HMC on Lefschetz thimbles-- A study of the residual sign problem

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in collaboration with

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based on

arXiv:1309.4371; JHEP10(2013)147

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Plan

I. To define lattice models on Lefschetz thimbles

- Morse function, Gradient flow, Lefschetz thimbles
- * extra phase factors / Tangent spaces >> residual sign problem

2. An algorithm of HMC on Lefschetz thimbles

- a. to parametrize/generate a thimble
- b. to formulate/solve the molecular dynamics on the thimble
- c. to measure observables by reweighting the residual phase

3. Test in the $\lambda \phi^4 \mu$ model

4. Summary & Discussion

Lattice models with complex-valued actions

- QCD with finite chemical potential,
- Chiral gauge theories,
- Chiral Yukawa theories, etc.

physically well-reasoned and -defined, but the state-of-art Monte Carlo methods do not apply straightforwardly

One possible approach based on complexification:

$$\phi_x \in \mathbb{R} \longrightarrow z_x \in \mathbb{C}$$
 $U_{x\mu} = e^{iA_{x\mu}^a T^a} \in SU(3) \longrightarrow e^{iZ_{x\mu}^a T^a} \in SL(3,\mathbb{C})$

complexified Langevin dynamics

G. Aarts, PRL 102(2009) 131601 ($\lambda \phi^4 \mu$) D. Sexty, arXiv:1307.7748 (QCD $_\mu$)

deformation of path-integration contours to Lefschetz thimbles

Analytic Continuation of Chern-Simons Theory

E.Witten, arXiv:1001.2933[hep-lat]

cf. Vanishing Homologies and the n variable saddle point method F. Pham (1983)

New approach of the sign problem in quantum field theories: High density QCD on a Lefschetz thimble (?)

AuroraScience Collaboration, Phys. Rev. D 86, 074506 (2012), arXiv:1205.3996 [hep-lat]

Lattice models on Lefschetz thimbles

$$Z = \int_{\mathcal{C}_{m}} \mathcal{D}[x] \exp\{-S[x]\} = \int_{\mathcal{C}} \mathcal{D}[z] \exp\{-S[z]\}$$
 ($\mathcal{D}[x] = d^{n}x$)

$$x \in \mathcal{C}_{\mathbb{R}} (\subseteq \mathbb{R}^n) \longrightarrow x + iy = z \in \mathbb{C}^n$$

$$S[x] \to S[x+iy] = S[z]$$

how to choose the contour of Path Integration?

Morse theory:

$$h \equiv -\text{Re}\,S[z]$$

$$\frac{d}{dt}z(t) = \frac{\partial \bar{S}[\bar{z}]}{\partial \bar{z}}, \qquad \frac{d}{dt}\bar{z}(t) = \frac{\partial S[z]}{\partial z}, \qquad t \in \mathbb{R}$$

critical points
$$\mathbf{z}_{\sigma}$$
: $\frac{\partial S[z]}{\partial z}\Big|_{z=z_{\sigma}} = 0$

$$\mathcal{C}_{\mathbb{R}} = \sum_{\sigma \in \Sigma} n_{\sigma} \mathcal{J}_{\sigma}, \qquad n_{\sigma} = \langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle$$

$$\frac{d}{dt}h = -\frac{1}{2} \left\{ \frac{\partial S[z]}{\partial z} \cdot \frac{d}{dt} z(t) + \frac{\partial \bar{S}[\bar{z}]}{\partial \bar{z}} \cdot \frac{d}{dt} \bar{z}(t) \right\} = - \left| \frac{\partial S[z]}{\partial z} \right|^2 \le 0$$

$$\frac{d}{dt}\operatorname{Im} S[z] = \frac{1}{2i} \left\{ \frac{\partial S[z]}{\partial z} \cdot \frac{d}{dt} z(t) - \frac{\partial \bar{S}[\bar{z}]}{\partial \bar{z}} \cdot \frac{d}{dt} \bar{z}(t) \right\} = 0$$

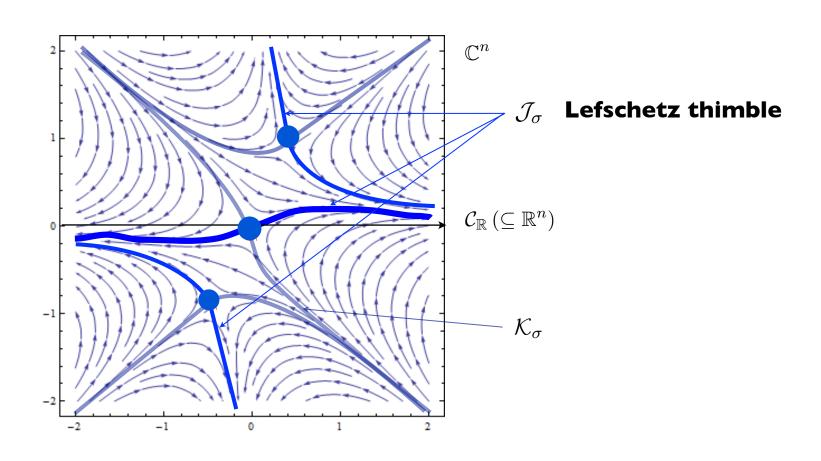
 \mathcal{J}_{σ} the union of all downward flows which trace back to z_{σ} in the limit $t \longrightarrow -\infty$ < **Lefschetz thimble** (n-dim. real mfd.)

 \mathcal{K}_{σ} the union of all upward flows which trace back to z_{σ} in the limit $t \longrightarrow -\infty$

$$\langle \mathcal{J}_{\sigma}, \mathcal{K}_{\tau} \rangle = \delta_{\sigma \tau}$$
 (orthogonarity by intersection numbers)

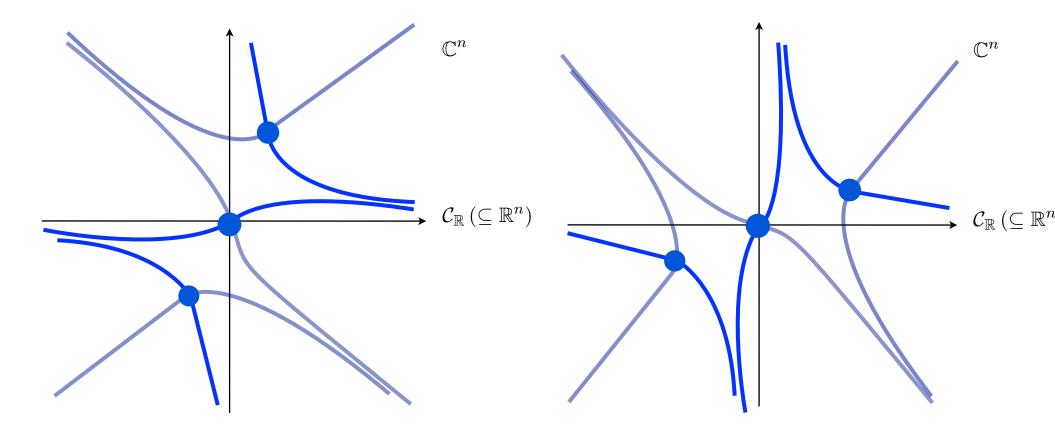
Simple example (n=1)

$$S[z] = \frac{\kappa}{2}z^2 + \frac{\lambda}{4}z^4 \qquad \kappa \in \mathbb{C}, \lambda > 0$$



