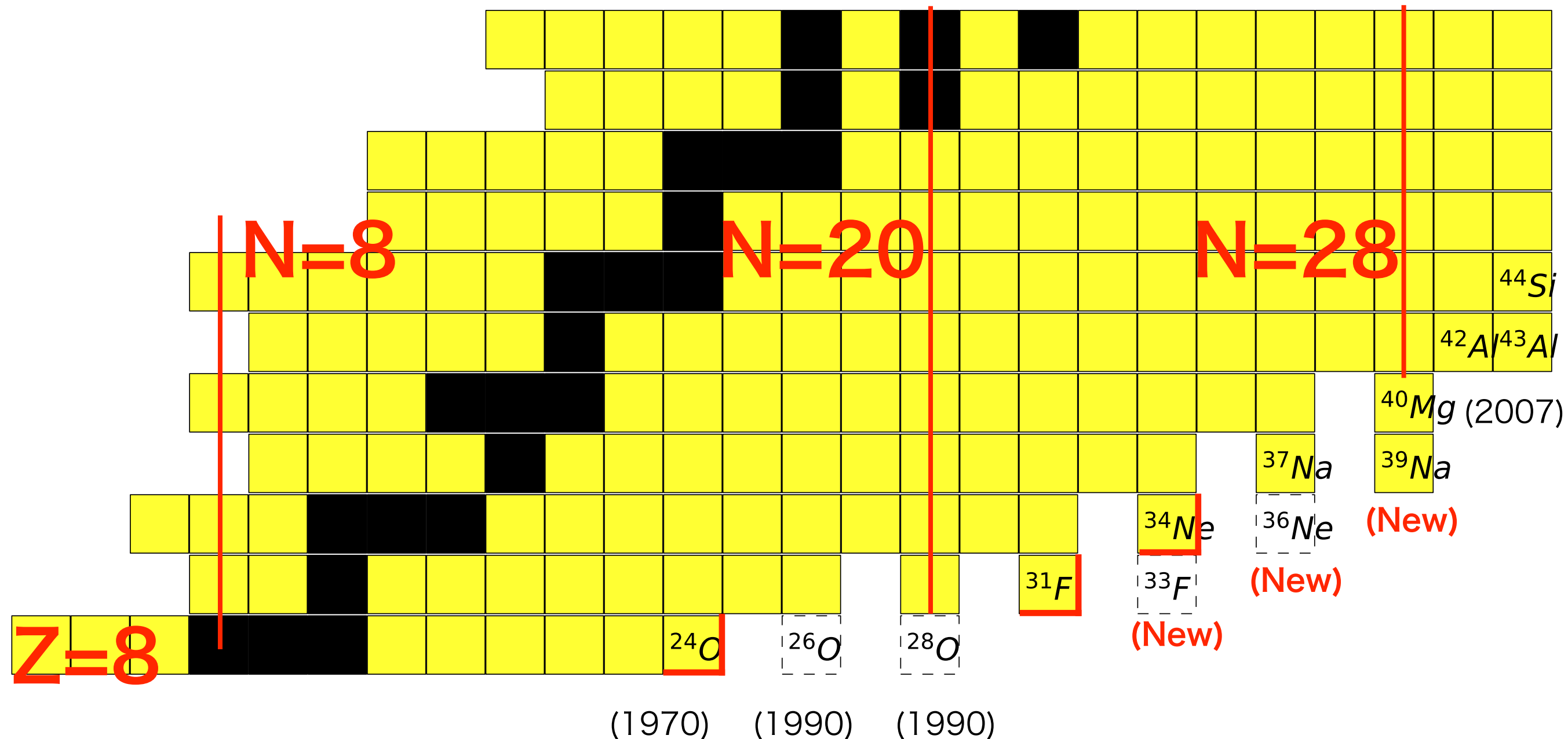
By The viewer - National Nuclear Data Center

加速器実験で測れるドリップライン (少し前)



- ドリップラインが確定しているのはOまで(軽い核のみ)
- F 以降はこの時点でまだ不明

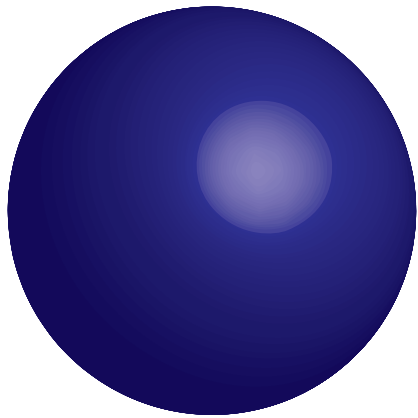
加速器実験で測れるドリップライン (現在)



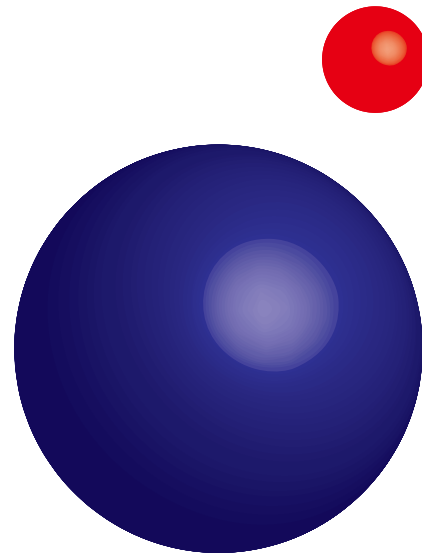
- F/Ne のドリップラインが確定 (^{31}F , ^{34}Ne)
- ^{39}Na の存在を確認 (ドリップラインはまだ不明)
- Mg 以降は次世代の実験
- Mechanism?

中性子ドリップライン

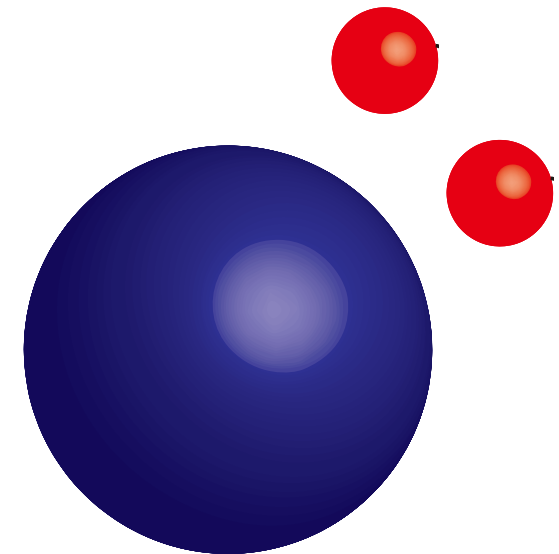
$E(Z, N)$: 結合エネルギー



$E(Z, N)$



$E(Z, N+1)$



$E(Z, N+2)$

$E(Z, N) < E(Z, N+1, 2) \Rightarrow (Z, N+1, 2)$ は **bound**

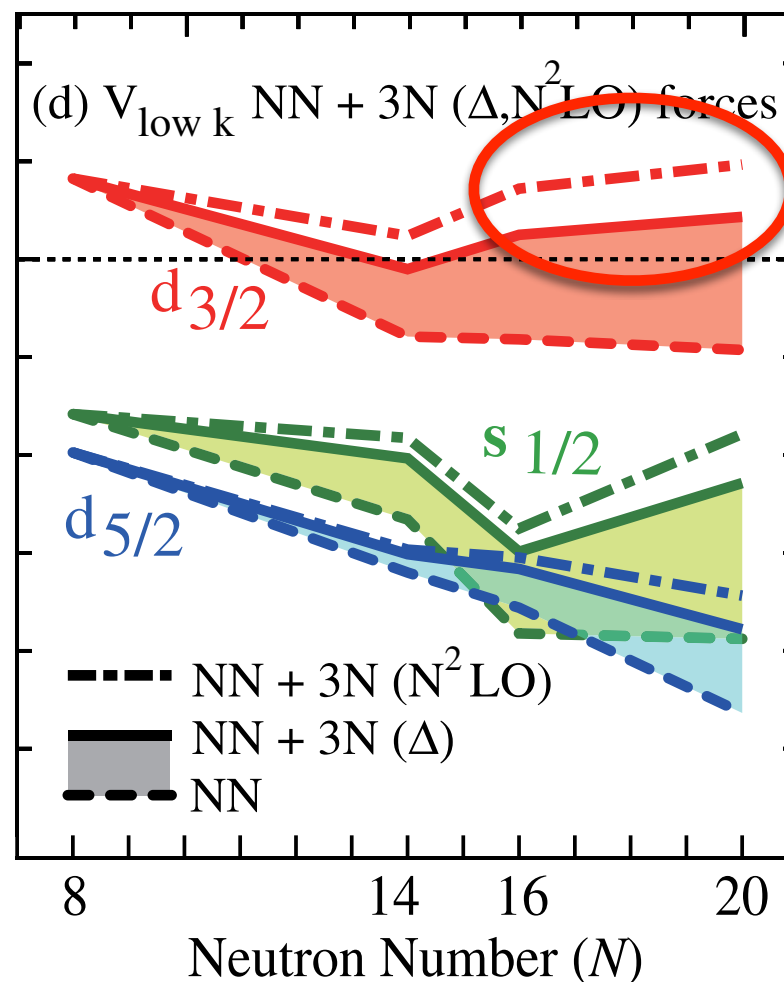
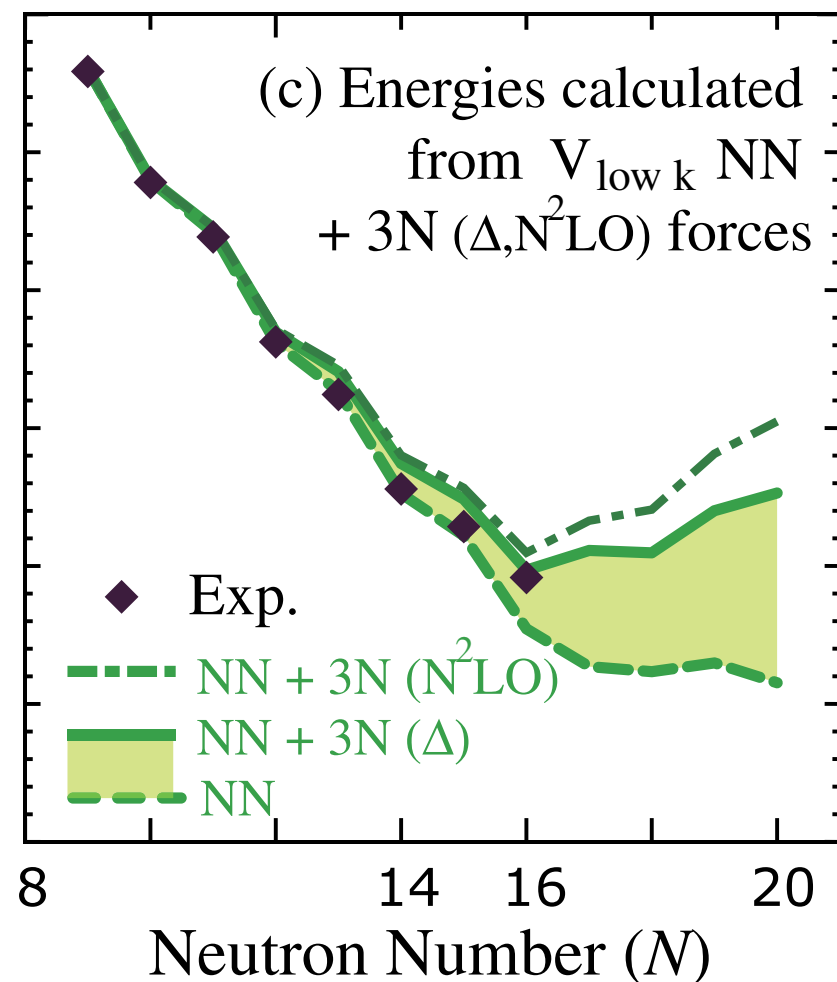
$E(Z, N) > E(Z, N+1, 2) \Rightarrow (Z, N+1, 2)$ は **unbound**

1つ、または2つ粒子を足した時、結合エネルギーを稼げるか？

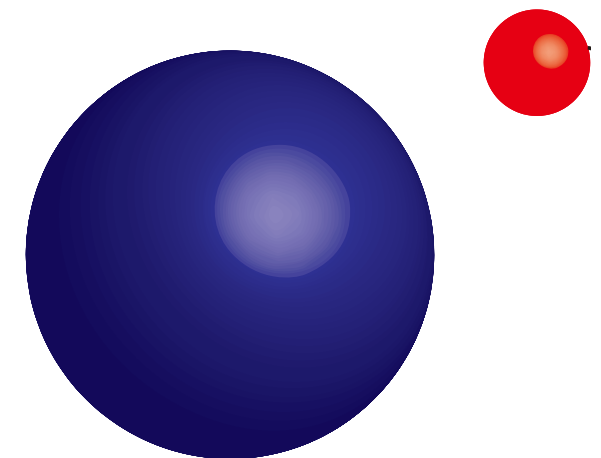
一般的な中性子ドリップラインの説明

Oxygen($Z=8$) case

T. Otsuka et al., Phys. Rev. Lett. 105, 032501 (2010).



positive ESPE

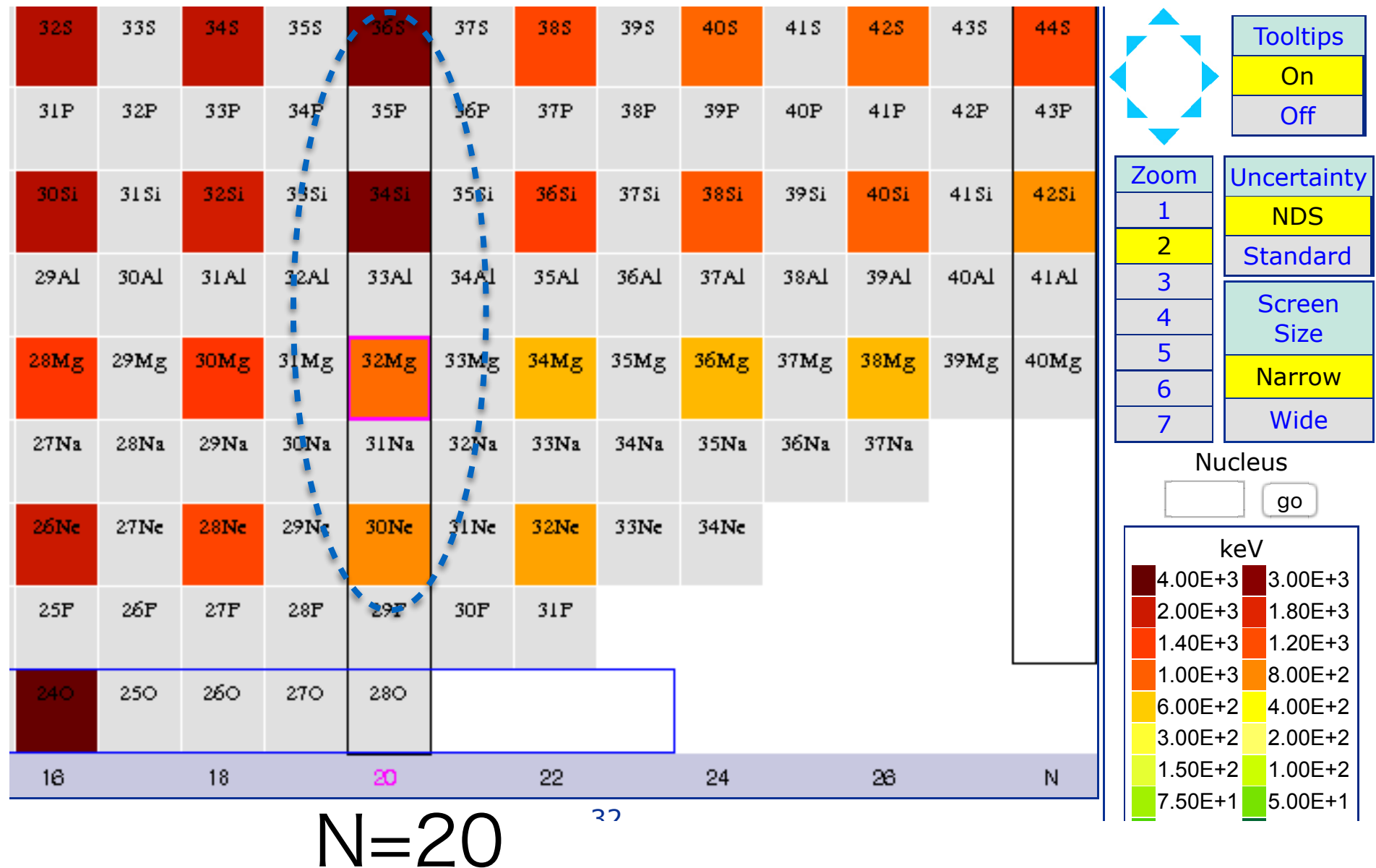


ESPE mostly determines drip line

Neutron-rich nuclei~ island of inversion

<http://www.nndc.bnl.gov/nudat2/reCenter.jsp?z=12&n=20>

shell gap
↑
breaking of
shell gap

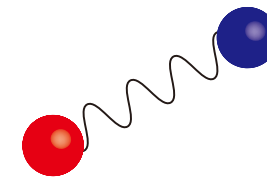


- $E(2^+) \sim 1$ MeV on $N=20$ indicate **breaking** of major shell gap
- Unified treatment of beyond and below the $N=20$ gap is necessary

Many body problem

Original Hamiltonian

$$H = T + \sum_{i,j=1}^N V_{NN} + \sum_{i,j,k}^N V_{NNN}$$



Shell model Hamiltonian

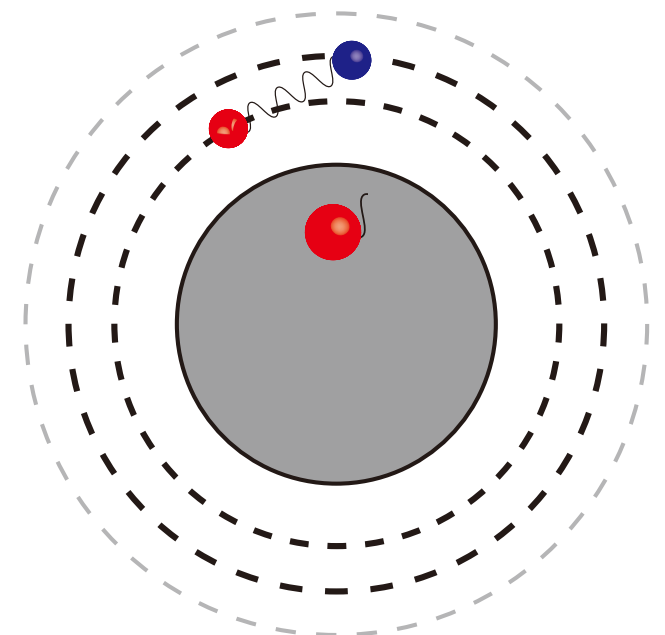
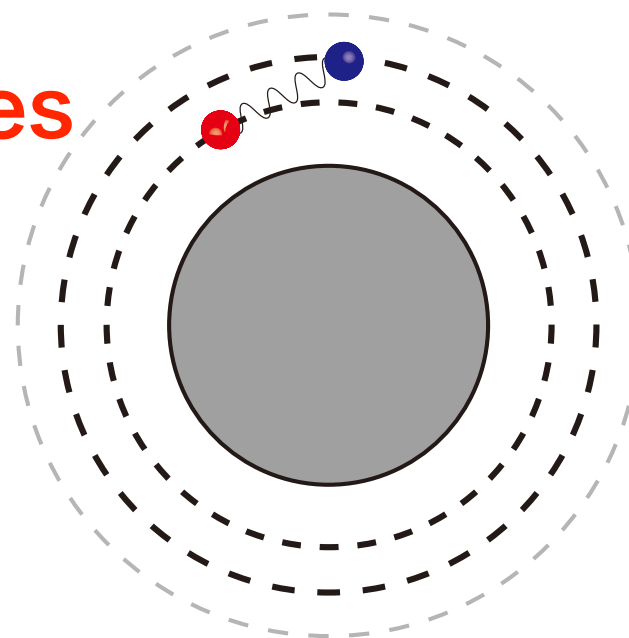
NN force

effective NN
from 3N

Single particle energies

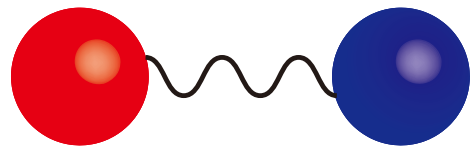
$$H = \sum_i \epsilon_i a_i^\dagger a_i + \sum_{ijkl} V_{ij,kl} a_i^\dagger a_j^\dagger a_l a_k.$$