Simple example (n=1)

$$S[z] = \frac{\kappa}{2}z^2 + \frac{\lambda}{4}z^4 \qquad \kappa \in \mathbb{C}$$



$$Z = \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma}, \qquad n_{\sigma} = \langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle$$
$$Z_{\sigma} = \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\}$$

Observables

$$\langle O[z] \rangle = \frac{1}{Z} \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma} \langle O[z] \rangle_{\mathcal{J}_{\sigma}}$$

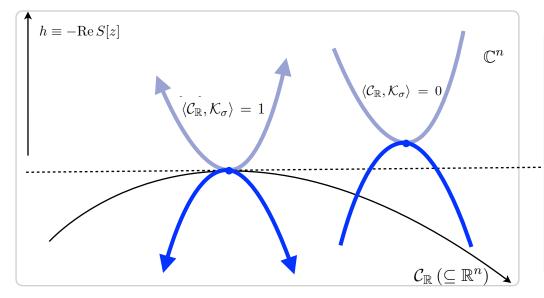
$$\langle O[z] \rangle_{\mathcal{J}_{\sigma}} = \frac{1}{Z_{\sigma}} \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\} O[z]$$

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$$\langle O[z] \rangle_{\mathcal{J}_{\sigma}} = \frac{1}{Z_{\sigma}} \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\} O[z]$$



$$\langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle = 0$$

 $\{z_{\sigma}\} \text{ satisfying } -\text{Re}S[z_{\sigma}] > \max\{-\text{Re}S[x]\} (x \in \mathcal{C}_{\mathbb{R}})$

$$\langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle = 1$$
 $\{z_{\sigma}\}$ in the original cycle $\mathcal{C}_{\mathbb{R}}$
the relative weights proportional to $\exp(-S[z_{\sigma}])$
 $z_{\text{vac}} \in \mathcal{C}_{\mathbb{R}} \quad -\text{Re}\,S[z_{\text{vac}}] = \max\{-\text{Re}S[x]\}\,(x \in \mathcal{C}_{\mathbb{R}})$

$$Z = \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma}, \qquad n_{\sigma} = \langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle$$
$$Z_{\sigma} = \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\}$$

Observables

$$\langle O[z] \rangle = \frac{1}{Z} \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma} \langle O[z] \rangle_{\mathcal{J}_{\sigma}}$$

$$\langle O[z] \rangle_{\mathcal{J}_{\sigma}} = \frac{1}{Z_{\sigma}} \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\} O[z]$$

The functional measure should be specified by **the tangent spaces** of the thimble.

Tangent spaces of Lefschetz thimbles

In the vicinity of critical point z_{σ}

linearized flow equation and its solution:

$$v_i^{\alpha} K_{ij} v_j^{\beta} = \kappa^{\alpha} \delta^{\alpha\beta}$$

$$\kappa^{\alpha} \ge 0 \ (\alpha = 1, \dots, n)$$

$$v_i^{\alpha} (\alpha = 1, \dots, n) \text{ are orthonormal}$$

At a generic point z on \mathcal{J}_{σ}

basis of tangent vectors $\{V_z^{\alpha}\}(\alpha=1,\cdots,n)$

$$\frac{d}{dt}V_{zi}^{\alpha}(t) = \bar{\partial}_i\bar{\partial}_j\bar{S}[\bar{z}] \ \bar{V}_{zj}^{\alpha}(t) \qquad (\alpha = 1, \cdots, n)$$

$$\bar{V}_{zi}^{\alpha}V_{zi}^{\beta} - \bar{V}_{zi}^{\beta}V_{zi}^{\alpha} = 0 \qquad (\alpha, \beta = 1, \dots, n)$$

 $V_{z}^{\alpha} = U_{z}^{\beta} E^{\beta \alpha} \quad \{U_{z}^{\alpha}\} \text{ is an orthonormal basis}$ E is a real upper triangle matrix

Residual phase factor

$$\delta z = U_z^{\alpha} \delta \xi^{\alpha} \qquad |\delta z|^2 = \delta \xi^2$$

$$d^n z |_{\mathcal{J}_{\sigma}} = d^n \delta \xi \det U_z$$

$$e^{i\phi_z} = \det U_z = \frac{\det V_z}{|\det V_z|}$$

$$Z = \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma}, \qquad n_{\sigma} = \langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle$$
$$Z_{\sigma} = \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\}$$

Observables

$$\langle O[z] \rangle = \frac{1}{Z} \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma} \langle O[z] \rangle_{\mathcal{J}_{\sigma}}$$

$$\langle O[z] \rangle_{\mathcal{J}_{\sigma}} = \frac{1}{Z_{\sigma}} \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\} O[z]$$

The functional measure should be specified by **the tangent spaces** of the thimble.

The functional measure may give rise to an extra phase factor!
>> residual sign problem

$$Z = \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma}, \qquad n_{\sigma} = \langle \mathcal{C}_{\mathbb{R}}, \mathcal{K}_{\sigma} \rangle$$

$$Z_{\sigma} = \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\}$$

Observables

$$\langle O[z] \rangle = \frac{1}{Z} \sum_{\sigma \in \Sigma} n_{\sigma} \exp\{-S[z_{\sigma}]\} Z_{\sigma} \langle O[z] \rangle_{\mathcal{J}_{\sigma}}$$

$$\langle O[z] \rangle_{\mathcal{J}_{\sigma}} = \frac{1}{Z_{\sigma}} \int_{\mathcal{J}_{\sigma}} \mathcal{D}[z] \exp\{-\operatorname{Re}(S[z] - S[z_{\sigma}])\} O[z]$$

This part is not straightforward to compute, in general

$$Z_{\sigma} = 1/\sqrt{\det K}$$

$$K_{ij} \equiv \left. \partial_i \partial_j S[z] \right|_{z=z_\sigma}$$

in the saddle point approximation

a possible approximation : take a single thimble \mathcal{J}_{vac}

$$\langle O[z] \rangle = \langle O[z] \rangle_{\mathcal{J}_{\text{vac}}}$$

(AuroraScience Collaboration)

This part may be evaluated by MC, but with the residual phase factor reweighted

Algorithm of HMC on Lefschetz thimbles

the saddle-point structures!

- a) To parametrize / generate a thimble unique parametrization in the asymptotic region close to the critical point / solve the flow eqs. by 4th-order RK
- b) To formulate / solve the molecular dynamics dynamical system constrained to the thimble by the forces in the normal directions / use 2nd-order constraint-preserving symmetric integrator
- c) To measure observables reweighting the residual sign factors

a) To parametrize/generate a thimble

Unique parametrization of z (field conf.) on the thimbles

Asymptotic solutions of gradient flow equations

$$z(t) \simeq z_{\sigma} + v^{\alpha} \exp(\kappa^{\alpha} t) e^{\alpha}; \qquad e^{\alpha} e^{\alpha} = n$$

 $V_{z}^{\alpha}(t) \simeq v^{\alpha} \exp(\kappa^{\alpha} t),$

the **direction** of the flow: e^{α} $(\alpha = 1, \dots, n; ||e||^2 = n)$

the **time** of the flow : $t' = t - t_0$

$$z[e, t'] : (e^{\alpha}, t') \to z \in \mathcal{J}_{\sigma}$$

$$z[e, t'] = z(t)|_{t=t'+t_0}$$

variation of z

$$\delta z[e, t'] = V_z^{\alpha}[e, t'] \left(\delta e^{\alpha} + \kappa^{\alpha} e^{\alpha} \delta t'\right)$$

$$\delta \dot{z}_i(t) = \bar{\partial}_i \bar{\partial}_j \bar{S}[\bar{z}] \, \delta z_j(t)$$

$$\delta z(t) = V_z^{\alpha}(t) \, \delta c^{\alpha}$$

$$\delta z(t) = v^{\alpha} \exp(\kappa^{\alpha} t) (\delta e^{\alpha} + e^{\alpha} \kappa^{\alpha} \delta t)$$

$$= V_z^{\alpha}(t) (\delta e^{\alpha} + \kappa^{\alpha} e^{\alpha} \delta t) \qquad (t \ll 0)$$

a) To parametrize/generate a thimble

Solve the flow equations by 4th order Runge-Kutta for a given (e^{α}, t') and $t_0 (\ll 0)$ with the initial conditions,

$$z_i(t_0) = z_{\sigma i} + v_i^{\alpha} \exp(\kappa^{\alpha} t_0) e^{\alpha}$$
$$V_{zi}^{\alpha}(t_0) = v_i^{\alpha} \exp(\kappa^{\alpha} t_0).$$

$$h = \tau/n_{lefs}$$

$$f_{1} = \bar{\partial}\bar{S}[z[t]]$$

$$f_{2} = \bar{\partial}\bar{S}[z[t] + (h/2)f_{1}]$$

$$f_{3} = \bar{\partial}\bar{S}[z[t] + (h/2)f_{2}]$$

$$f_{4} = \bar{\partial}\bar{S}[z[t] + hf_{3}]$$

$$z(t+h) = z(t) + \frac{h}{6}(f_{1} + 2f_{2} + 2f_{3} + f_{4})$$

$$F_{1}^{\alpha} = \bar{V}^{\alpha}(t)\bar{\partial}\bar{\partial}\bar{S}[z[t]]$$

$$F_{2}^{\alpha} = \{\bar{V}^{\alpha}(t) + (h/2)\bar{F}_{1}^{\alpha}\}\bar{\partial}\bar{\partial}\bar{S}[z[t] + (h/2)f_{1}]$$

$$F_{3}^{\alpha} = \{\bar{V}^{\alpha}(t) + (h/2)\bar{F}_{2}^{\alpha}\}\bar{\partial}\bar{\partial}\bar{S}[z[t] + (h/2)f_{2}]$$

$$F_{4}^{\alpha} = \{\bar{V}^{\alpha}(t) + h\bar{F}_{3}^{\alpha}\}\bar{\partial}\bar{\partial}\bar{S}[z[t] + hf_{3}]$$

$$V^{\alpha}(t+h) = V^{\alpha}(t) + \frac{h}{6}(F_{1}^{\alpha} + 2F_{2}^{\alpha} + 2F_{3}^{\alpha} + F_{4}^{\alpha})$$

To check the solutions

$$\bar{\partial}_i \bar{S} \left[\bar{z}[e, t'] \right] - V_{zi}^{\alpha}[e, t'] \kappa^{\alpha} e^{\alpha} = 0.$$

b) To formulate/solve Molecular Dynamics on the thimble

Constrained dynamical system

Equations of motion:

$$\dot{z}_i = w_i,$$

$$\dot{w}_i = -\bar{\partial}_i \bar{S}[\bar{z}] - iV_{zi}^{\alpha} \lambda^{\alpha} \qquad \lambda^{\alpha} \in \mathbb{R} \ (\alpha = 1, \dots, n)$$

Constraints:

$$z_i = z_i[e, t']$$
 $w_i = V_{zi}^{\alpha}[e, t'] w^{\alpha}, \quad w^{\alpha} \in \mathbb{R}$

the set of normal vectors

$$\{iU_z^{\alpha}\}\ {\rm or}\ \{iV_z^{\alpha}\}(\alpha=1,\cdots,n)$$

$$\bar{V}_{zi}^{\alpha}V_{zi}^{\beta} - \bar{V}_{zi}^{\beta}V_{zi}^{\alpha} = 0 \qquad (\alpha, \beta = 1, \dots, n)$$

$$\operatorname{Re}\left\{(-i)\bar{V}_{zi}^{\alpha}\,V_{zi}^{\beta}\right\} = 0$$

A conserved Hamiltonian:

$$H = \frac{1}{2}\bar{w}_i w_i + \frac{1}{2} \left\{ S[z] + \bar{S}[\bar{z}] \right\}$$

$$\dot{H} = \frac{1}{2} \{ \dot{\bar{w}}_i w_i + \bar{w}_i \dot{w}_i \} + \frac{1}{2} \{ \partial_i S[z] \dot{z}_i + \bar{\partial}_i \bar{S}[\bar{z}] \dot{\bar{z}}_i \}$$

$$= \frac{1}{2} \{ (+i\bar{V}_{zi}^{\alpha} \lambda^{\alpha}) w_i + \bar{w}_i (-iV_{zi}^{\alpha} \lambda^{\alpha}) \}$$

$$= \frac{i}{2} \lambda^{\alpha} w^{\beta} \{ \bar{V}_{zi}^{\alpha} V_{zi}^{\beta} - \bar{V}_{zi}^{\beta} V_{zi}^{\alpha} \} = 0.$$

b) To formulate/solve Molecular Dynamics on the thimble

Second-order constraint-preserving symmetric integrator

$$z^{n} = z[e^{(n)}, t'^{(n)}],$$

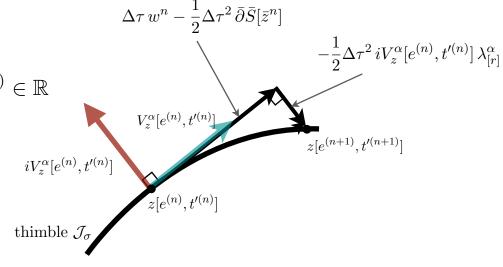
 $w^{n} = V_{z}^{\alpha}[e^{(n)}, t'^{(n)}] w^{\alpha(n)}, \quad w^{\alpha(n)} \in \mathbb{R},$