Two-body matrix elements

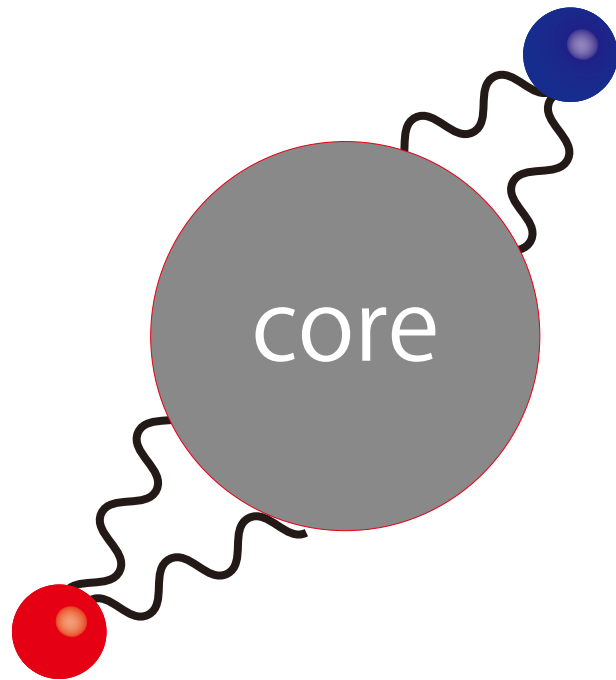


Effective interaction

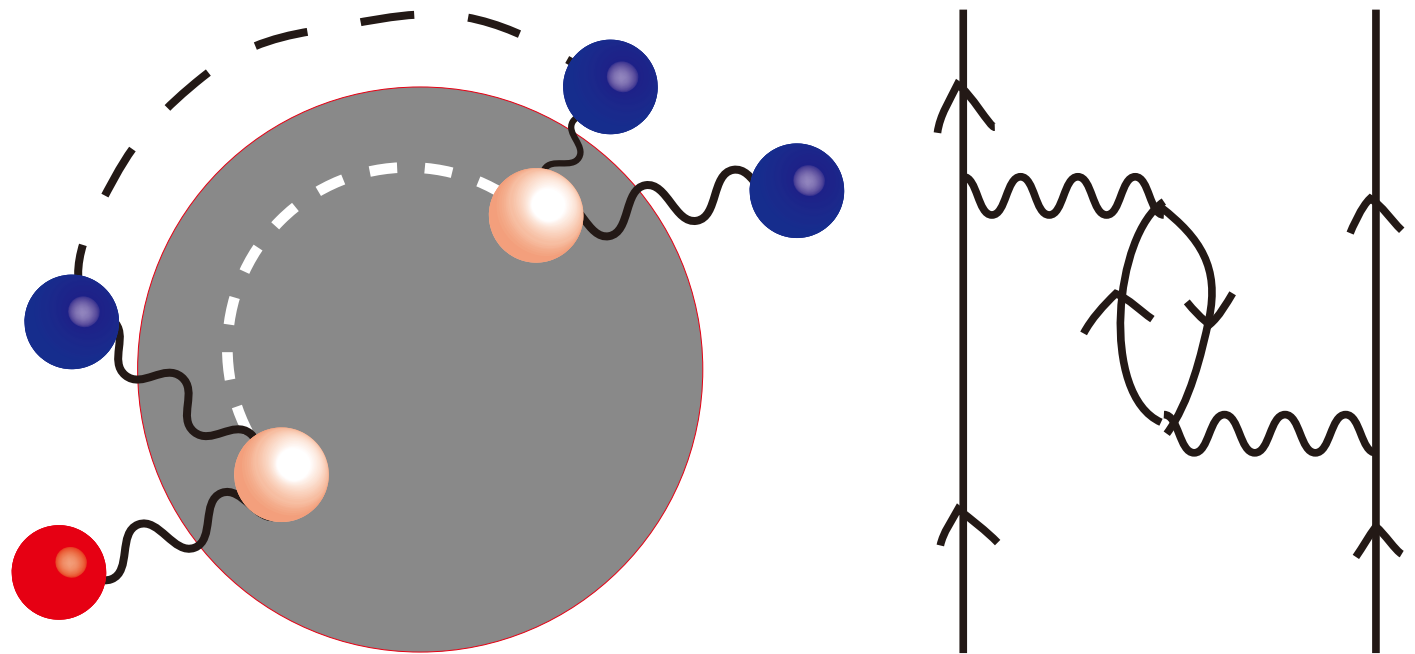
Nuclear force in vacuum



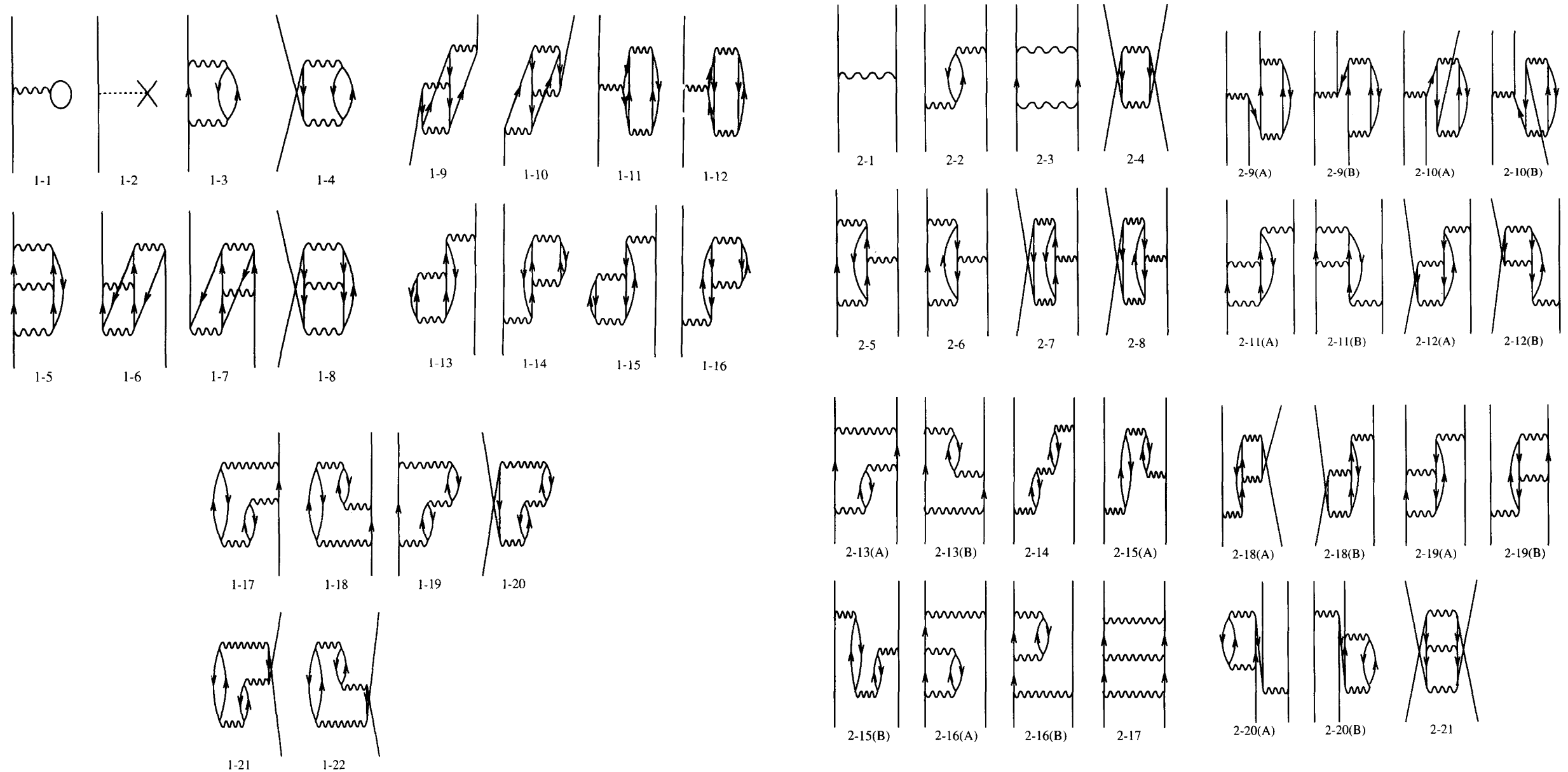
Medium effect



Core polarization (3p1h)



Many-body perturbation theory



Extended KK method and conventional KK method

EKK method

New parameter E (arbitrary parameter)

$$H = H'_0 + V'$$

$$= \begin{pmatrix} E & 0 \\ 0 & QH_0Q \end{pmatrix} + \begin{pmatrix} P\tilde{H}P & PVQ \\ QVP & QVQ \end{pmatrix},$$

$$H_{\text{BH}}(E) = PHP + PVQ \frac{1}{E - QH_0Q} QVP.$$

$$\tilde{H}_{\text{eff}}^{(n)} = \tilde{H}_{\text{BH}}(E) + \sum_{k=1}^{\infty} \hat{Q}_k(E) \{\tilde{H}_{\text{eff}}^{(n-1)}\}^k.$$

KK method (conventional)

$$H = H_0 + V$$

$$= \begin{pmatrix} PH_0P & 0 \\ 0 & QH_0Q \end{pmatrix} + \begin{pmatrix} PVP & PVQ \\ QVP & QVQ \end{pmatrix}$$

$$\hat{Q}(E) = PVP + PVQ \frac{1}{E - QH_0Q} QVP$$

$$V_{\text{eff}}^{(n)} = \hat{Q}(\epsilon_0) + \sum_{k=1}^{\infty} \hat{Q}_k(\epsilon_0) \{V_{\text{eff}}^{(n-1)}\}^k.$$

- EKK method enable us to construct effective interaction for multi-major shell

N. Tsunoda, K. Takayanagi, M. Hjorth-Jensen, and T. Otsuka, Phys. Rev. C 89, 024313 (2014).

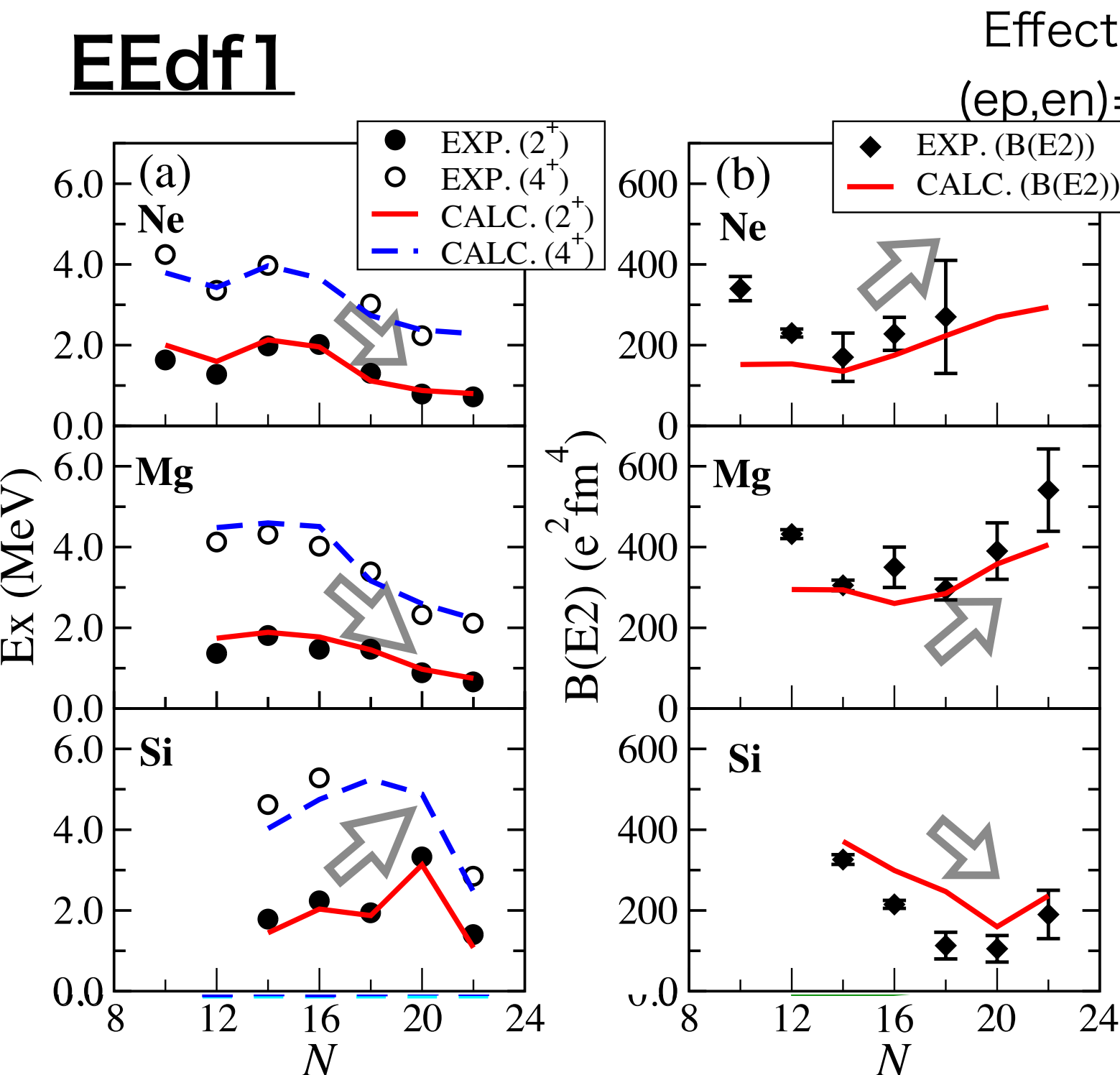
K. Takayanagi, Annals of Physics 350, 501 (2014).

K. Takayanagi, Nucl. Phys. A 852, 61 (2011).

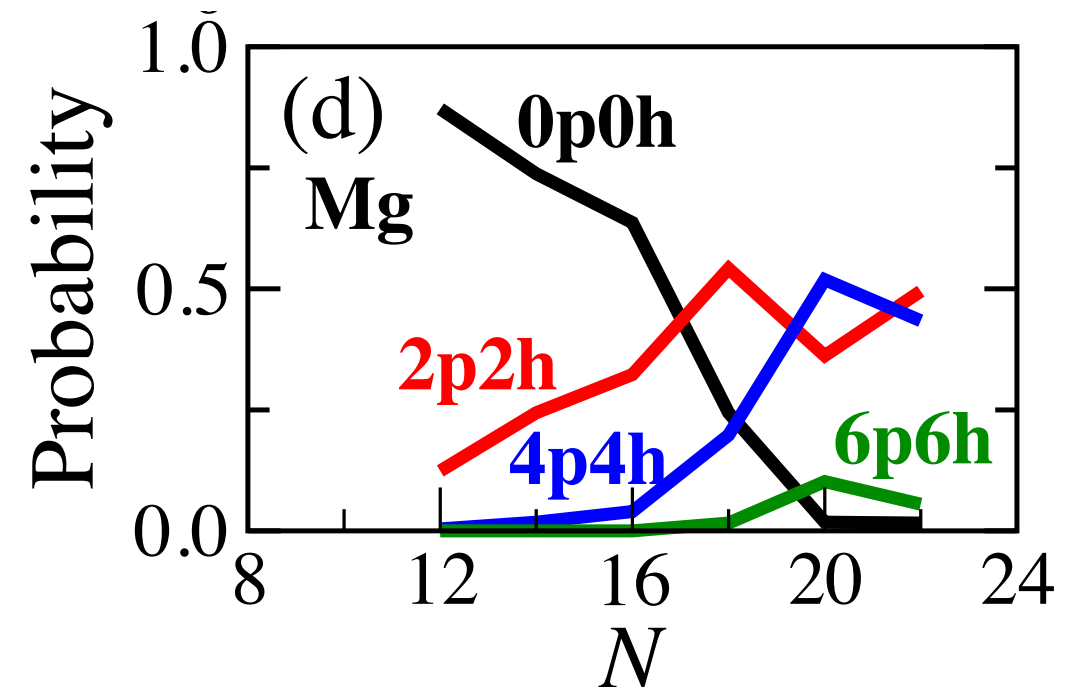
- Effective interaction for **island of inversion**
- Effective interaction designed for **sd+pf** shell
- TBMEs are determined by **EKK method**
- Effective 2NF from **3NF**(Fujita-Miyazawa type) force is added
- SPEs are fitted to experimental data

Shell structure in “island of inversion”