$$\begin{split} w^{n+1/2} &= w^n & -\frac{1}{2} \Delta \tau \, \bar{\partial} \bar{S}[\bar{z}^n] - \frac{1}{2} \Delta \tau \, i V_z^{\alpha}[e^{(n)}, t'^{(n)}] \, \lambda_{[r]}^{\alpha}, \\ z^{n+1} &= z^n & + \Delta \tau \, w^{n+1/2}, \\ w^{n+1} &= w^{n+1/2} - \frac{1}{2} \Delta \tau \, \bar{\partial} \bar{S}[\bar{z}^{n+1}] - \frac{1}{2} \Delta \tau \, i V_z^{\alpha}[e^{(n+1)}, t'^{(n+1)}] \, \lambda_{[v]}^{\alpha} \end{split}$$

$$\lambda^{\alpha}_{[r]}$$
 and $\lambda^{\alpha}_{[v]}$ are fixed by

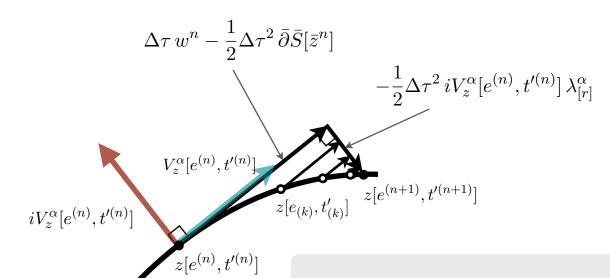
$$\begin{split} z^{n+1} &= z[e^{(n+1)}, t'^{(n+1)}], \\ w^{n+1} &= V_z^{\alpha}[e^{(n+1)}, t'^{(n+1)}] \, w^{\alpha(n+1)}, \quad w^{\alpha(n+1)} \in \mathbb{R} \end{split}$$

To solve the constraints ...



Solving the constraints

thimble \mathcal{J}_{σ}



the sequences $(e_{(k)}^{\alpha}, t_{(k)}')$ $(k = 0, 1, \cdots)$ with $(e_{(0)}^{\alpha}, t_{(0)}') = (e^{\alpha(n)}, t'^{(n)})$

$$\Delta e_{(k)}^{\alpha} = e_{(k+1)}^{\alpha} - e_{(k)}^{\alpha}, \qquad \sum_{\alpha=1} \Delta e_{(k)}^{\alpha} e^{\alpha(n)} = 0,$$

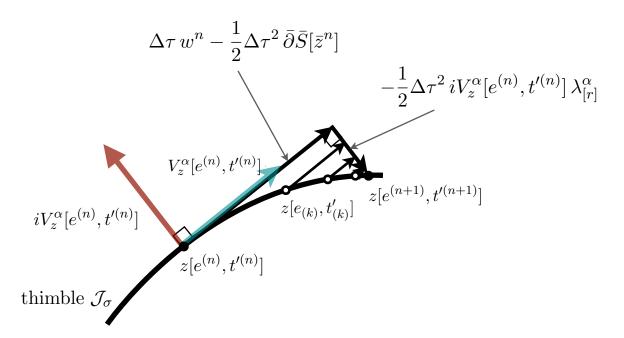
$$\Delta t_{(k)}' = t_{(k+1)}' - t_{(k)}',$$

$$\Delta e^{\alpha}{}_{(k)} + e^{\alpha(n)} \kappa^{\alpha} \Delta t'_{(k)} = \text{Re} \left[\{ V_z^{-1} [e^{(n)}, t'^{(n)}] \}_i^{\alpha} \times \left(z_i [e^{(n)}, t'^{(n)}] + \Delta \tau \, w_i^n - \frac{1}{2} \Delta \tau^2 \, \bar{\partial}_i \bar{S}[\bar{z}^n] - z_i [e_{(k)}, t'_{(k)}] \right) \right]$$

$$\frac{1}{2}\Delta \tau^2 \lambda_{[r](k)}^{\alpha} = \operatorname{Im}\left[\left\{V_z^{-1}[e^{(n)}, t'^{(n)}]\right\}_i^{\alpha} \left(z_i[e^{(n)}, t'^{(n)}] - z_i[e_{(k)}, t'_{(k)}]\right)\right]$$

$$\left\| V_z^{\alpha}[e^{(n)}, t'^{(n)}] \left(\Delta e^{\alpha}_{(k)} + e^{\alpha(n)} \kappa^{\alpha} \Delta t'_{(k)} \right) \right\|^2 \le n \epsilon'^2$$

Solving the constraints



$$\frac{1}{2}\Delta\tau \,\lambda_{[v]}^{\alpha} = \operatorname{Im}\left[\left\{V_{z}^{-1}[e^{(n+1)}, t'^{(n+1)}]\right\}_{i}^{\alpha}\left(w_{i}^{n+1/2} - \frac{1}{2}\Delta\tau \,\bar{\partial}_{i}\bar{S}[\bar{z}^{n+1}]\right)\right]$$

a HMC update

A hybrid Monte Carlo update then consists of the following steps for a given trajectory length τ_{traj} and a number of steps n_{step} :

1. Set the initial field configuration z_i :

$$\{e^{\alpha(0)}, t'^{(0)}\} = \{e^{\alpha}, t'\}, \qquad z^0 = z[e, t'].$$

2. Refresh the momenta w_i by generating n pairs of unit gaussian random numbers (ξ_i, η_i) , setting tentatively $w_i = \xi_i + i\eta_i$, and chopping the non-tangential parts:

$$w^{0} = V_{z}^{\alpha} \operatorname{Re}[\{V_{z}^{-1}\}_{j}^{\alpha}(\xi_{j} + i\eta_{j})] = U_{z}^{\alpha} \operatorname{Re}[\{U_{z}^{-1}\}_{j}^{\alpha}(\xi_{j} + i\eta_{j})].$$

- 3. Repeat n_{step} times of the second order symmetric integration the step size $\Delta \tau = \tau_{\text{traj}}/n_{\text{step}}$.
- 4. Accept or reject by $\Delta H = H[w^{n_{\text{step}}}, z^{n_{\text{step}}}] H[w^0, z^0].$

As for the initialization procedure, one may generate unit gaussian random numbers $\eta^{\alpha}(\alpha=1,\cdots,n)$, set

$$e^{\alpha} = \eta^{\alpha} \sqrt{\frac{n}{\sum_{\beta=1}^{n} \eta^{\beta} \eta^{\beta}}}, \qquad t' = -t_0,$$

and then prepare z[e,t'], $\{V_z^{\alpha}[e,t']\},$ and the inverse matrix $V_z^{-1}[e,t'].$

c) To measure observables by reweighting the residual sign factors

$$\langle o[z] \rangle'_{\mathcal{J}_{\sigma}} = \frac{1}{N_{\text{conf}}} \sum_{k=1}^{N_{\text{conf}}} o[z^{(k)}]$$

$$\langle O[z] \rangle_{\mathcal{J}_{\sigma}} = \frac{\langle e^{i\phi_z} O[z] \rangle_{\mathcal{J}_{\sigma}}'}{\langle e^{i\phi_z} \rangle_{\mathcal{J}_{\sigma}}'} \qquad e^{i\phi_z} = \det U_z = \frac{\det V_z}{|\det V_z|}$$

 $\{\langle \mathrm{e}^{i\phi_z}\rangle'_{\mathcal{J}_\sigma}\}(\sigma\in\Sigma)$ should not be vanishingly small

A possible sign problem !!