EEdf1



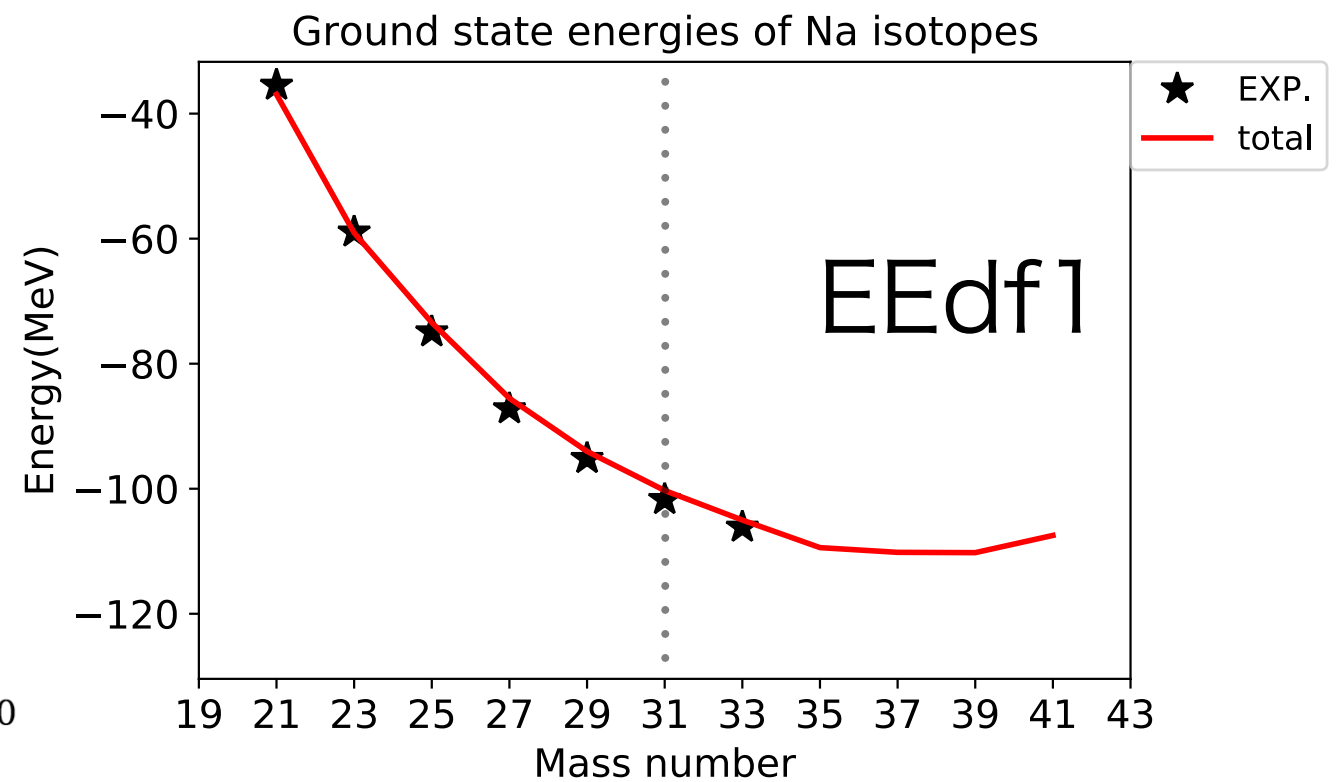
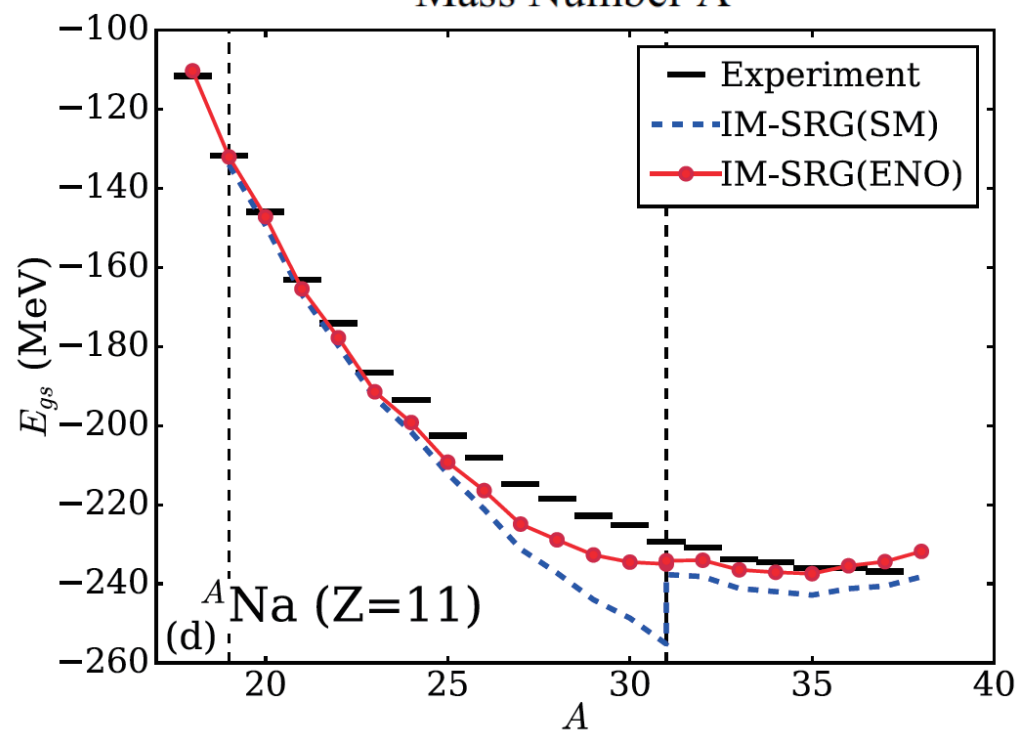
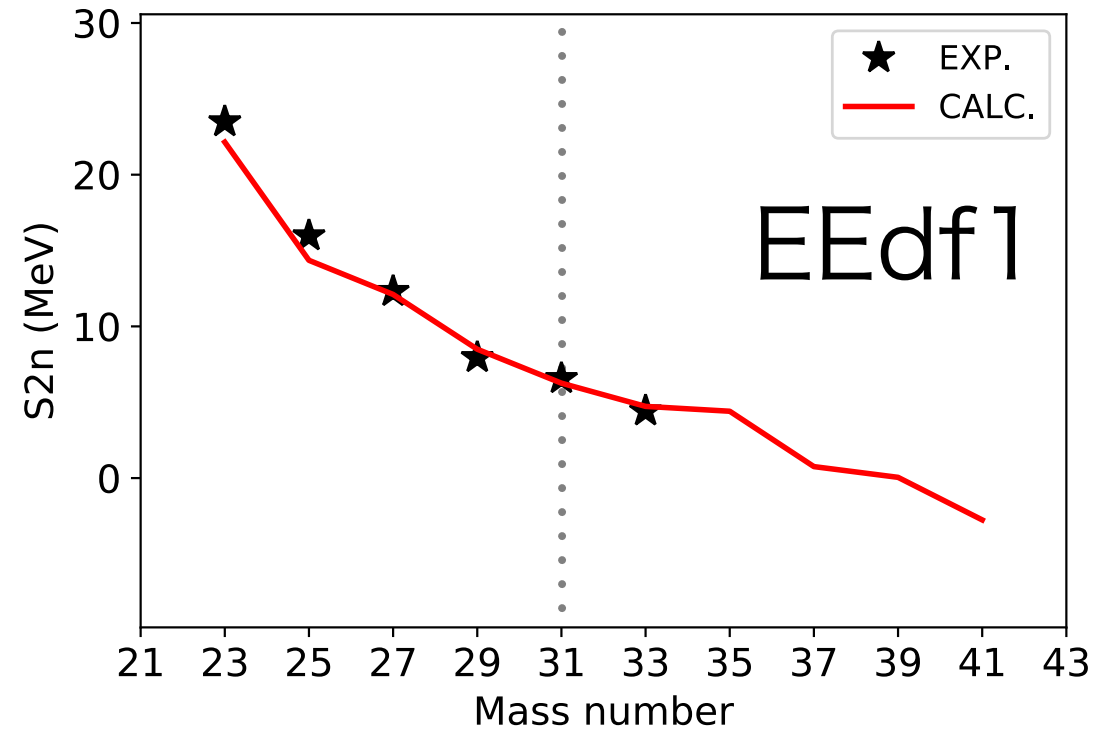
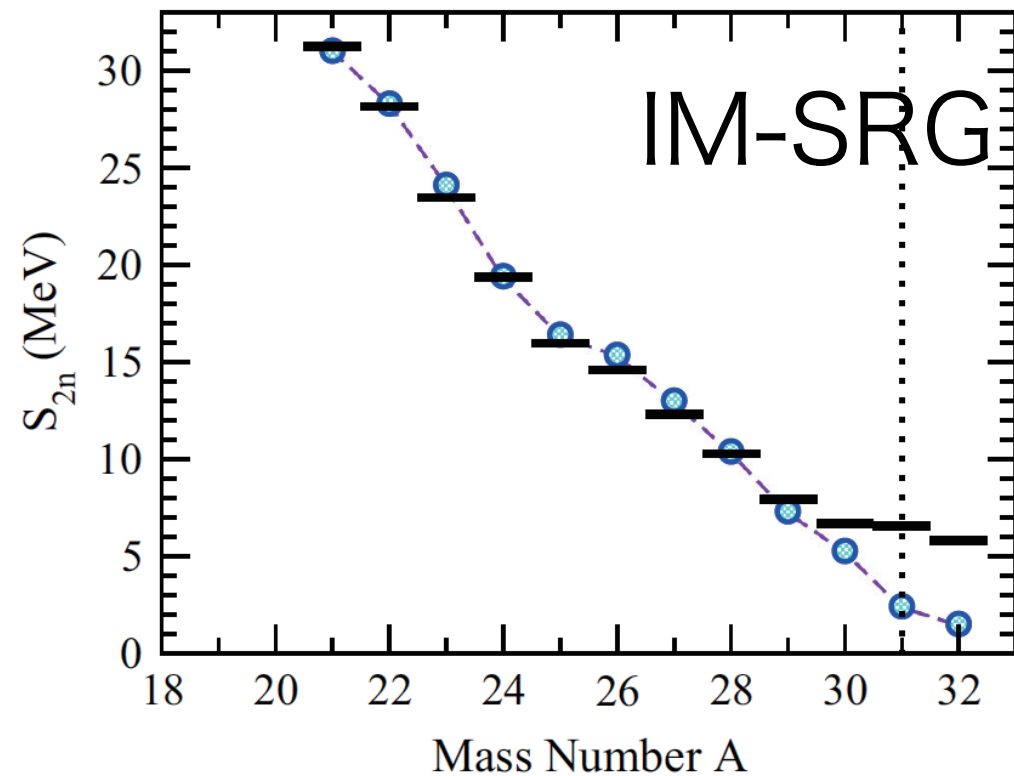
N.T, et al. Phys. Rev. C 95, 021304 (2017)



Similar ph excitation pattern to:
A. O. Macchiavelli, et al., Phys. Rev. C 94, 051303 (2016).

For larger N , we are now working on...
(e.g. 40Mg 2^+ by Crawford's talk)

Comparison to ab initio calculations



S. R. Stroberg, et al., Phys. Rev. Lett. 118, 032502 (2017).

J. Simonis, et al., Phys. Rev. C 96, 014303 (2017).

その他最新の実験との比較

- ^{27}Ne C. Leolius et al., accepted to PRL
- ^{30}Mg B. Fernandez-Dominguez et al., PLB **779**, 124(2018)
- ^{31}Na , ^{31}Mg H. Nishibata et al., 論文投稿中
- ^{32}Ne I. Murray et al., PRC accepted(2019)
- ^{34}Ne , ^{39}Na Deuksoon Ahn et al., 論文準備中
- ^{34}Al Z. Xu et al. PLB **782**, 619(2018)

Comparisons to experimental data are successful !

Anatomy of interaction

bare SPE $\sum \epsilon_i a_i^\dagger a_i$

monopole $\sum_{i,j} V_{\text{mono}}^{ab} a_i^\dagger a_j^\dagger a_j a_i$ $V_{\text{mono}}^{ab} = \sum_J \frac{(2J+1) \langle ab|V|ab \rangle_J}{2J+1}$

pairing J=0 (monopole removed)

multipole other (QQ etc.)

Anatomy of interaction

bare SPE

$$\sum \epsilon_i a_i^\dagger a_i$$

bare and effective SPE

monopole

$$\sum_{i,j} V_{\text{mono}}^{ab} a_i^\dagger a_j^\dagger a_j a_i$$

$$V_{\text{mono}}^{ab} = \sum_J \frac{(2J+1) \langle ab|V|ab \rangle_J}{2J+1}$$

pairing

J=0 (monopole removed)

multipole

other (QQ etc.)

Anatomy of interaction

bare SPE

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bare and effective SPE

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pairing

J=0 (monopole removed)

correlation

multipole

other (QQ etc.)

Anatomy of interaction

bare SPE

$$\sum \epsilon_i a_i^\dagger a_i$$

bare and effective SPE

monopole

$$\sum_{i,j} V_{\text{mono}}^{ab} a_i^\dagger a_j^\dagger a_j a_i$$

$$V_{\text{mono}}^{ab} = \sum_J \frac{(2J+1) \langle ab|V|ab \rangle_J}{2J+1}$$

pairing

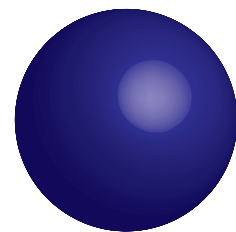
J=0 (monopole removed)

correlation

multipole

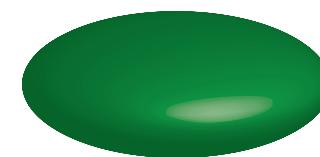
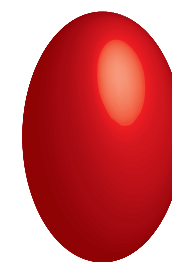
other (QQ etc.)

including deformation



monopole

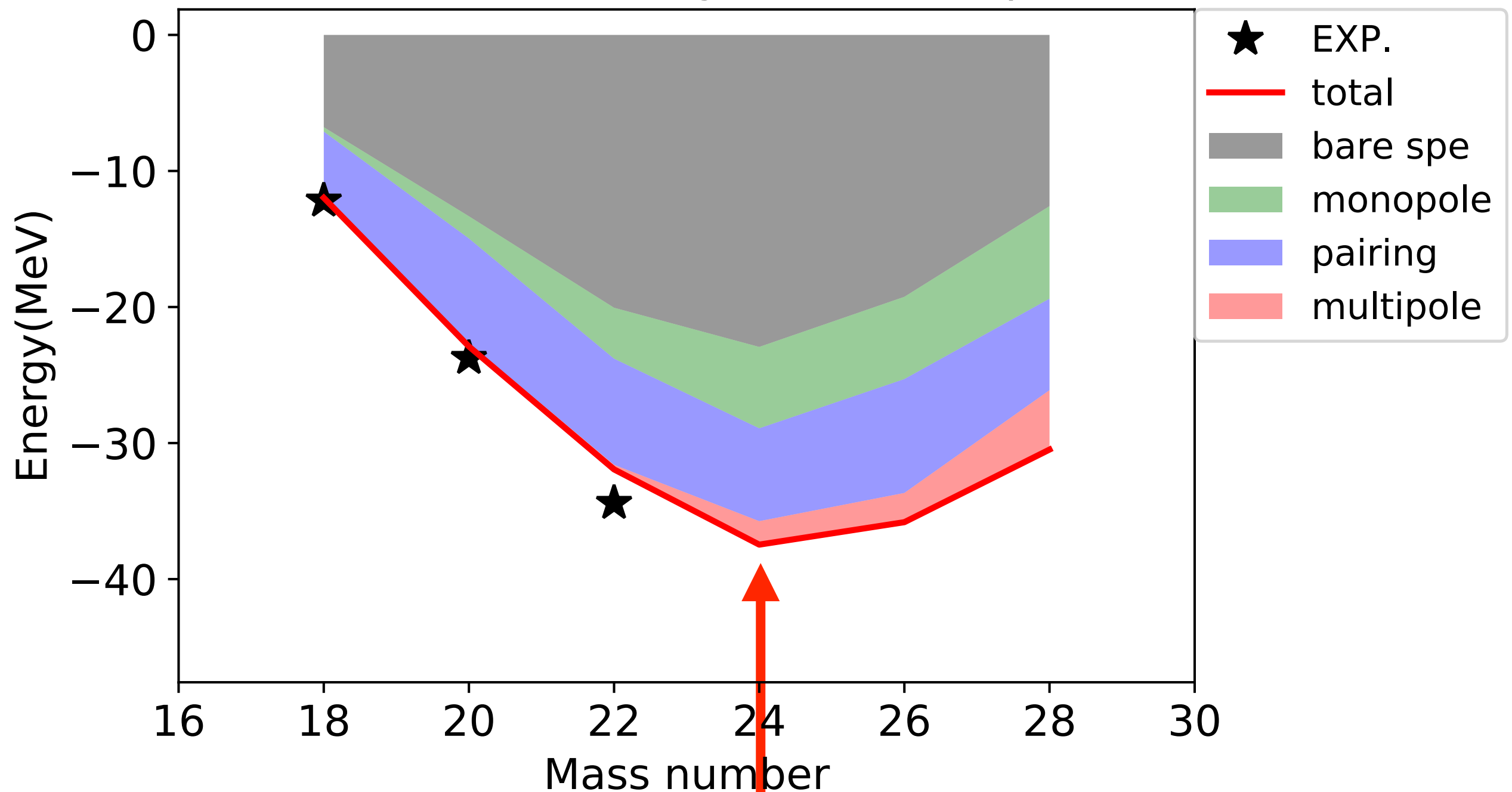
multipole



O isotope ($Z=8$)

(all the SPEs shifted by 0.9 MeV)

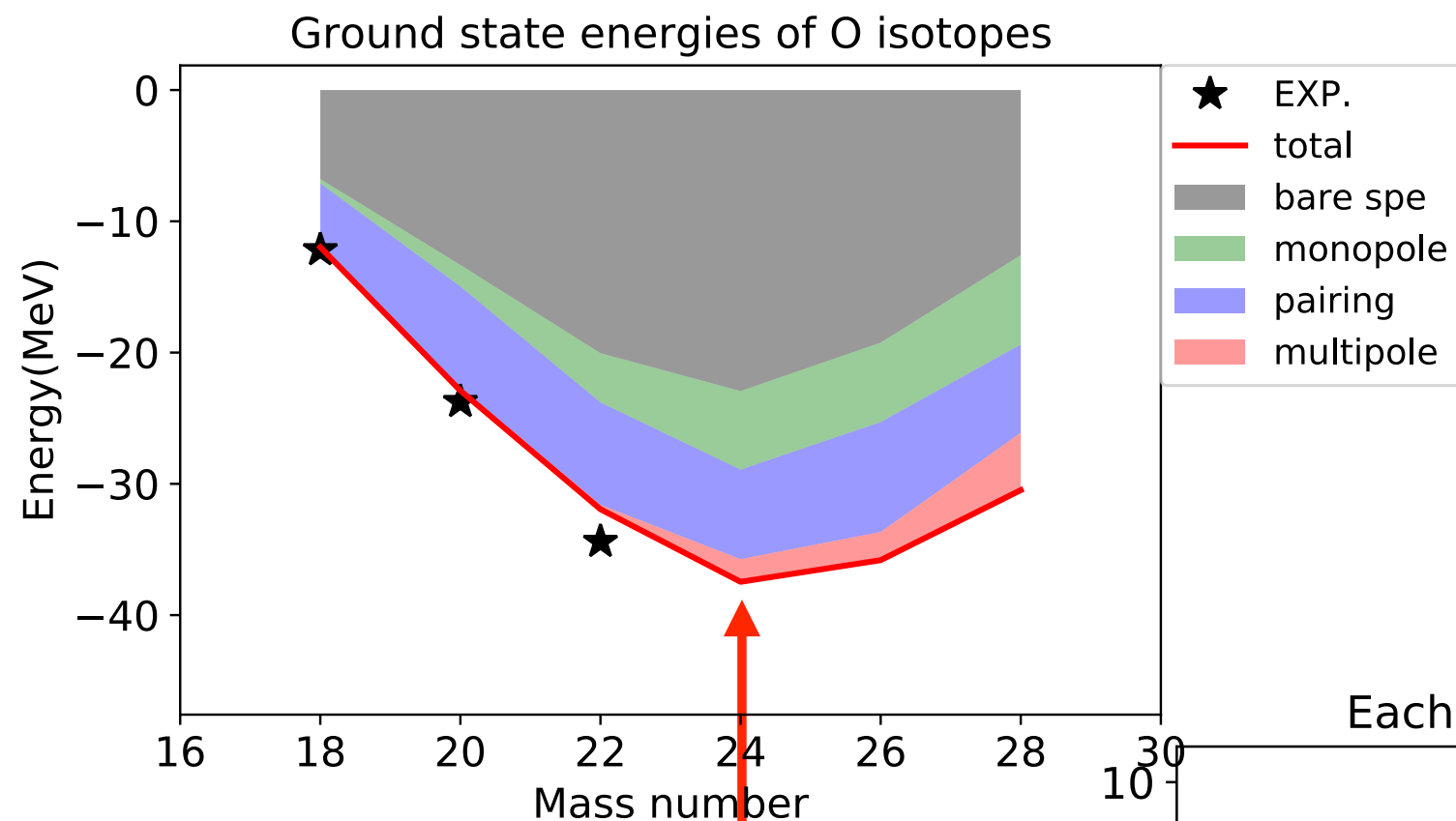
Ground state energies of O isotopes



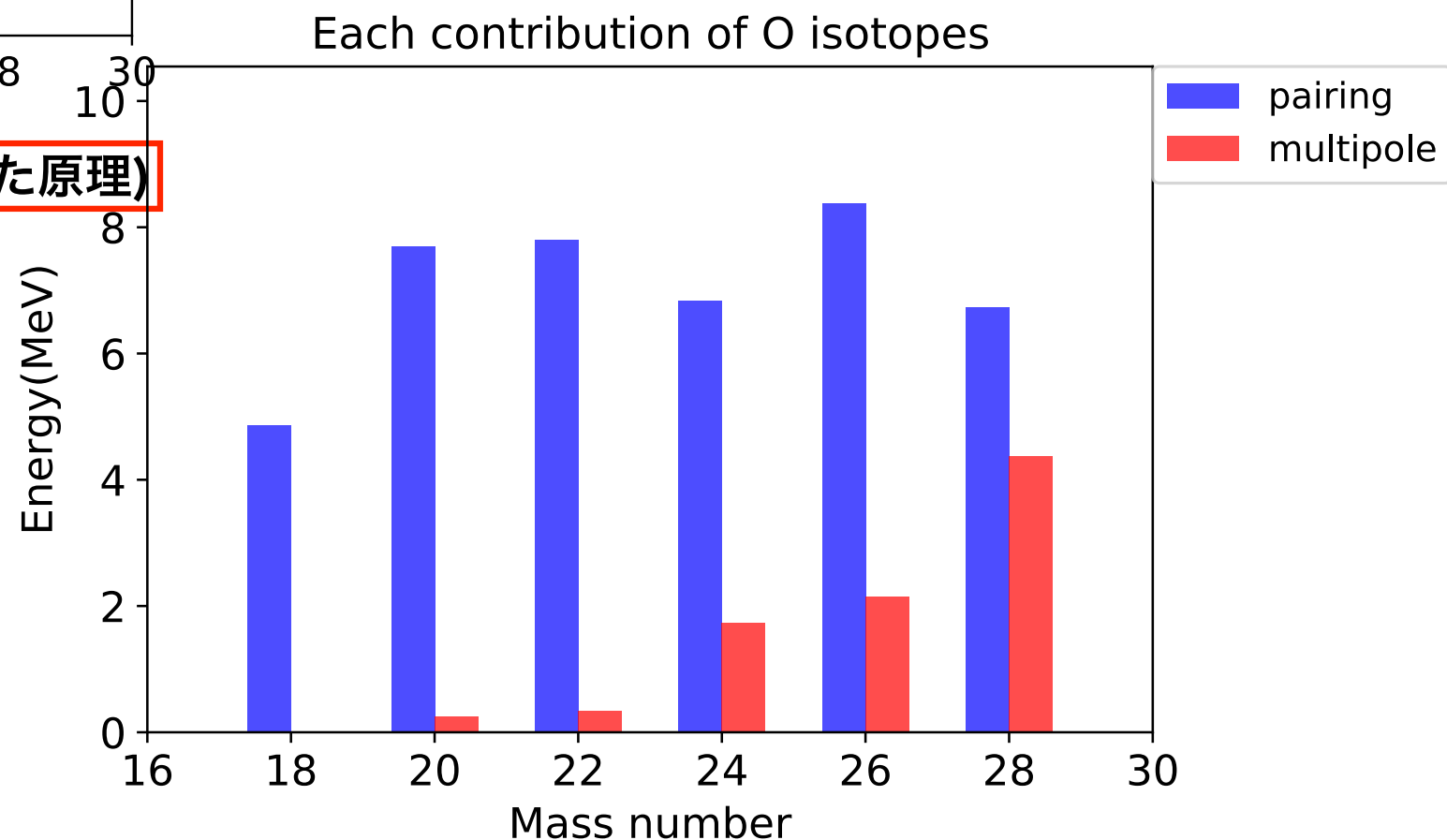
Dripline 24O (今まで知られてきた原理)

O isotope ($Z=8$)

(all the SPEs shifted by 0.9 MeV)



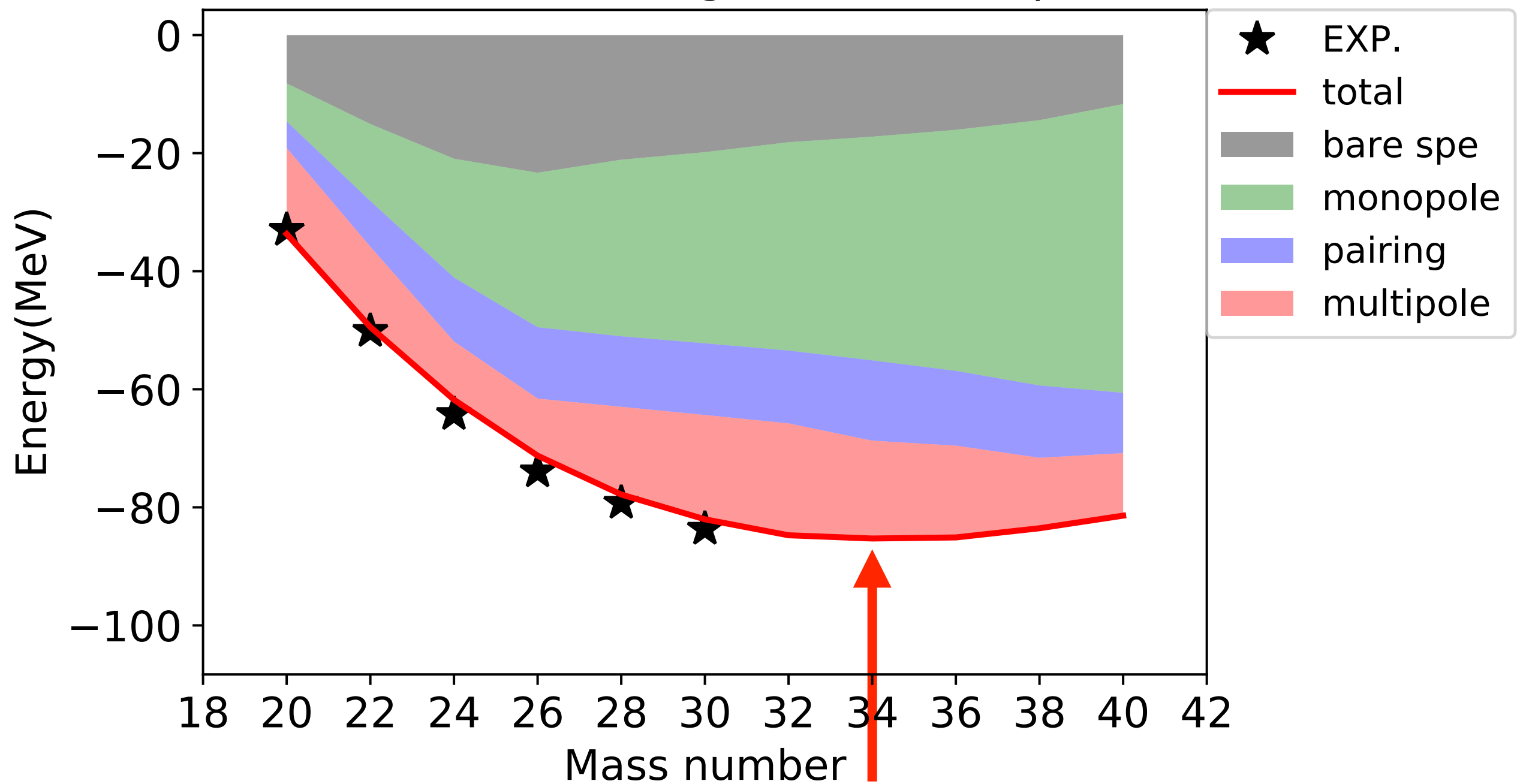
Dripline 24O (今まで知られてきた原理)



Ne isotope ($Z=10$)

(all the SPEs shifted by 0.9 MeV)

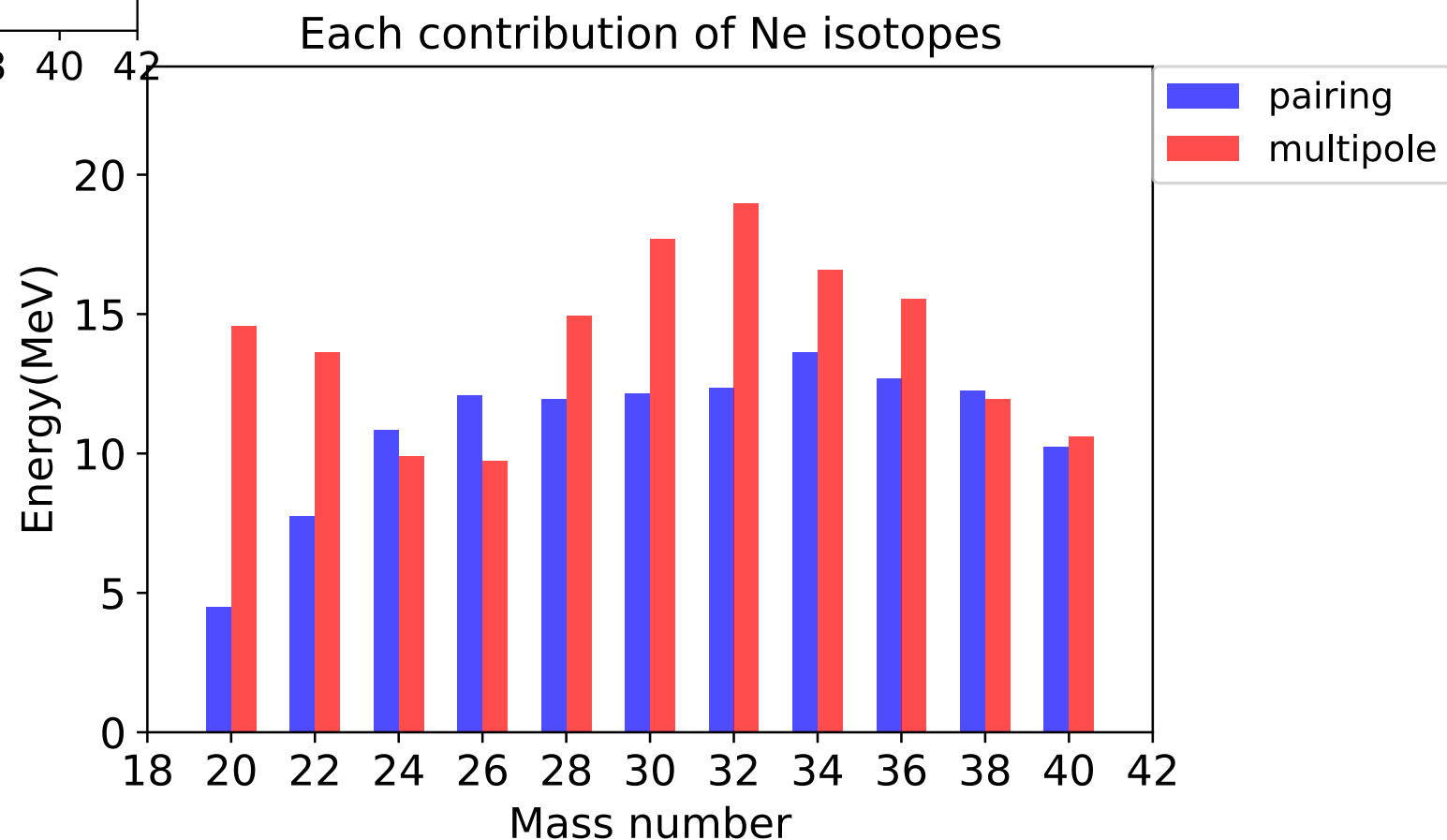
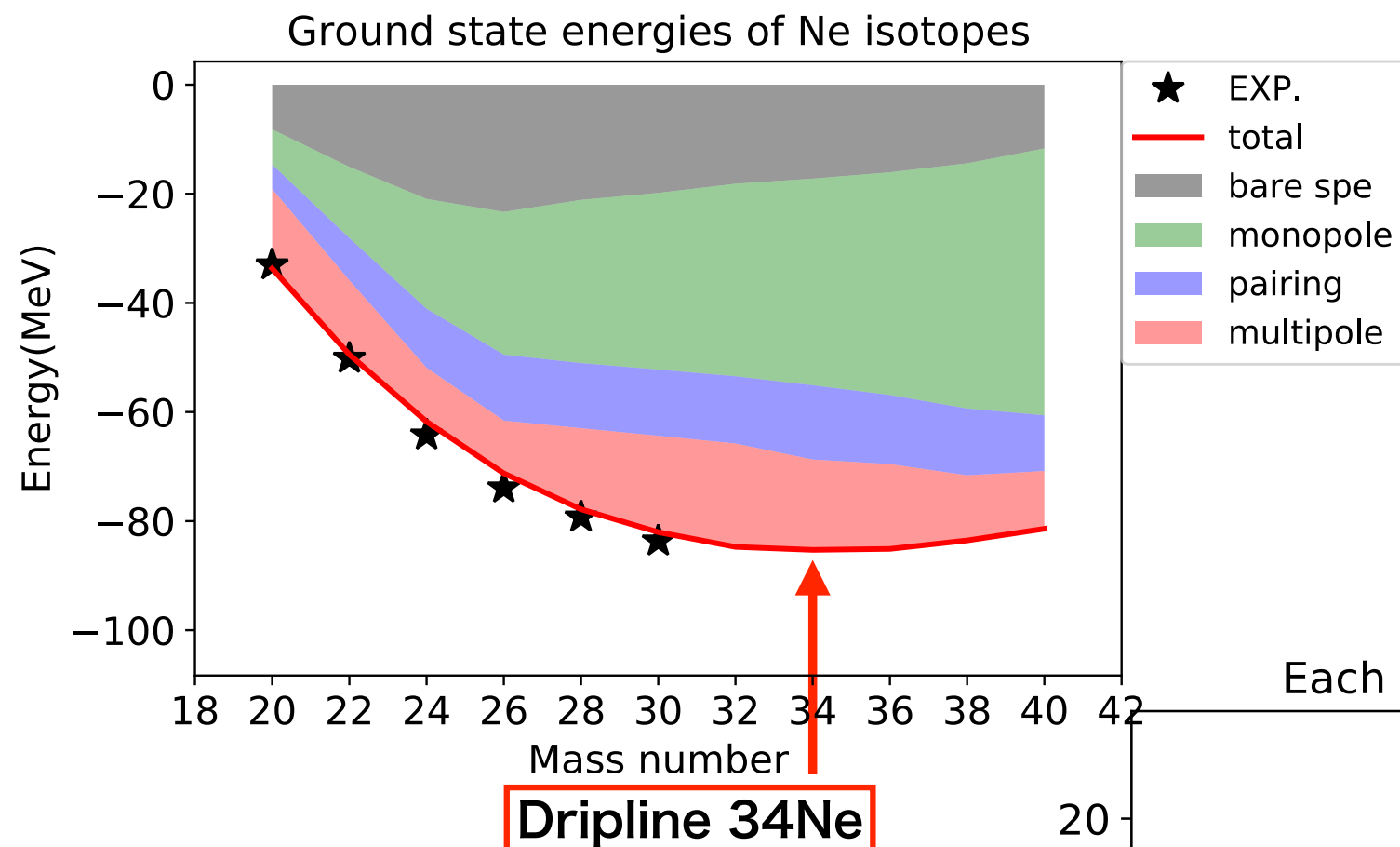
Ground state energies of Ne isotopes



Dripline ^{34}Ne

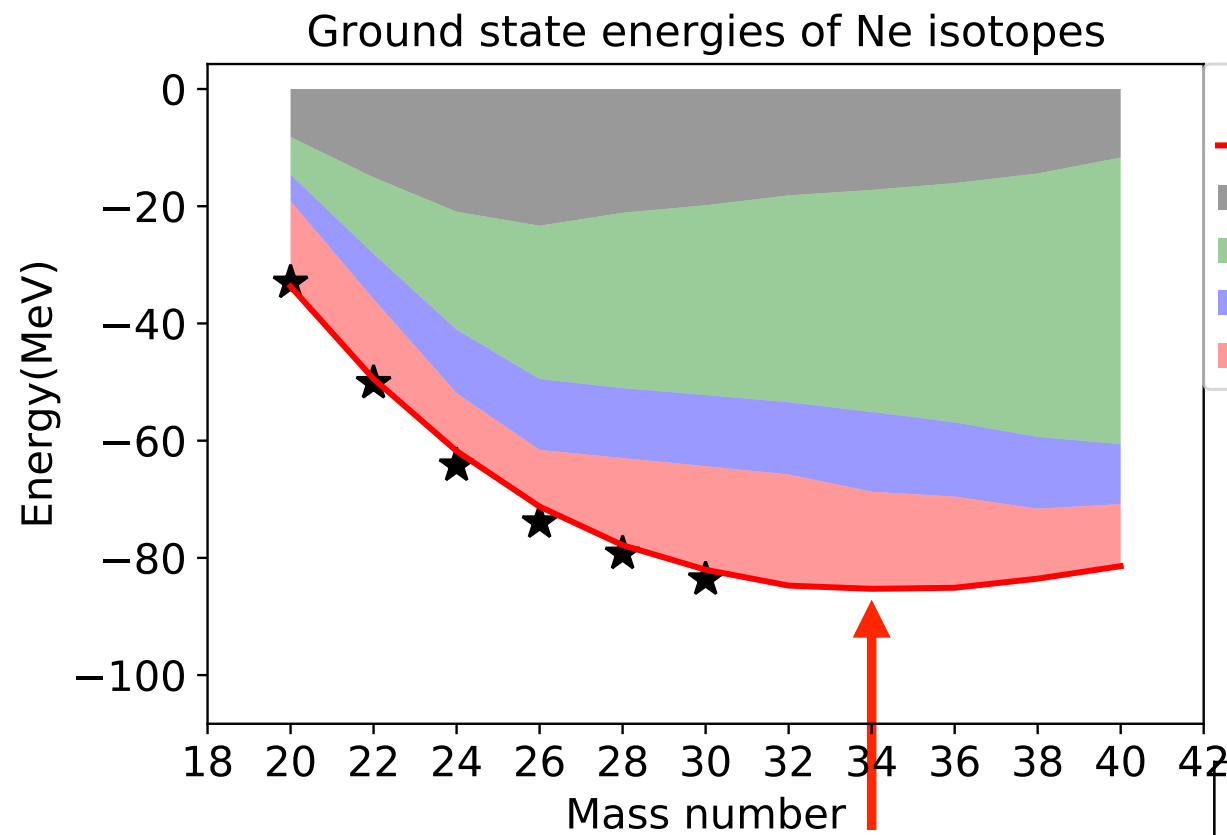
Ne isotope ($Z=10$)

(all the SPEs shifted by 0.9 MeV)

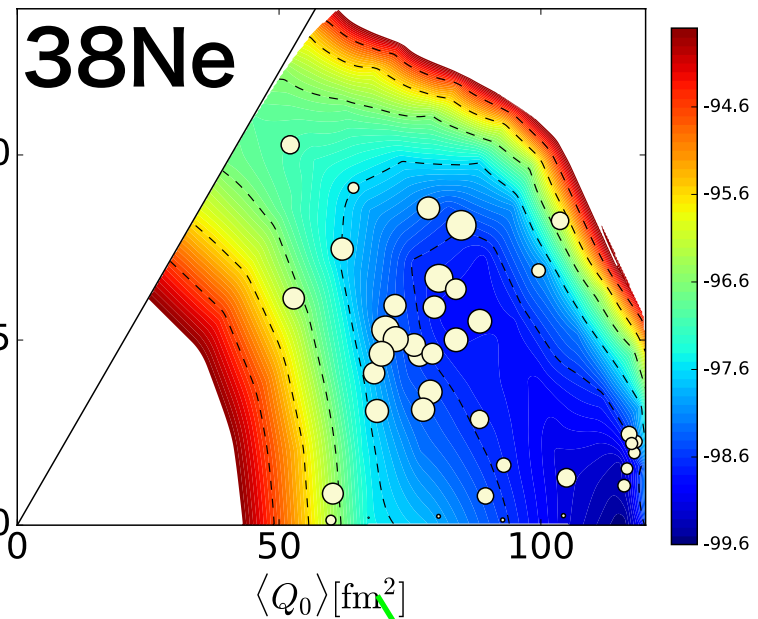


Ne isotope (Z=10)