Need a careful and systematic study !!

Test in the $\lambda \phi^4 \mu$ model

$$S = \sum_{x \in \mathbb{L}^4} \left\{ \left(\varphi^{\dagger}(x+\hat{0}) e^{+\mu} - \varphi^{\dagger}(x) \right) \left(e^{-\mu} \varphi(x+\hat{0}) - \varphi(x) \right) \right.$$

$$\left. + \sum_{k=1}^3 |\varphi(x+\hat{k}) - \varphi(x)|^2 + \frac{\kappa}{2} \varphi^{\dagger}(x) \varphi(x) + \frac{\lambda}{4} \left(\varphi^{\dagger}(x) \varphi(x) \right)^2 \right\}$$

$$= \sum_{x \in \mathbb{L}^4} \left\{ -\phi_a(x) \phi_b(x+\hat{0}) \left[\delta_{ab} \cosh(\mu) - i \epsilon_{ab} \sinh(\mu) \right] \right.$$

$$\left. - \sum_{k=1}^3 \phi_a(x) \phi_a(x+\hat{k}) + \frac{(8+\kappa)}{2} \phi_a(x) \phi_a(x) + \frac{\lambda}{4} \left(\phi_a(x) \phi_a(x) \right)^2 \right\}$$

$$\varphi(x) = (\phi_1(x) + i\phi_2(x))/\sqrt{2}$$
$$\phi_a(x) \in \mathbb{R} \ (a = 1, 2)$$

$$\phi_a(x) \to z_a(x) \in \mathbb{C} \ (a=1,2)$$

$$S[z] = \sum_{x \in \mathbb{L}^4} \left\{ + \frac{1}{2} z_a(x) z_a(x) + \frac{\lambda_0}{4} \left(z_a(x) z_a(x) \right)^2 - K_0 \sum_{k=1}^3 z_a(x) z_a(x + \hat{k}) - K_0 z_a(x) z_b(x + \hat{0}) \left[\delta_{ab} \cosh(\mu) - i \epsilon_{ab} \sinh(\mu) \right] \right\}.$$

where
$$K_0 = \frac{1}{(2D+\kappa)}$$
, $\lambda_0 = K_0^2 \lambda$

K=1.0, λ =1.0, μ =0.0~1.8

cf. G.Aarts (Complex Langevin simulation)

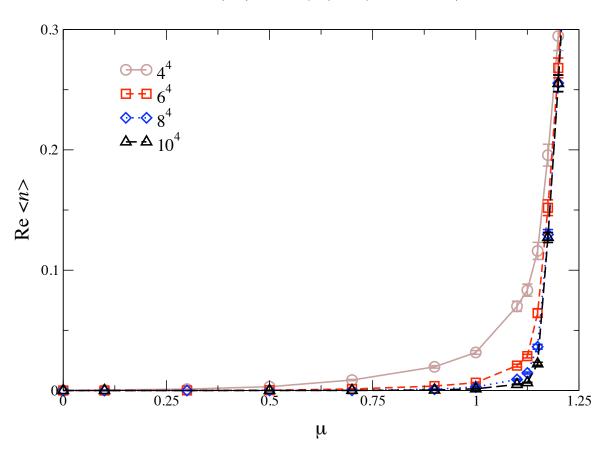
Can stochastic quantization evade the sign problem?
-- the relativistic Bose gas at finite chemical potential
G.Aarts, PRL 102:131601, 2009 arXiv:0810.2089

cf. G.Aarts (Complex Langevin simulation)

RELATIVISTIC BOSE GAS

COMPLEX LANGEVIN

density
$$\langle n \rangle = (1/\Omega) \partial \ln Z / \partial \mu$$



second order phase transition in thermodynamic limit

cf. G.Aarts PRL 102(2009) 131601 (Complex Langevin simulation)

SILVER BLAZE AND THE SIGN PROBLEM

COMPLEX VS PHASE QUENCHED

density $\Theta - \Theta 4^4$ **△ △** 10⁴ $\triangle \triangle 10^4$ 0.2 <*n*> pq Re <*n*> 0.1 0.1 0.75 1.25 μ μ phase quenched

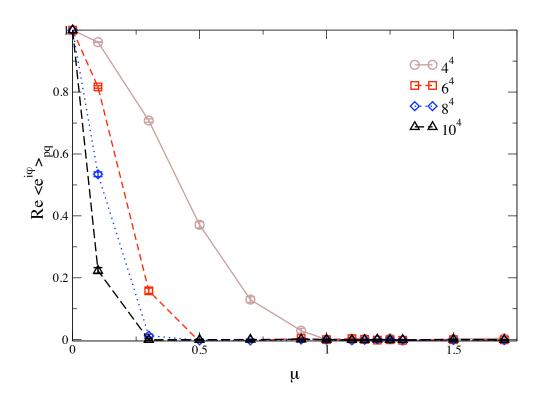
complex

phase $e^{i\varphi} = e^{-S}/|e^{-S}|$ does precisely what is expected

cf. G.Aarts (Complex Langevin simulation)

HOW SEVERE IS THE SIGN PROBLEM?