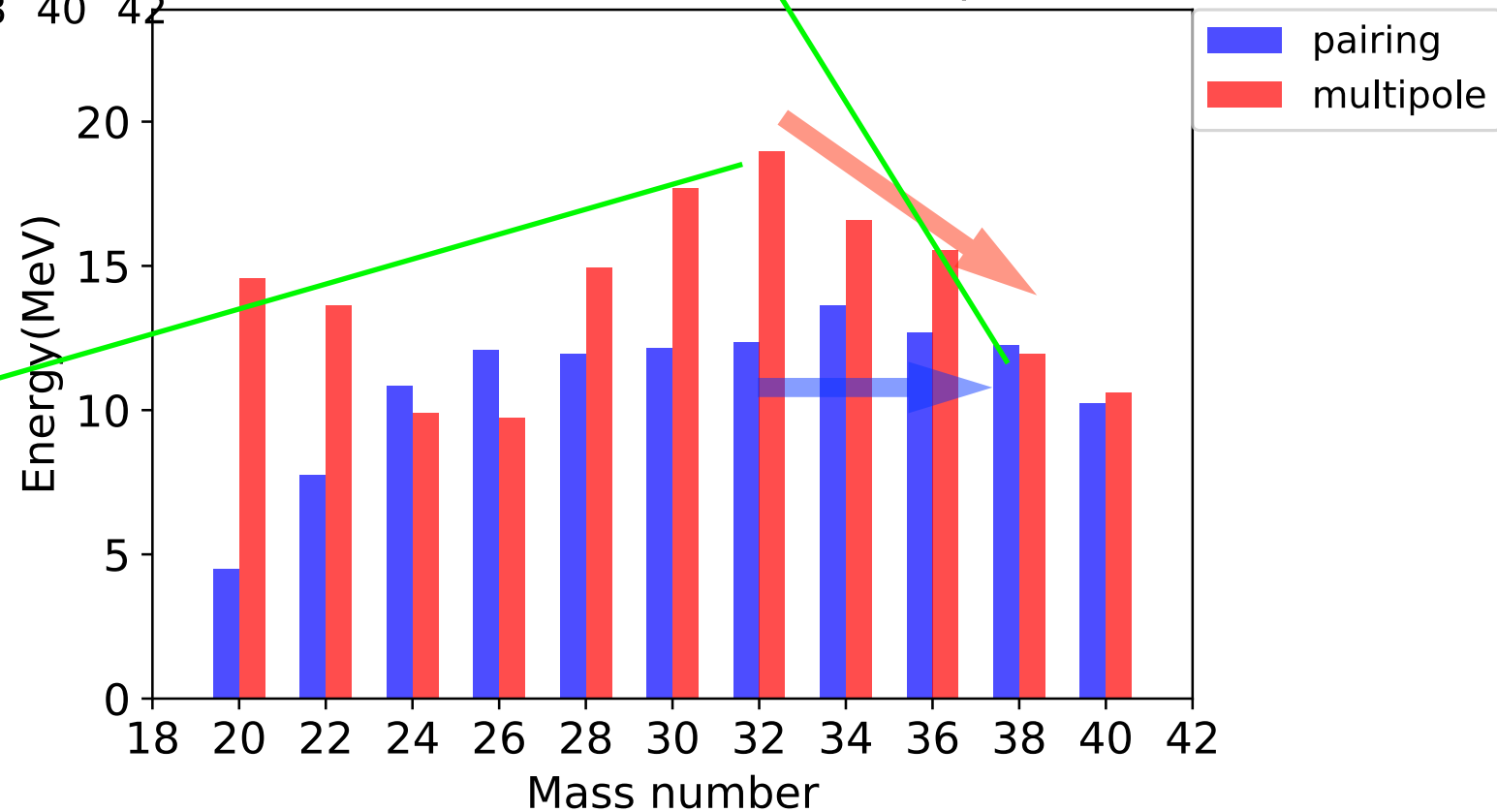
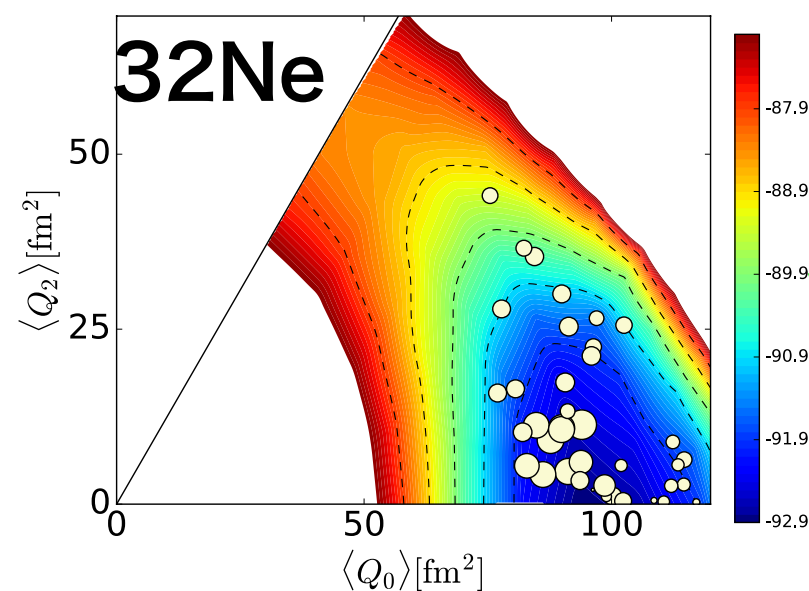
(all the SPEs shifted by 0.9 MeV)



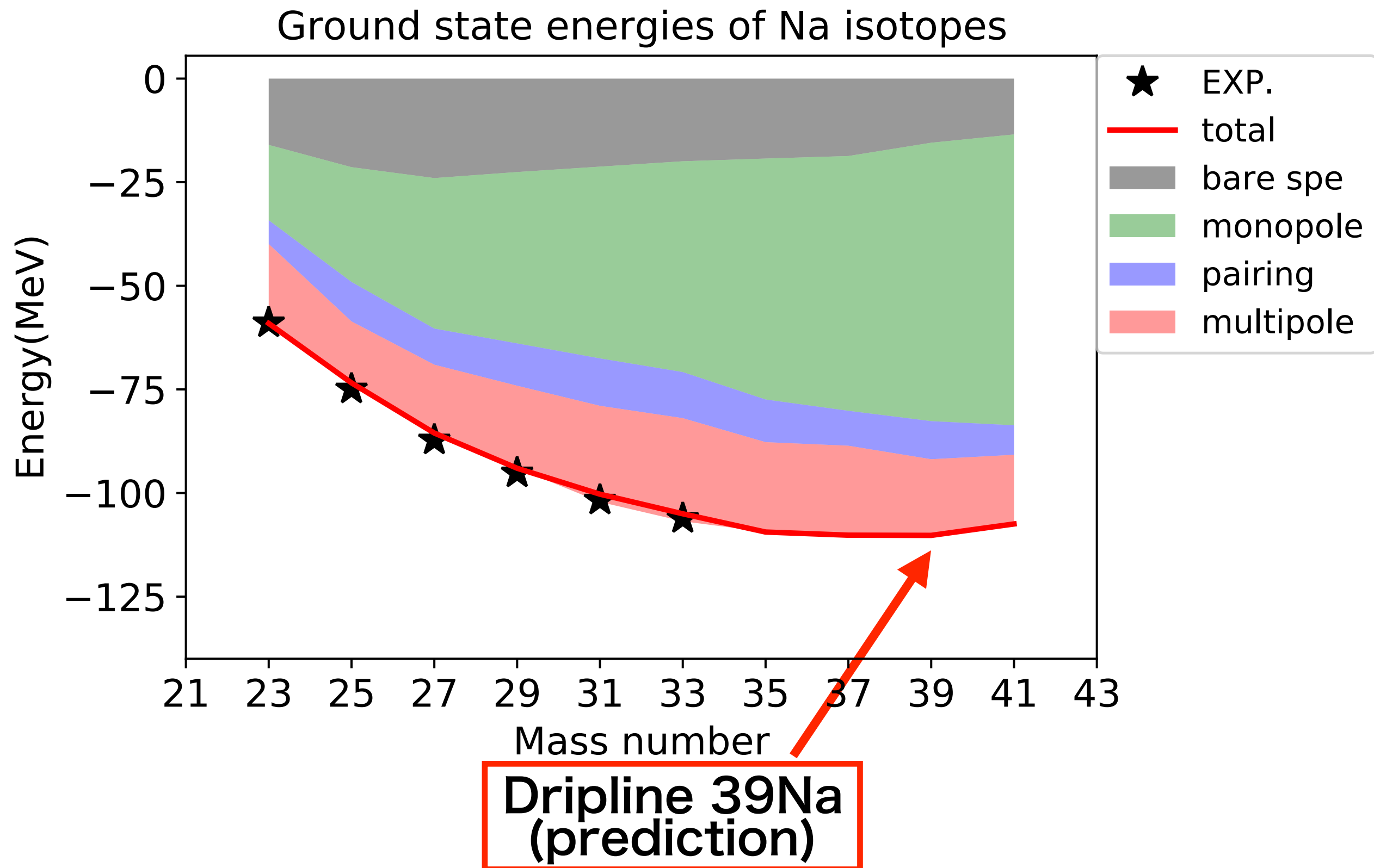
Dripline 34Ne



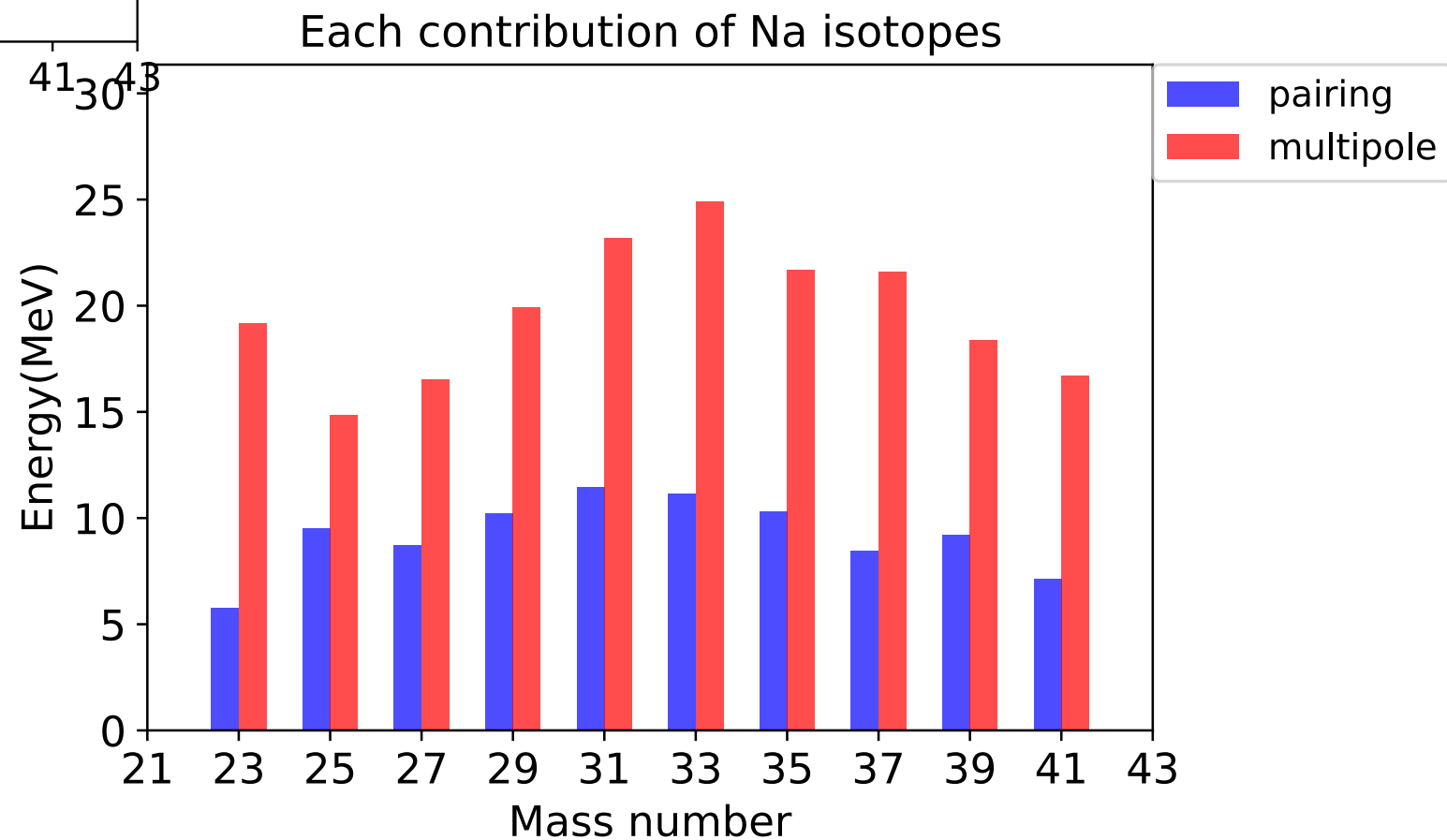
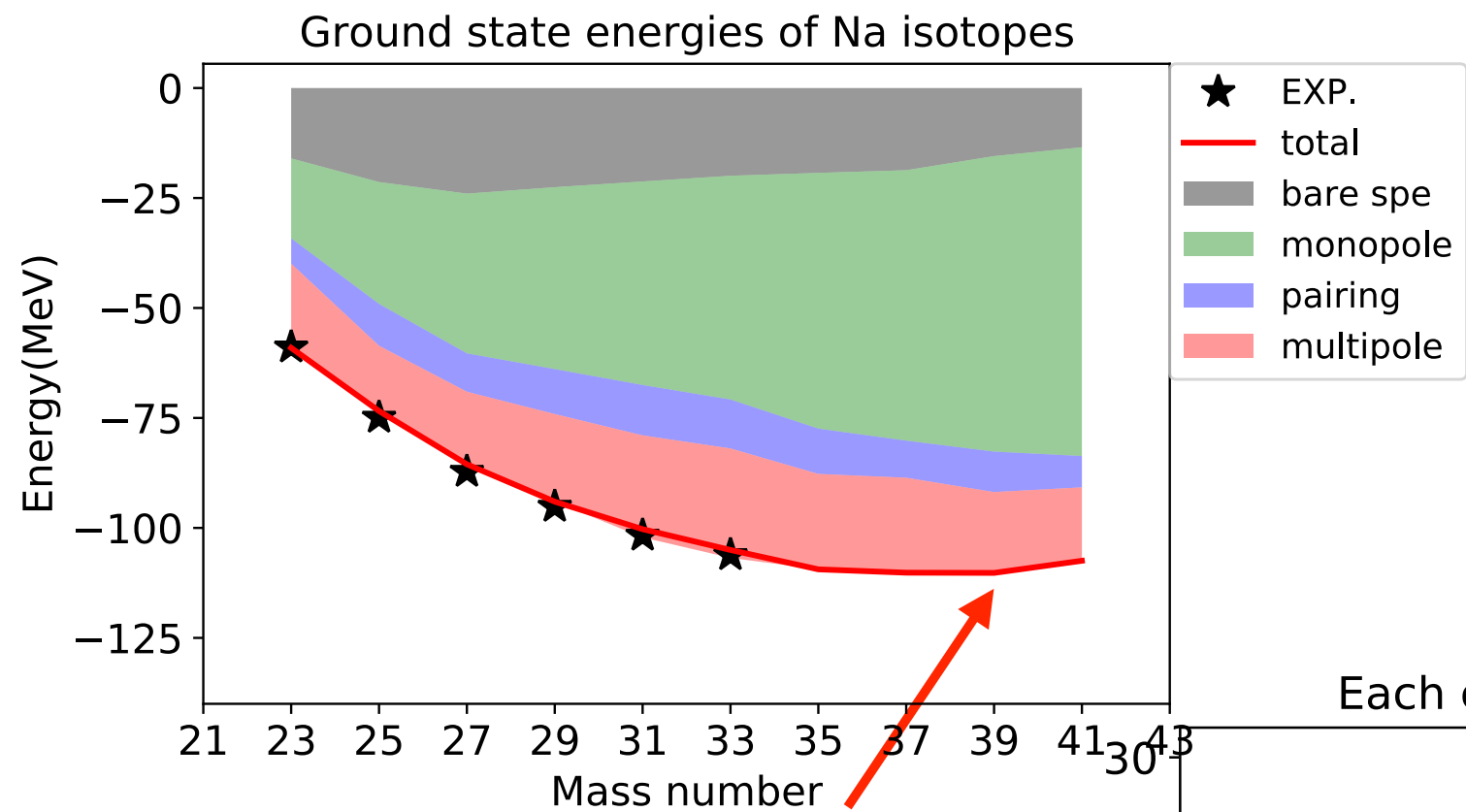
Each contribution of Ne isotopes



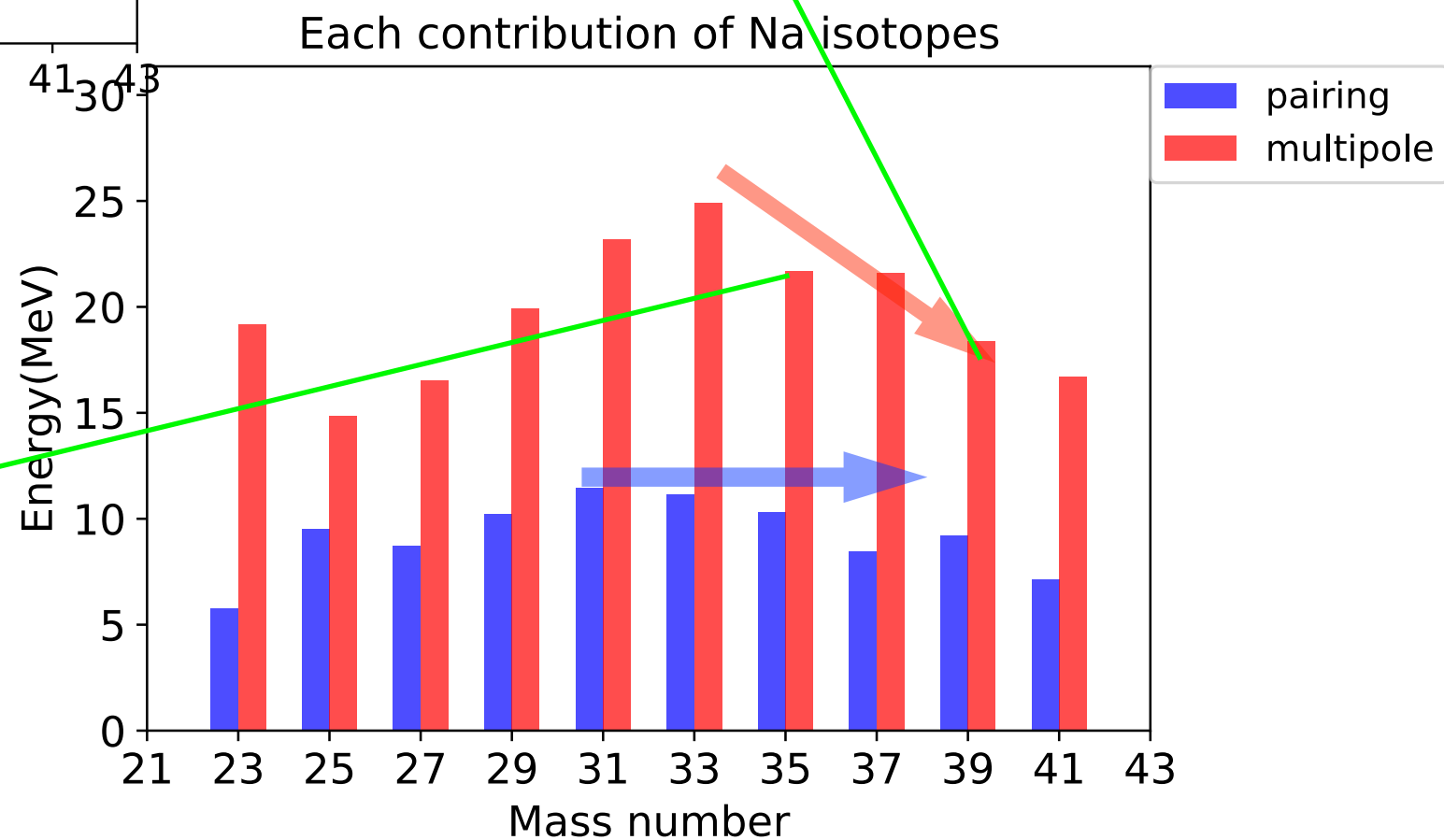
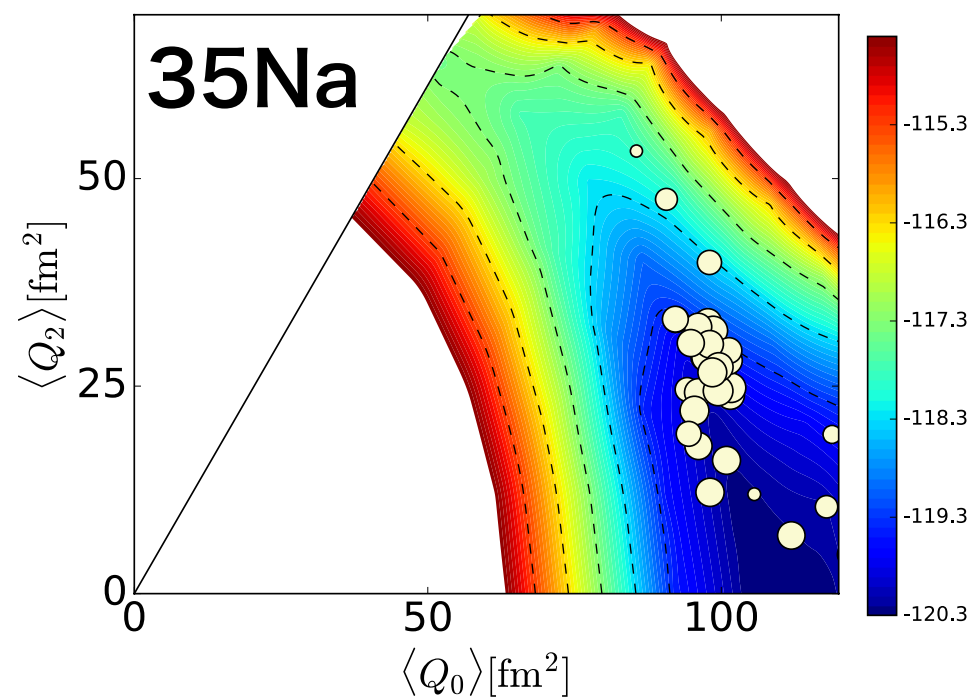
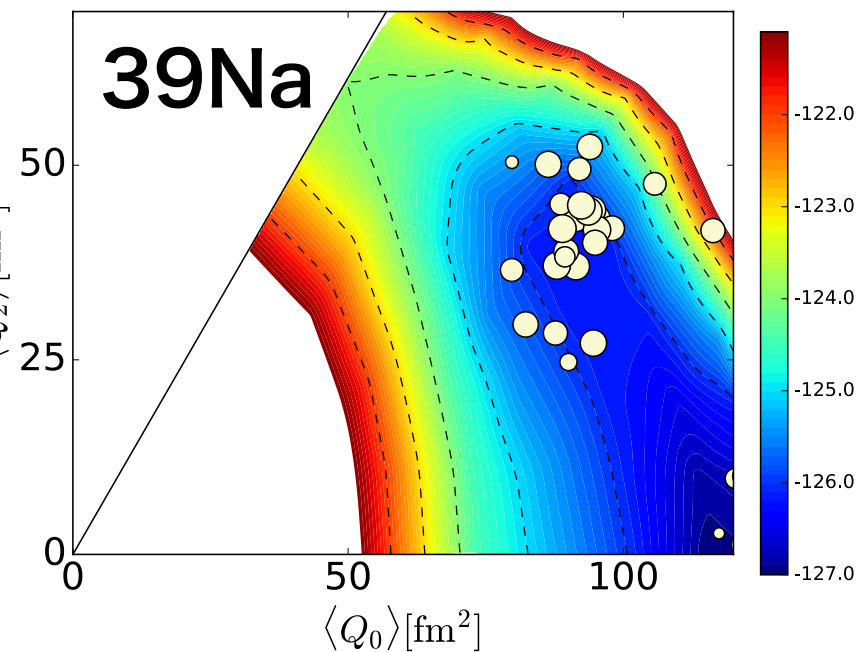
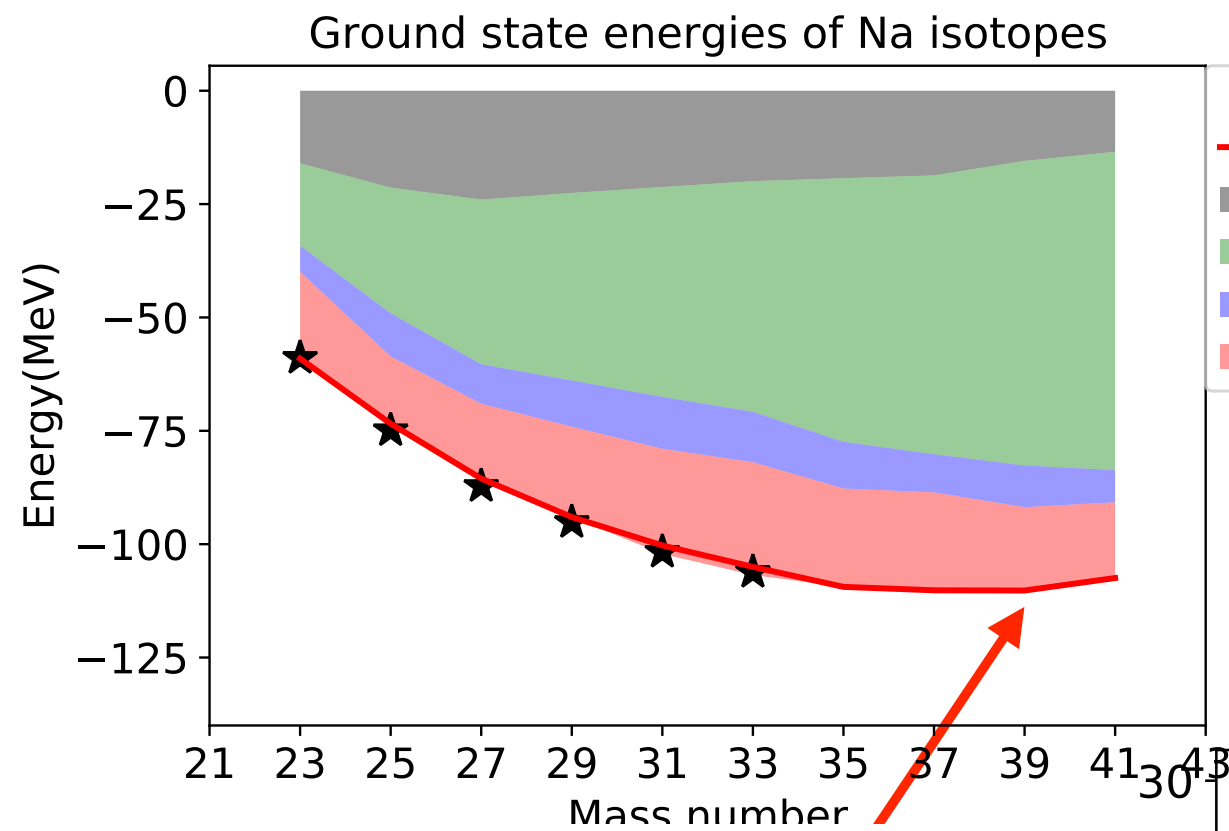
Na isotope ($Z=11$)



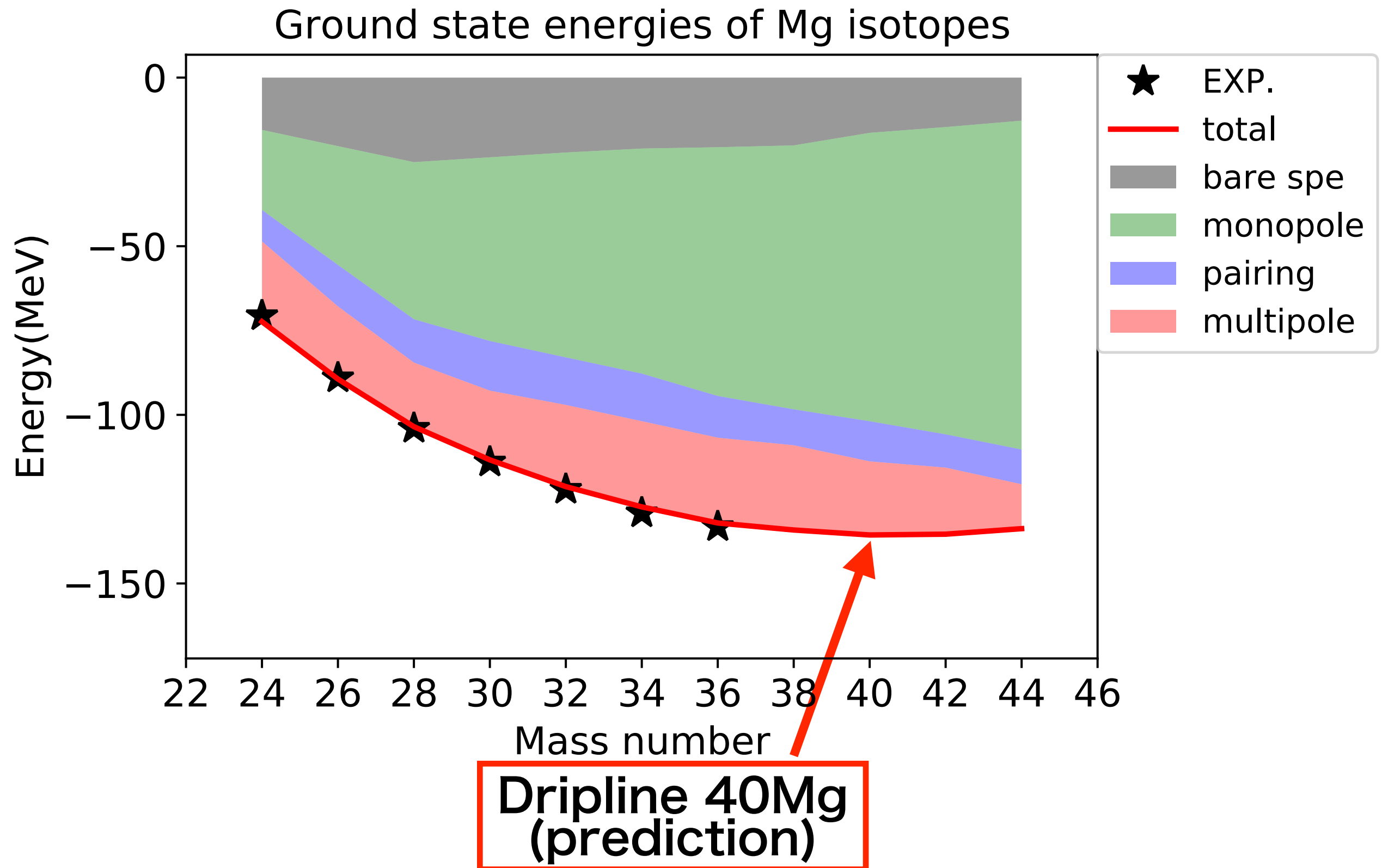
Na isotope (Z=11)



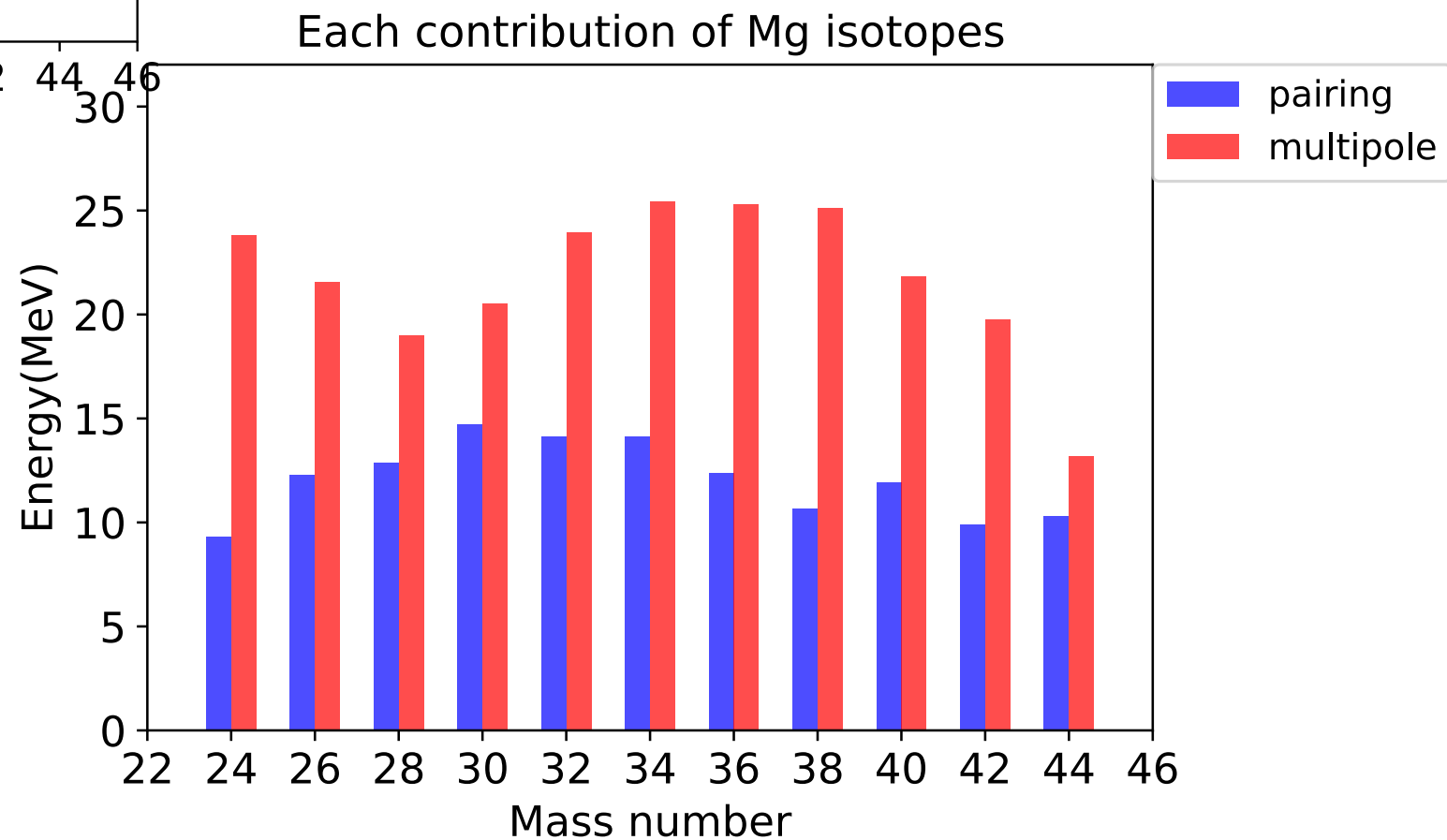
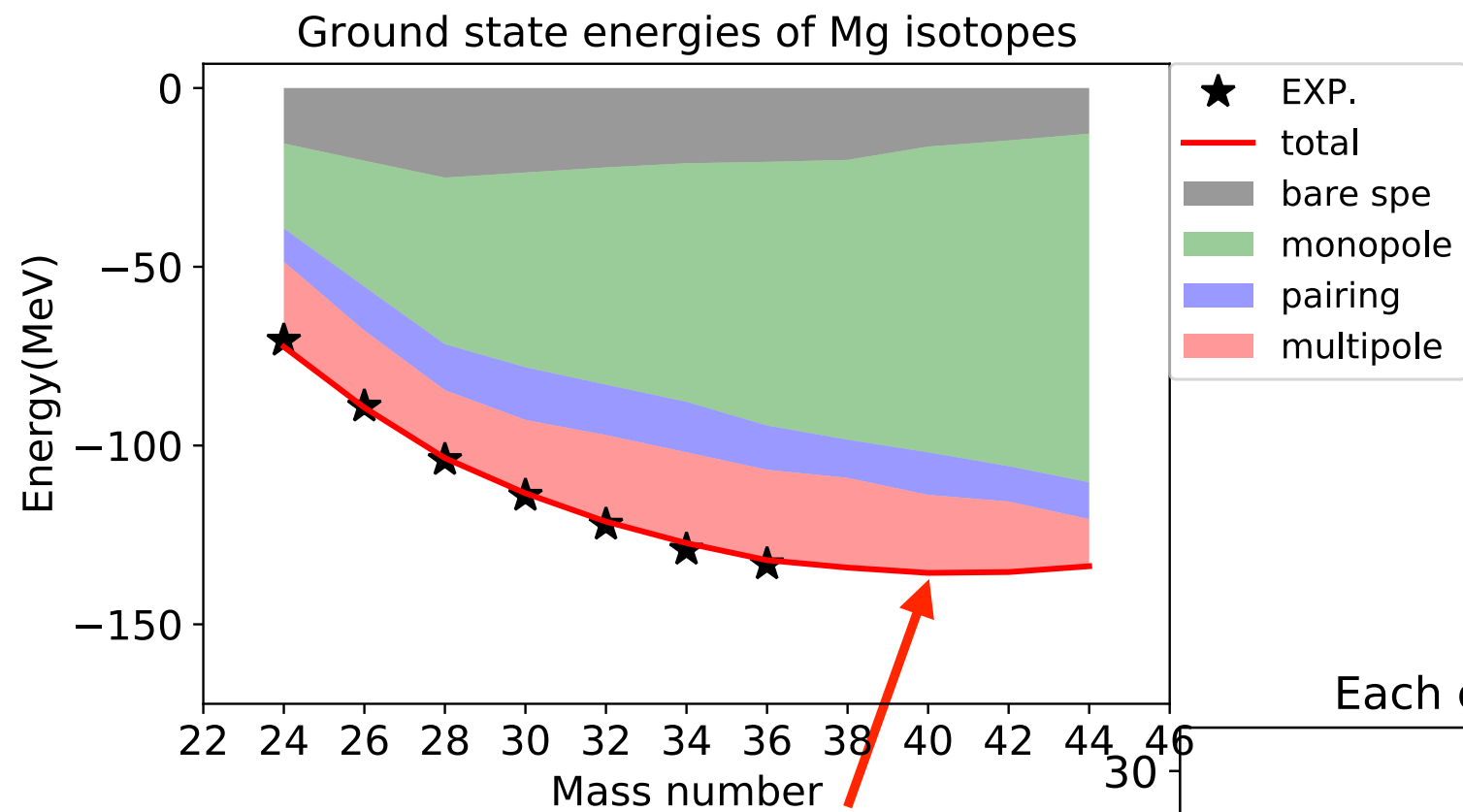
Na isotope (Z=11)