AVERAGE PHASE FACTOR



average phase factor $\langle e^{i\varphi} \rangle_{pq}$

Test in the $\lambda \phi^4 \mu$ model (cont'd)

critical points with constant field $z_a(x)=z_a$

$$\left. \frac{\partial S[z]}{\partial z_a(x)} \right|_{z_a(x)=z_a} = (1 - 6K_0 - 2K_0 \cosh(\mu)) z_a + \lambda_0 (z_1^2 + z_2^2) z_a = 0 \quad (a = 1, 2).$$

critical value of
$$\mu$$
 (classical) $\tilde{\mu}_c = \ln \left[\left(\frac{1-6K_0}{2K_0} \right) + \sqrt{\left(\frac{1-6K_0}{2K_0} \right)^2 - 1} \right]$

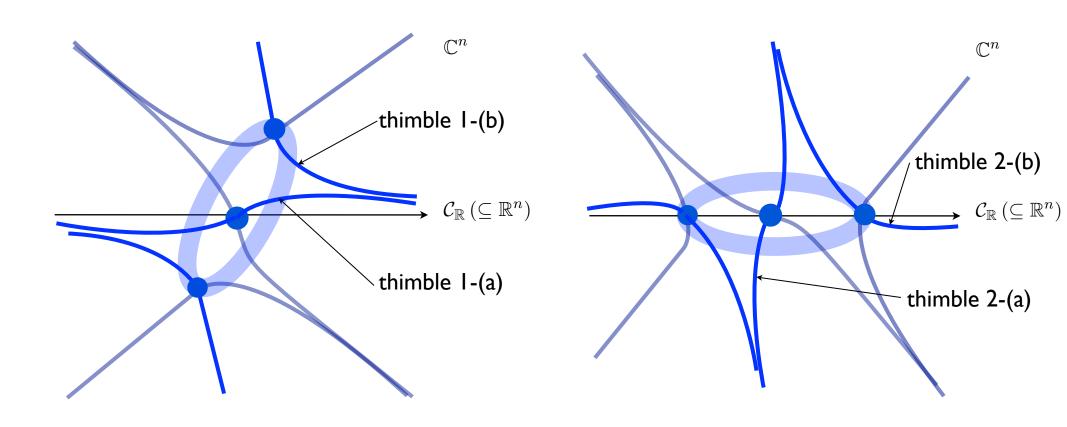
- 1. For $\mu \leq \tilde{\mu}_c$,
 - (a) $z_1 = z_2 = 0$; S[z] = 0,
 - (b) $z_1 = i\phi_0 \cos \theta$, $z_2 = i\phi_0 \sin \theta$; $S[z] = -L^4 \frac{\lambda_0}{4} \phi_0^4$, where $\phi_0 = \sqrt{\frac{+(1-6K_0-2K_0 \cosh(\mu))}{\lambda_0}}$.
- 2. For $\mu > \tilde{\mu}_c$,
 - (a) $z_1 = z_2 = 0$; S[z] = 0,
 - (b) $z_1 = \phi_0 \cos \theta$, $z_2 = \phi_0 \sin \theta$; $S[z] = -L^4 \frac{\lambda_0}{4} \phi_0^4$, where $\phi_0 = \sqrt{\frac{-(1 - 6K_0 - 2K_0 \cosh(\mu))}{\lambda_0}}$.

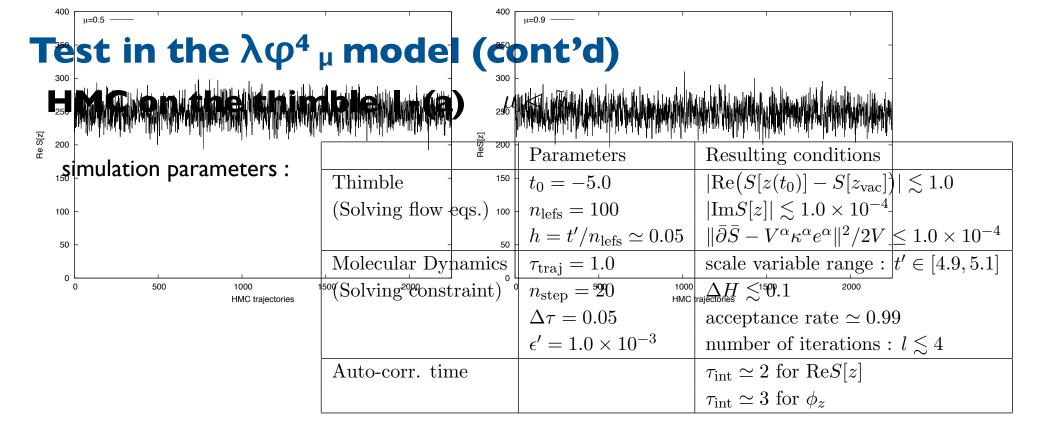
 \rightarrow the thimble I-(a)

 \rightarrow the thimble 2-(b)

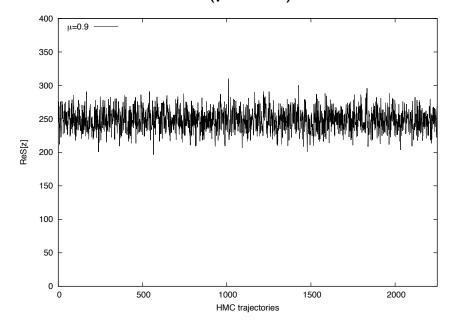
$$\mu < \tilde{\mu}_c$$

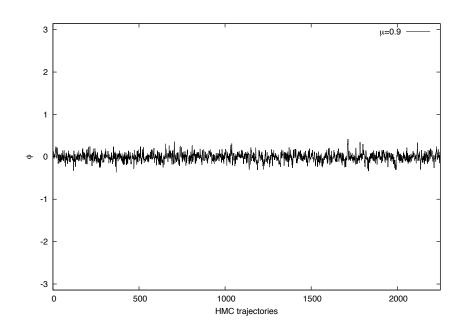
$$\mu > \tilde{\mu}_c$$

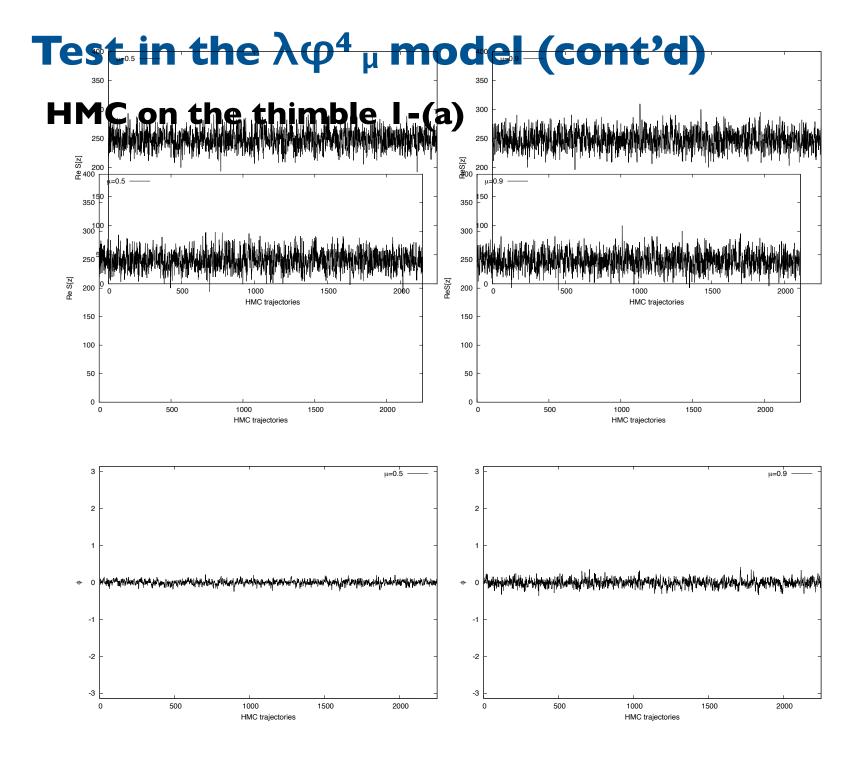




HMC histories ($\mu = 0.9$)