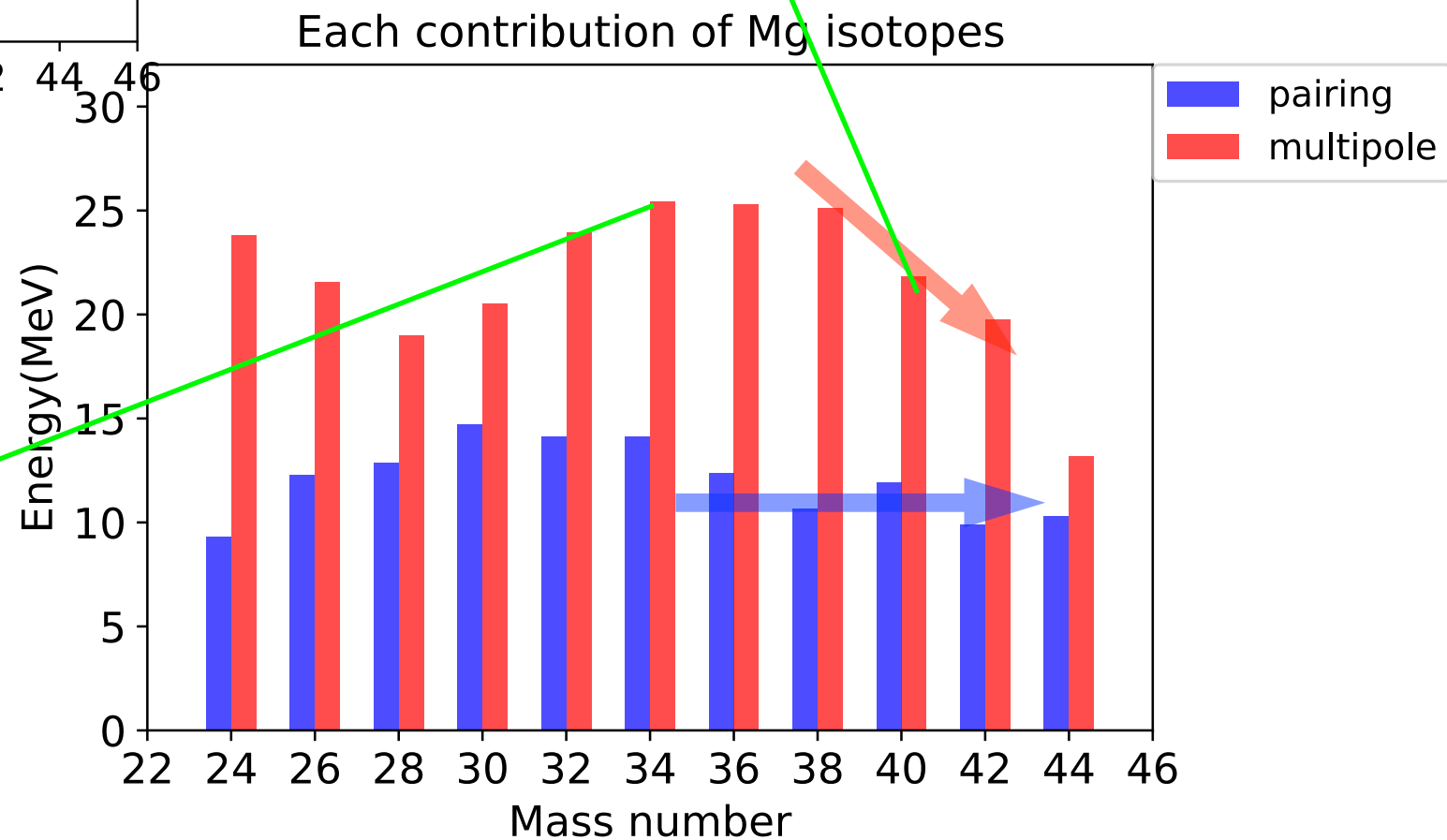
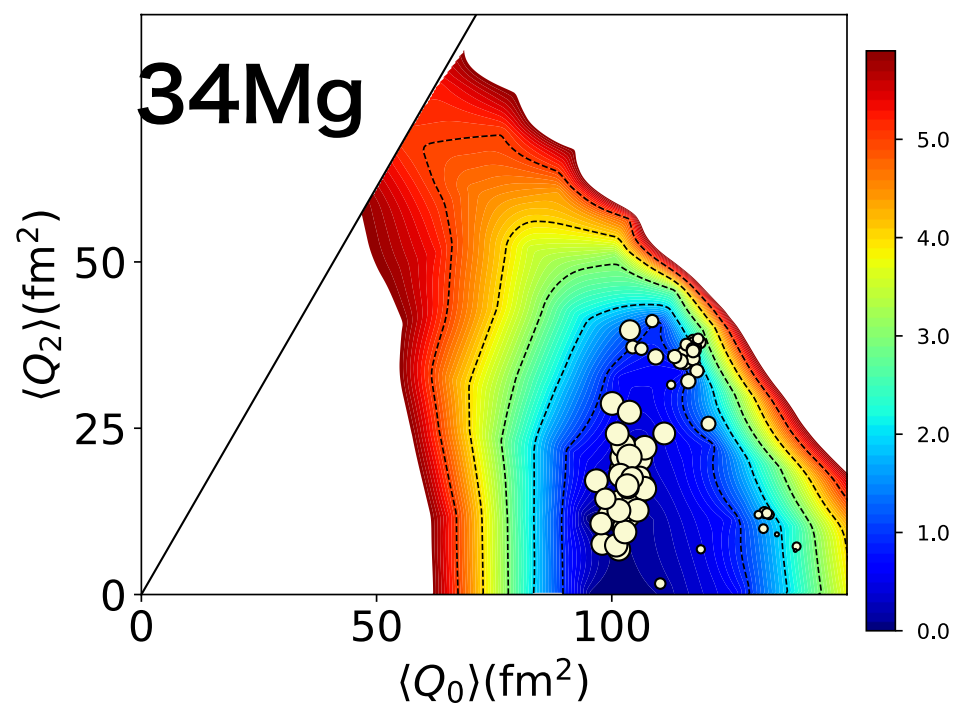
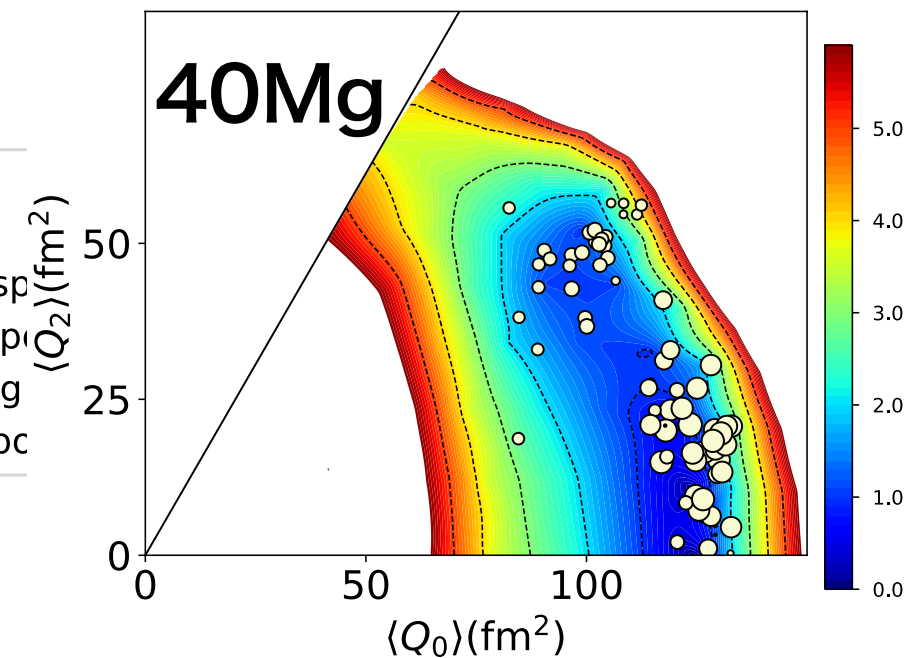
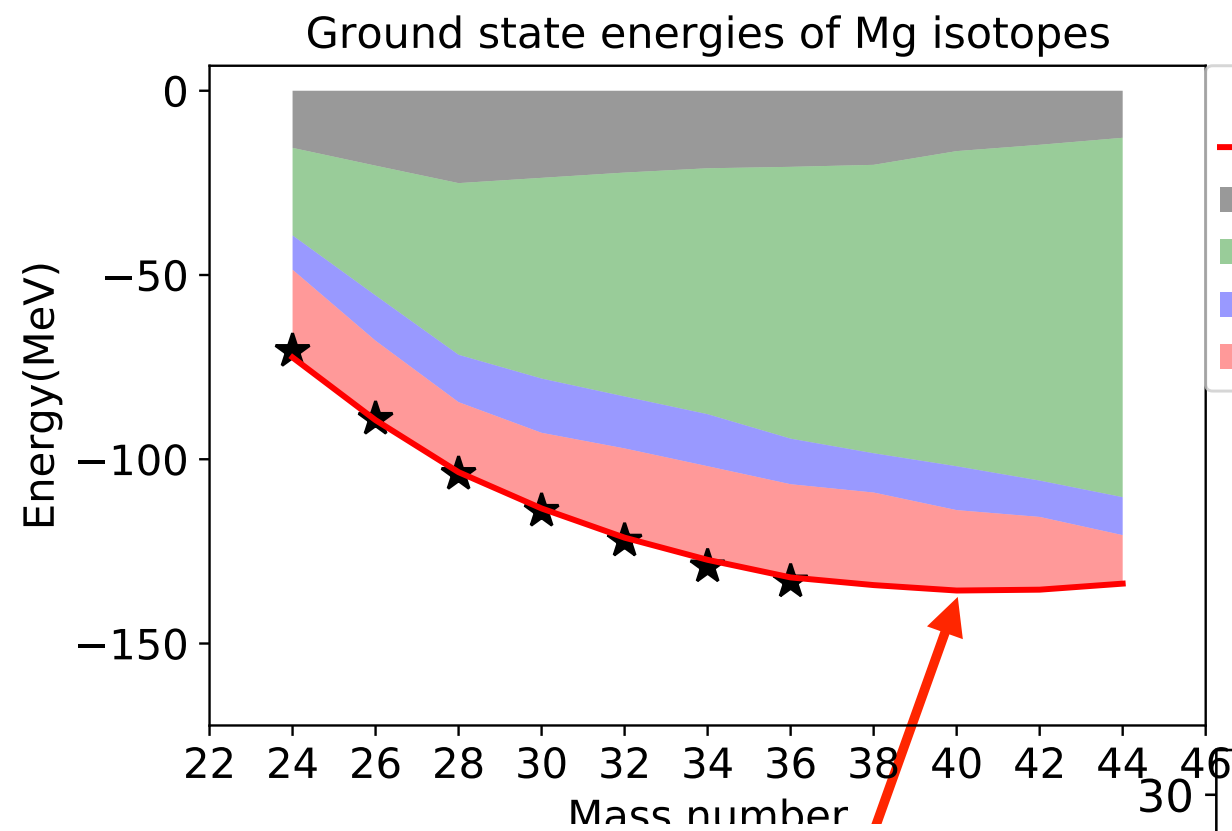
Mg isotope ($Z=12$)



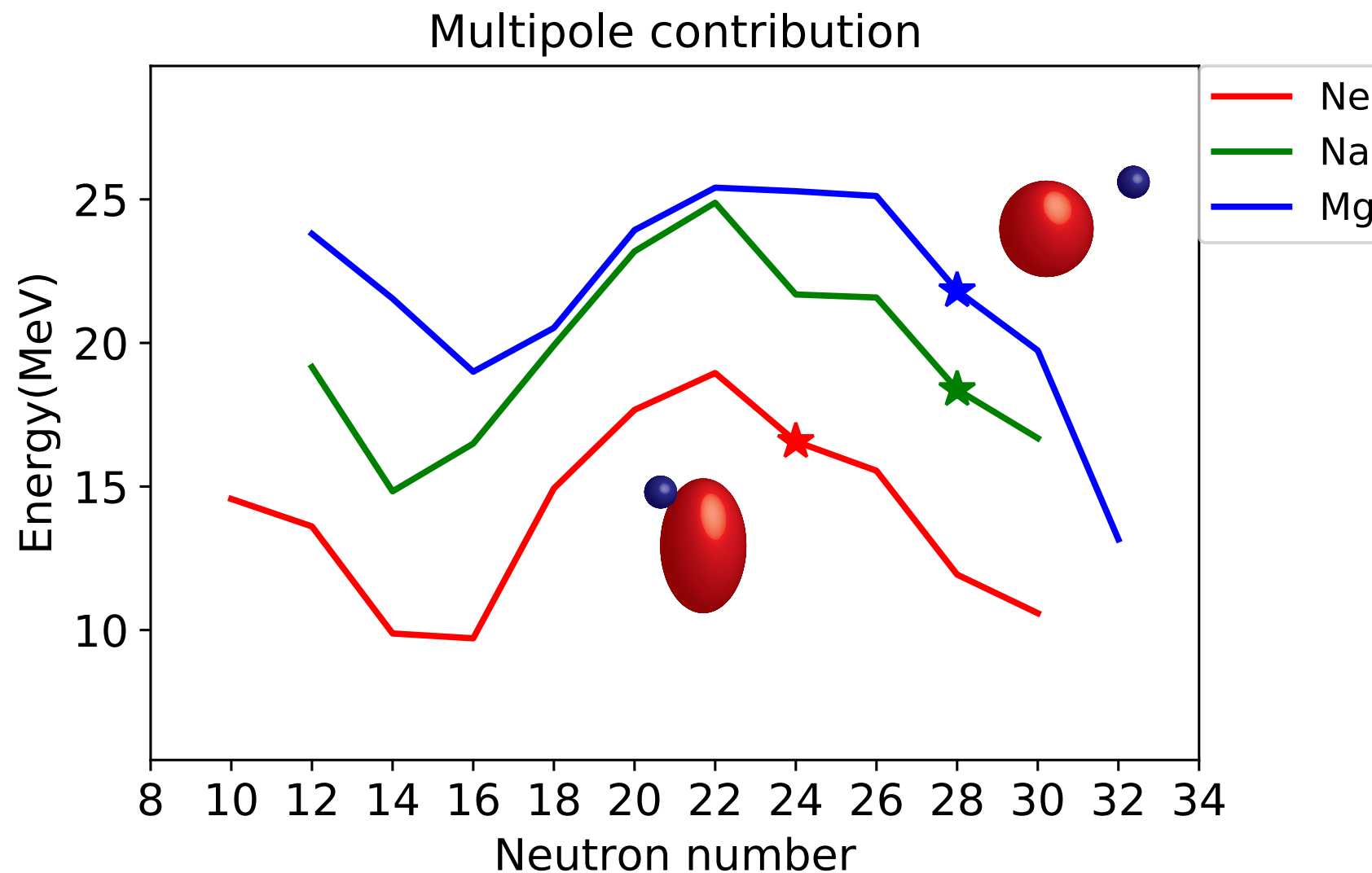
Mg isotope (Z=12)



Mg isotope (Z=12)



Multipole contribution (変形の効果)



A new mechanism of dripline

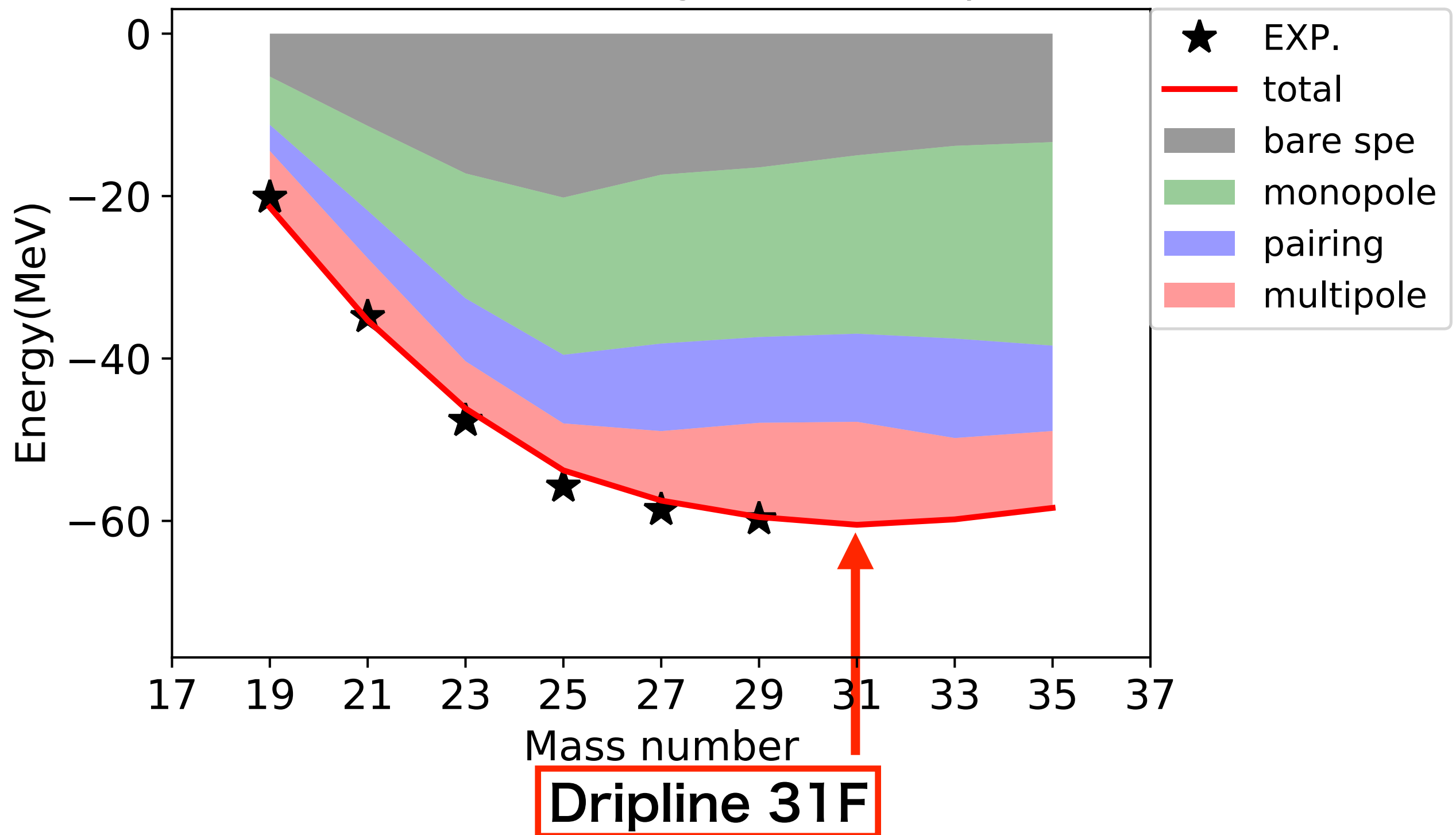
=> **competition** between **EPSE** and **Deformation negative ESPE** but **less deformation energy**

粒子をくっつけると、変形エネルギーが減ってしまい、
束縛できない

F isotope (二つの原理の切り替わり $Z=9$)

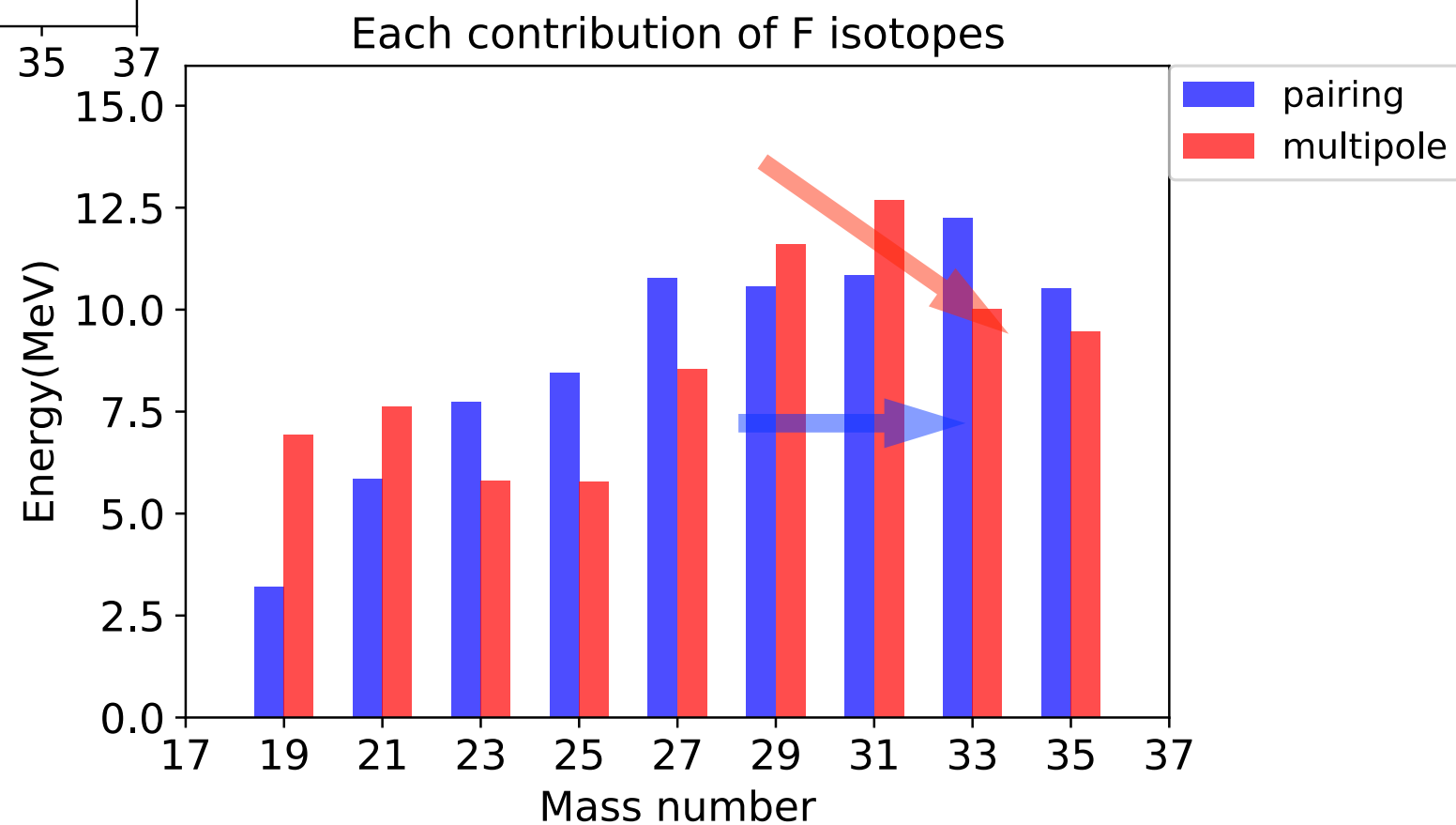
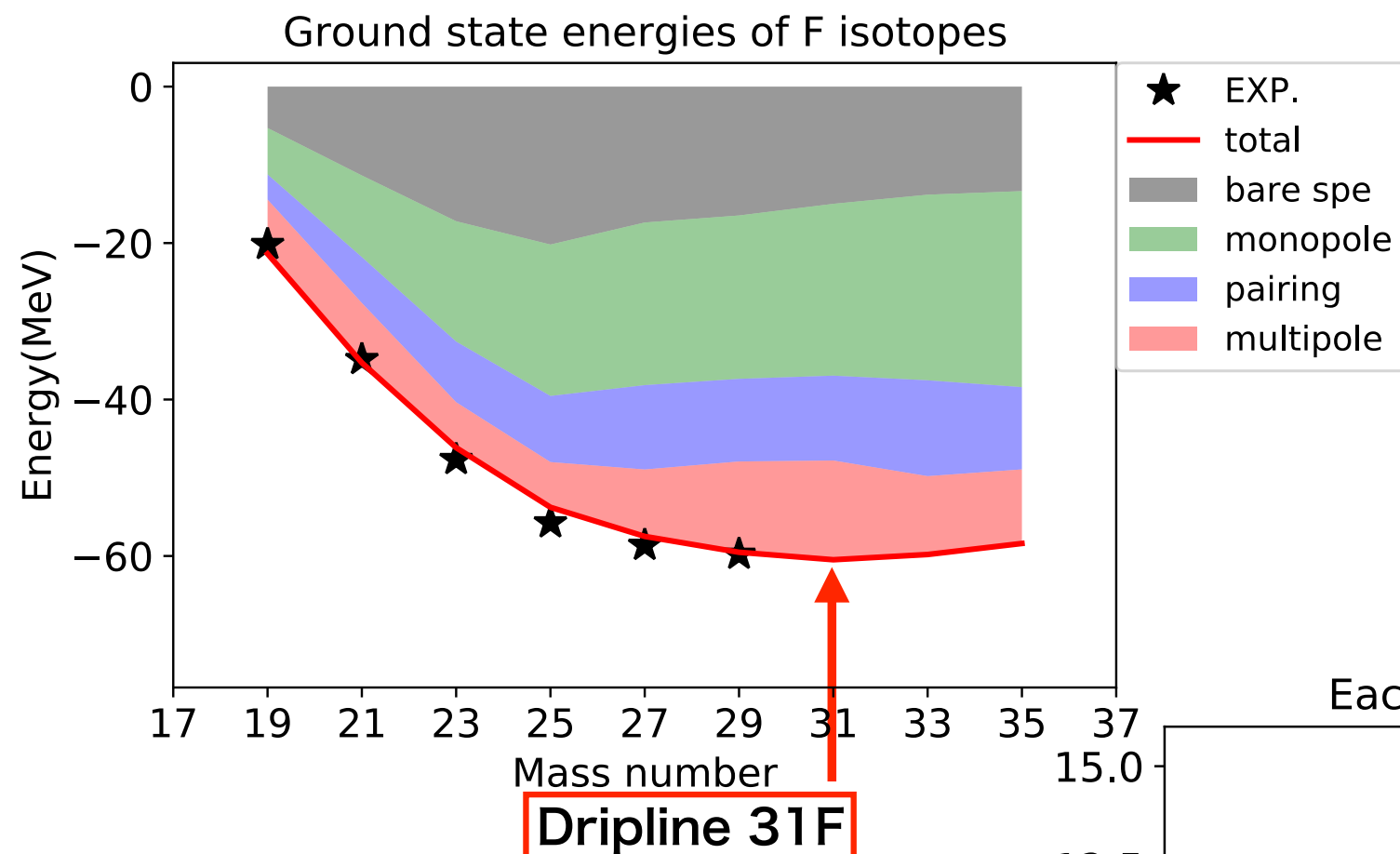
(all the SPEs shifted by 0.9 MeV)

Ground state energies of F isotopes

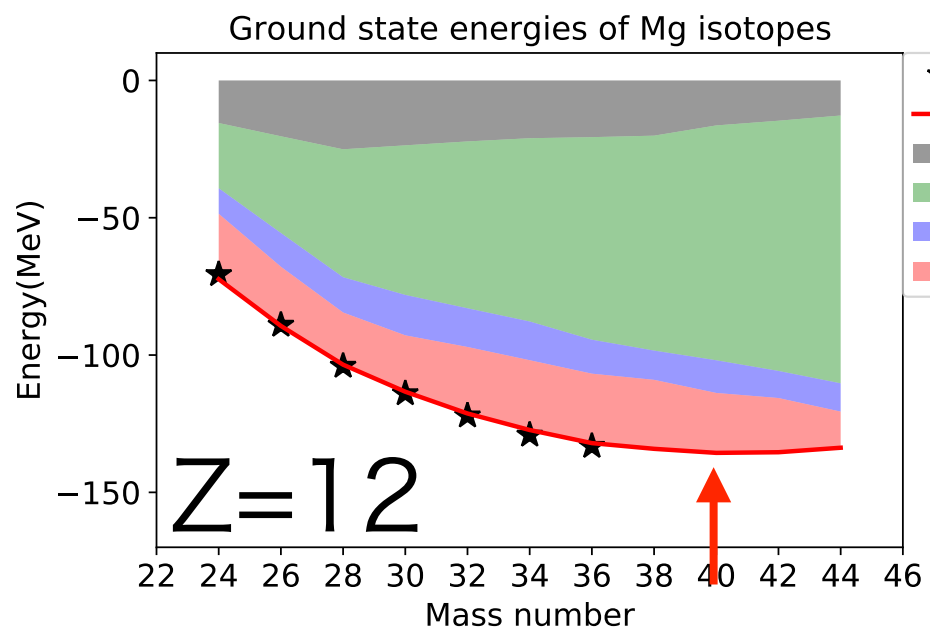
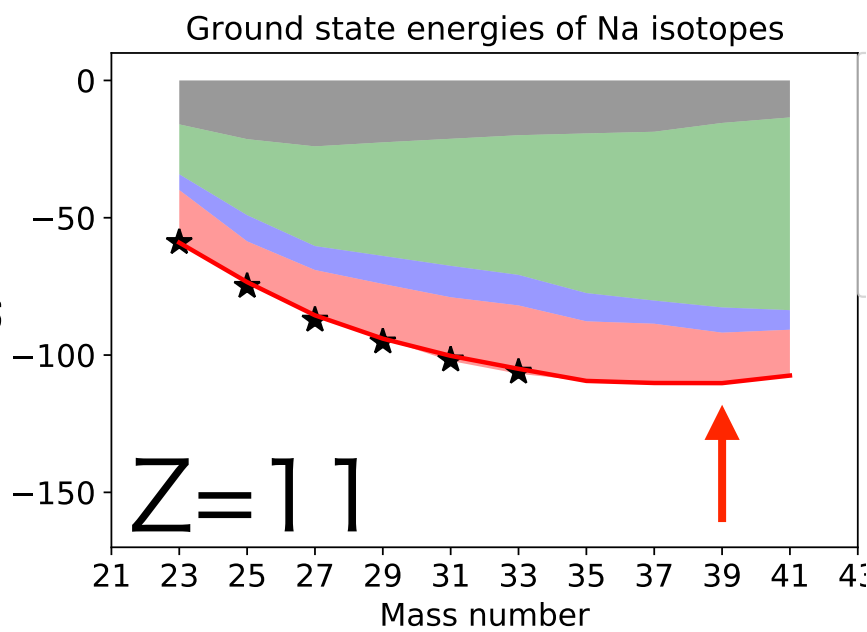
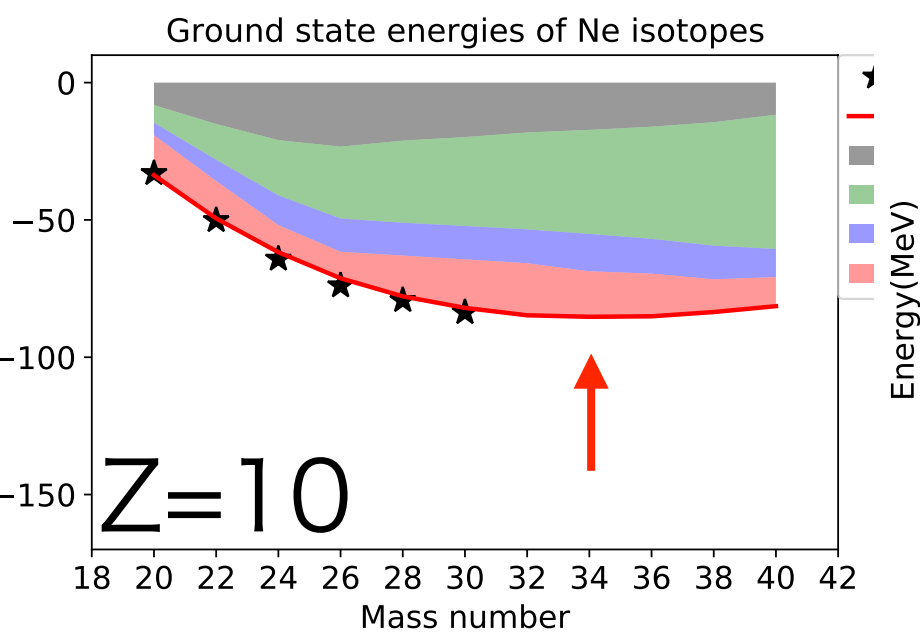
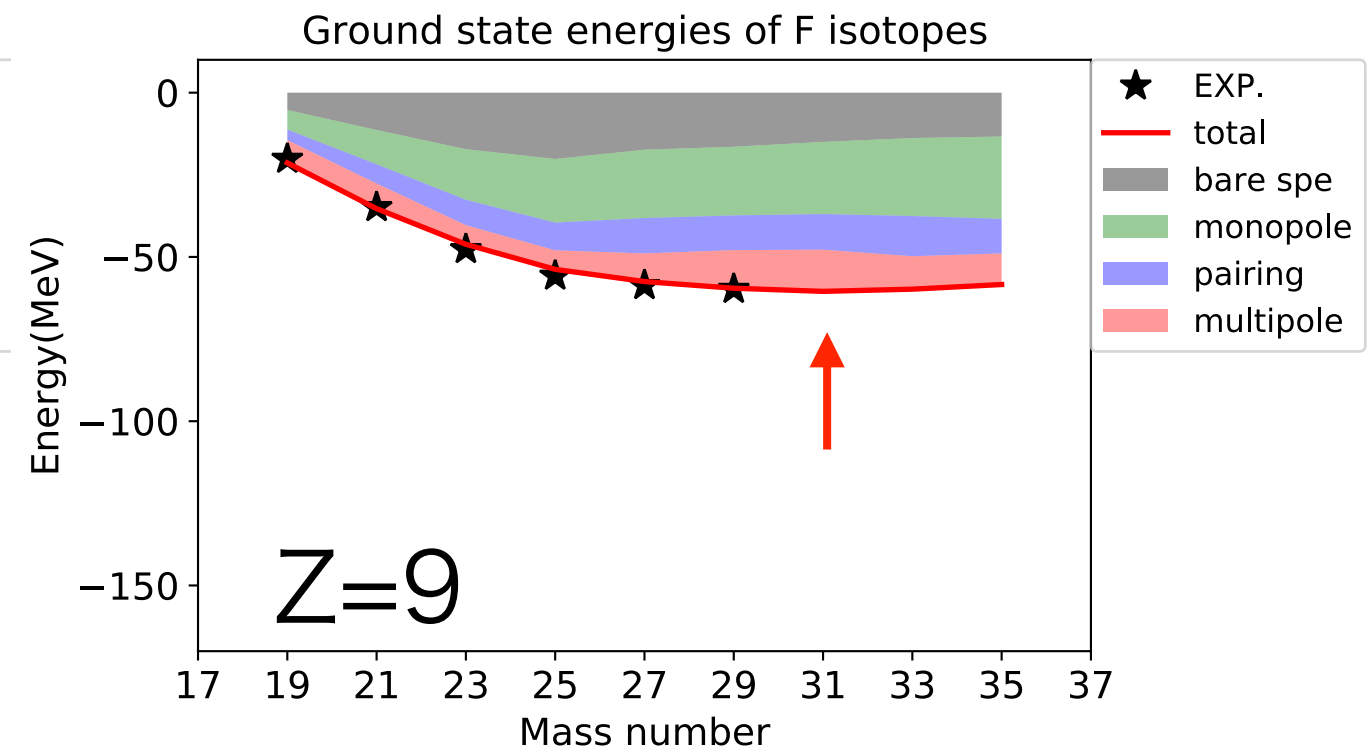
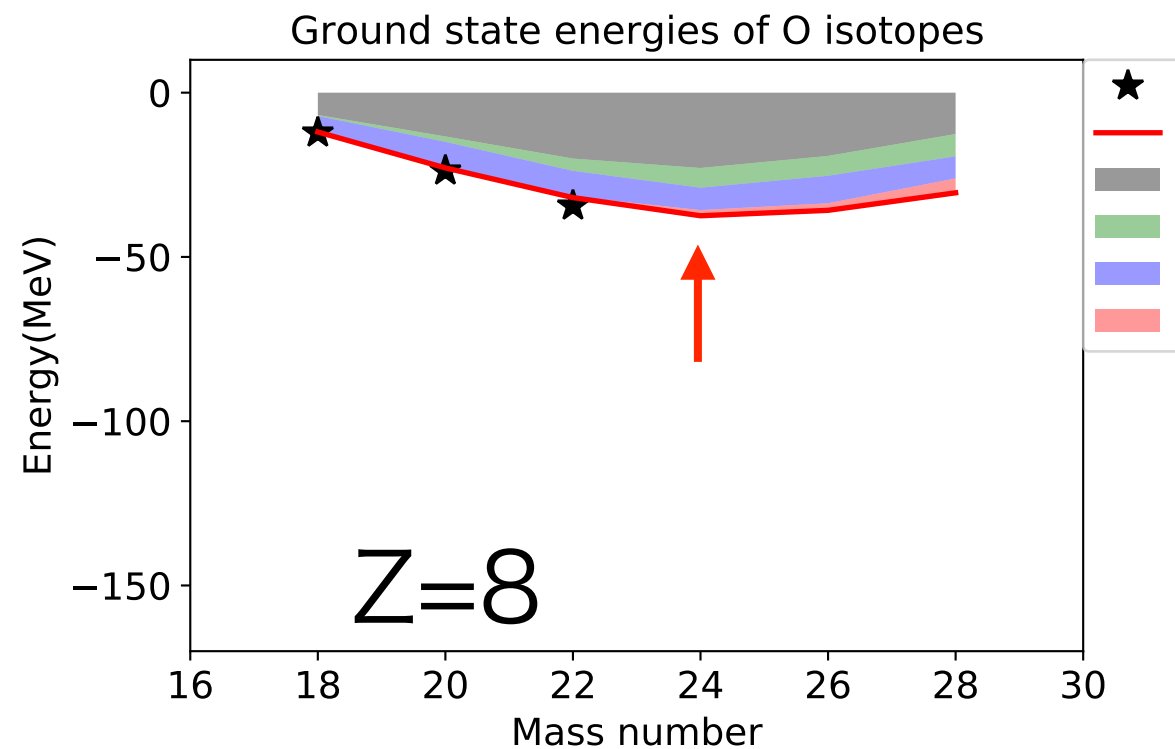


F isotope (二つの原理の切り替わり $Z=9$)

(all the SPEs shifted by 0.9 MeV)



All isotopes



prediction

prediction

- ドリップラインの解析は原子核物理にとって重要な使命
- 「京」を用いた $Z=8-12$ のドリップラインの理論的解析
- 実験との一致
- 新しいドリップラインの原理の発見
- Ne、Na、Mg のドリップラインは**変形エネルギー**が小さくなることによって出現 (0の場合には、変形エネルギーではなく、ナイーブなsingle particle energy によってドリップラインが決定されることと対照的)

- Takaharu Otsuka
- Noritaka Shimizu
- Kazuo Takayanagi
- Morten Hjorth-Jensen
- Toshio Suzuki
- DeukSoon Ahn (and her collaborators)
- Hiroki Nishibata (and his collaborators)
- B. Fernández-Domínguez (and her collaborators)
- Ian Murray (and his collaborators)