Test in the $\lambda \varphi^4$ model (cont'd)

HMC on the thimble I-(a)

residual phase averages:

HMC trajectories

2000 0 500 HMC trajectories

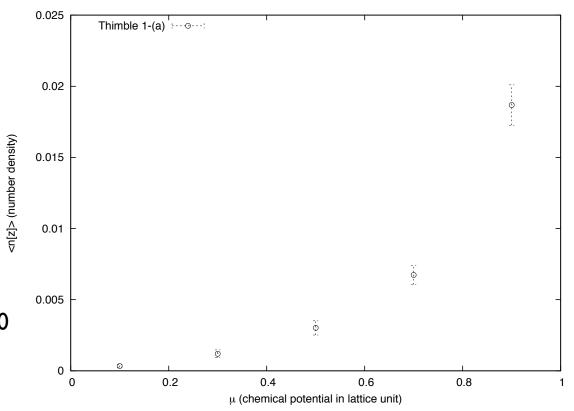
number density:

$$e^{i\phi_z} = \det U_z = \frac{\det V_z}{|\det V_z|}$$

$$n[z] = \frac{1}{L^4} \sum_{x} K_0 z_a(x) z_b(x + \hat{0}) \left[\delta_{ab} \sinh(\mu) - i\epsilon_{ab} \cosh(\mu) \right]$$

μ	$\langle { m e}^{i\phi_z} angle_{{\mathcal J}_{ m vac}}'$		
0.1	$(9.99e-01, -1.15e-03) \pm (5.7e-02, 7.4e-04)$		
0.3	$(9.99e-01, -1.03e-03) \pm (5.7e-02, 2.1e-03)$		
0.5	$(9.98e-01, -2.68e-03) \pm (5.7e-02, 3.3e-03)$		
0.7	$(9.97e-01, 5.24e-04) \pm (5.7e-02, 4.3e-03)$		
0.9	$(9.94e-01, -7.40e-03) \pm (5.7e-02, 5.9e-03)$		

generated 4,250 traj. sampling 300 conf. with the separation of 10



Test in the $\lambda \phi^4 \mu$ model (cont'd)

HMC on the thimble 2-(b)

Critical region of real dimension one : $\theta \in [0, 2\pi]$

$$z_a(x;t) \simeq R_{ab}(\theta) \left\{ \delta_{b1} \phi_0 + \sum_{\beta=1}^{2V-1} v_b(x)^{\beta} \exp(\kappa^{\beta} t) e^{\beta} \right\} \qquad (t \ll 0)$$

$$\delta z_a(x;t) = V_a(x;t)^0 \left(\phi_0 \sqrt{V} \delta \theta\right) + \sum_{\beta=1}^{2V-1} V_b(x;t)^\beta \left(\delta e^\beta + \kappa^\beta e^\beta \delta t\right)$$

Critical fluctuation : lowest mode $\frac{\kappa^1=2\lambda_0\phi_0^2}{v_a(x)^1=\delta_{a1}/\sqrt{V}}$ gets very light ! $(\mu\gtrsim\tilde{\mu}_c)$

$$z_a(x;t) \simeq R_{ab}(\theta) \left\{ \delta_{b1} \frac{\phi_0}{\sqrt{1 - \frac{2}{\sqrt{V}\phi_0}} e^1 \exp(\kappa^1 t)} + \sum_{\beta=2}^{2V-1} v_b(x)^{\beta} \exp(\kappa^{\beta} t) e^{\beta} \right\} \qquad \sum_{\beta=2}^{2V-1} e^{\beta} e^{\beta} = 2V-2$$

$$V_a(x;t)^0 \simeq R_{ab}(\theta) v_b(x)^0 \frac{1}{\sqrt{1 - \frac{2}{\sqrt{V}\phi_0} e^1 \exp(\kappa^1 t)}},$$

$$V_a(x;t)^1 \simeq R_{ab}(\theta) v_b(x)^1 \frac{\exp(\kappa^1 t)}{\left(1 - \frac{2}{\sqrt{V}\phi_0} e^1 \exp(\kappa^1 t)\right)^{3/2}},$$

$$V_a(x;t)^\beta \simeq R_{ab}(\theta) v_b(x)^\beta \exp(\kappa^\beta t) \qquad (\beta = 2, \dots, 2V - 1)$$

zero mode

$$\kappa^0 = 0$$
$$v_a(x)^0 = \delta_{a2}/\sqrt{V}$$

$$\sum_{\beta=2}^{2V-1} e^{\beta} e^{\beta} = 2V - 2$$

the global flow mode $z_a(x;t) = z_a(t)$

$$\frac{d}{dt}z_a(t) = \bar{\partial}_{ax}\bar{S}[\bar{z}]\big|_{z_a(x;t)=z_a(t)}$$

$$= \lambda_0 (\bar{z}_b(t)\bar{z}_b(t) - \phi_0^2)\bar{z}_a(t)$$

Test in the $\lambda \phi^4 \mu$ model (cont'd)

HMC on the thimble 2-(b) $\mu > \tilde{\mu}_c$

simulation parameters:

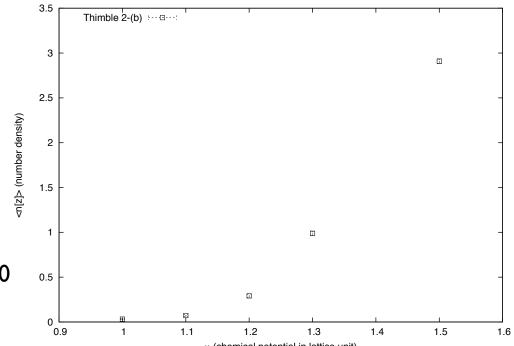
	Parameters	Resulting conditions	
Thimble	$t_0 = -3.0$	$ \operatorname{Re}(S[z(t_0)] - S[z_{\operatorname{vac}}]) \lesssim 2.0 \times 10^1$	
	$n_{\mathrm{lefs}} = 100$	$ \operatorname{Im}(S[z] - S[z_{\text{vac}}]) \lesssim 5.0 \times 10^{-2}$	
	$h = t'/n_{\rm lefs} \simeq 0.03$	$\ \bar{\partial}\bar{S} - V^{\alpha}\kappa^{\alpha}e^{\alpha}\ ^2/2V \le 3.0 \times 10^{-2}$	
MD	$\tau_{\mathrm{traj}} = 0.3$	$t' \in [2.5, 3.5]$	
	$n_{\text{step}} = 10, 30 \ (\mu = 1.0, 1.1)$	$\Delta H \lesssim 0.05$	
	$\Delta \tau = 0.03, 0.01 \ (\mu = 1.0, 1.1)$	Acceptance rate $\simeq 0.99$	
	$\epsilon' = \sqrt{10} \times 10^{-3}$	$l \lesssim 4, 6 \ (\mu = 1.0), 14 \ (\mu = 1.1)$	
Auto-corr. time	(for $ReS[z]$)	$\tau_{\rm int} \simeq 10, 14 \; (\mu = 1.0, 1.1)$	
	$(\text{for }\phi_z)$	$\tau_{\text{int}} \simeq 15, 14 \ (\mu = 1.0), 28 \ (\mu = 1.1)$	

residual phase averages:

μ	$\langle { m e}^{i\phi_z} angle_{{\cal J}_{ m vac}}'$		
1.0	$(9.94e-01, -8.77e-03) \pm (3.1e-02, 3.1e-03)$		
1.1	$(9.94e-01, -3.21e-03) \pm (3.1e-02, 3.4e-03)$		
1.2	$(9.95e-01, -8.25e-04) \pm (3.1e-02, 3.0e-03)$		
1.3	$(9.97e-01, -3.08e-03) \pm (3.1e-02, 2.2e-03)$		
1.5	$(9.99e-01, -1.06e-03) \pm (3.1e-02, 1.0e-03)$		

generated 11,250 traj. sampling 1,000 conf. with the separation of 10

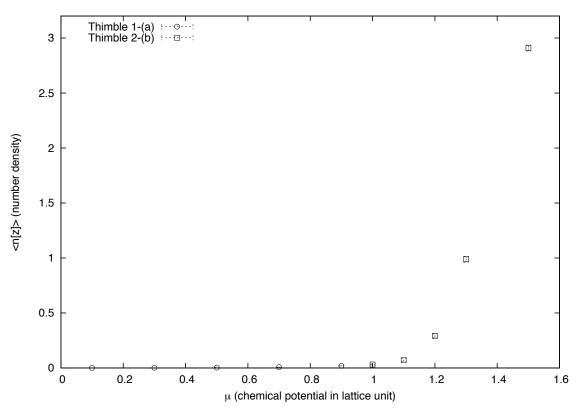
number density:



Test study in the $\lambda \phi^4$ μ model (contid) μ (chemical potential in lattice unit)

HMC on the thimbles I-(a) & 2-(b)

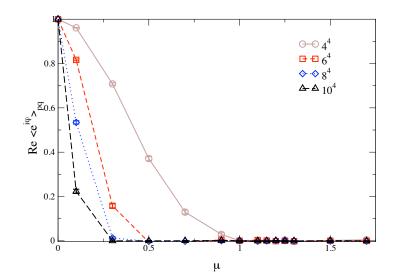
number density:



μ	Re $\langle n[z] \rangle_{\mathcal{J}_{\text{vac}}}$ (jk. error)	Re $\langle e^{i\phi_z} n[z] \rangle'_{\mathcal{J}_{vac}}$	Re $\langle n[z] \rangle'_{\mathcal{J}_{\text{vac}}}$
0.1	3.34e-04 (9.2e-05)	3.35e-04	2.15e-04
0.3	1.20e-03 (2.7e-04)	1.19e-03	8.56e-04
0.5	3.02e-03 (5.0e-04)	3.01e-03	2.44e-03
0.7	6.74e-03 (6.7e-04)	6.71e-03	5.91e-03
0.9	1.89e-02 (1.4e-03)	1.85 e-02	1.73e-02
1.0	3.14e-02 (4.3e-03)	3.12e-02	3.00e-02
1.1	7.17e-02 (1.3e-02)	7.12e-02	7.01e-02
1.2	2.92e-01 (1.8e-02)	2.90e-01	2.90e-01
1.3	9.88e-01 (2.6e-02)	9.85 e-01	9.87e-01
1.5	2.91e-00 (2.7e-02)	2.90e-00	2.90e-00

HOW SEVERE IS THE SIGN PROBLEM

AVERAGE PHASE FACTOR



Test in the $\lambda \phi^4 \mu$ model (cont'd)

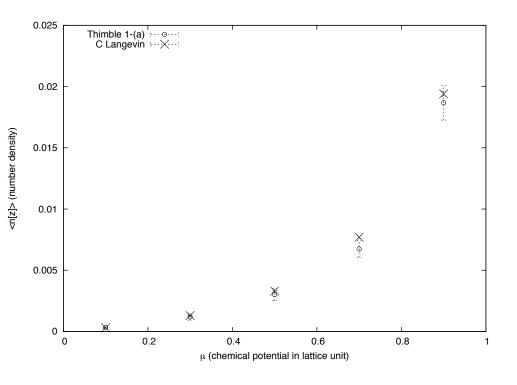
Comparison to Complex Langevin simulations

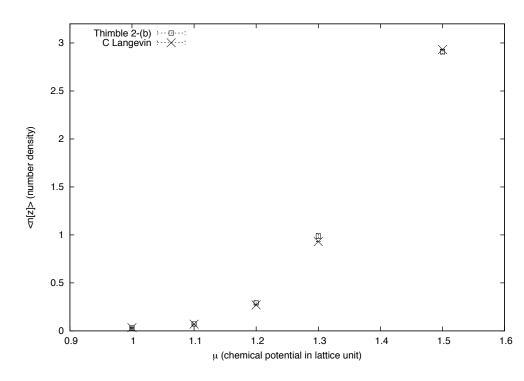
$$\frac{dz(t)}{dt} = -\frac{\partial S[z]}{\partial z} + \eta(t); \quad \langle \eta(t)\eta(t') \rangle = 2\delta(t - t')$$

$$\langle \mathcal{O} \rangle = \lim_{t \to \infty} \frac{1}{t} \int_0^t dt' \, \mathcal{O}(z(t'))$$

parameters of CL simulations: step size ϵ =5.0 × 10⁻⁵ , 5,000,000 time steps sampling 10,000 configurations with the separation of 500

number density:





Summary & Discussions

- We have formulated a HMC algorithm which is applicable to lattice models defined on Lefschetz thimbles
- We have tested the algorithm in the $\lambda \phi^4$ μ model on the lattice V=4⁴
 - the thimbles associated with the classical vacua
 - the residual phase factors reweighted successfully
 - known results of the number density reproduced (cf. CL)
 - Need the careful study of the systematic errors
 - setup of the asymptotic regions
 - contributions of other thimbles, ex. thimble 2-(a), ...
 - Need the study of the residual sign problem on larger lattices
 - numerical cost per traj.
 - $O(V^2 \times n_{Lefs} \times n_{step})$ (tangent vectors), $O(V^3 \times n_{step})$ (V-1, detV)
- Possible applications to QCD μ cf. D. Sexty, arXiv:1307.7748