

# 核力から出発した非現象論的な中重原子核構造 の研究~多体摂動論と大規模殻模型計算を用いて

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2017/02/16, 「京からポスト京に向けて」 シンポジウム

N.T., K. Takayanagi, M. Hjorth-Jensen, and T. Otsuka, Phys. Rev. C **89**, 024313 (2014).

N.T., T. Otsuka, N. Shimizu, M. Hjorth-Jensen, K. Takayanagi, T. Suzuki, Phys. Rev. C. Rapid. accepted.

## ✳ Introduction

✳ Construction of Effective interaction

✳ Application to island of inversion (京でやったこと)

✳ Application to pf+sdg-shell nuclei (ポスト京に向けて)

✳ Conclusion

# Construction of “Bridges”

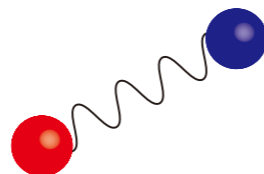
QCD



Lattice QCD

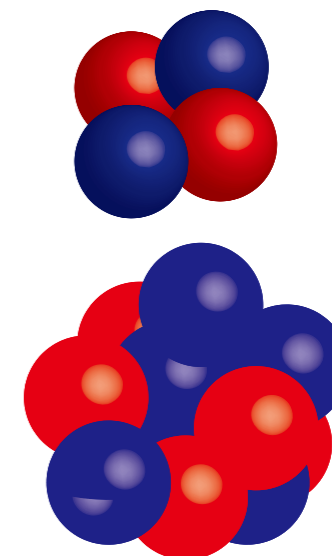
Effective Field Theory

Nuclear force



Few body techniques  
No core shell model  
and many others...

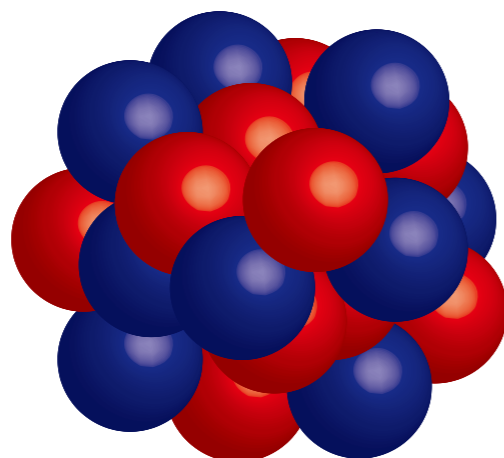
Light nuclei  $\sim A \approx 10-20$



This work



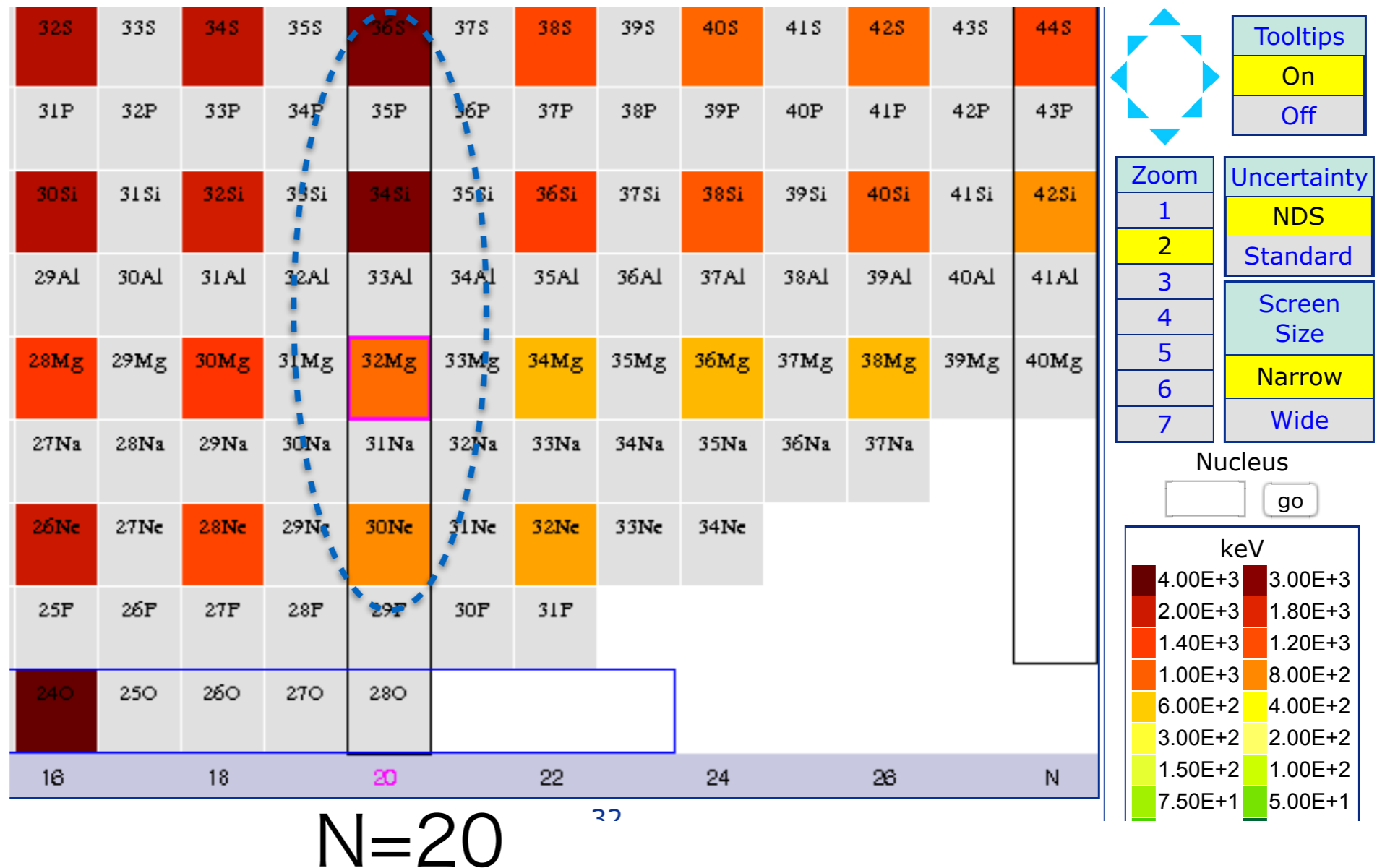
Medium mass nuclei  $\sim A \approx 20-100$



shell model with core  
via the **effective interaction**  
derived from **nuclear force**

# Neutron-rich nuclei~ island of inversion

<http://www.nndc.bnl.gov/nudat2/reCenter.jsp?z=12&n=20>



- $E(2+) \sim 1$  MeV on  $N=20$  indicate **breaking** of major shell gap
- Unified treatment of beyond and below the  $N=20$  gap is necessary
- And this is one of many examples...



# Nuclear force and Nuclear shell model

## Single particle energies

Shell model Hamiltonian

$$H = \sum_i \epsilon_i a_i^\dagger a_i + \sum_{ijkl} V_{ij,kl} a_i^\dagger a_j^\dagger a_l a_k.$$

two-body matrix elements

$i, j, k, l$  : relevant degrees of freedom = **model space**

What we aim at is to:

- **derive** Shell model Hamiltonian based on **Nuclear force**
- **diagonalize this Hamiltonian** to know nuclear properties

What we need is:

- **Many-body theory** to derive Hamiltonian for **desired** model space
  - neutron-rich nuclei need large model space
- **Large scale shell model calculation (Large scale Lanczos diagonalization and Monte Carlo Shell model)**

✱ Introduction

✱ **Theory of Effective interaction**

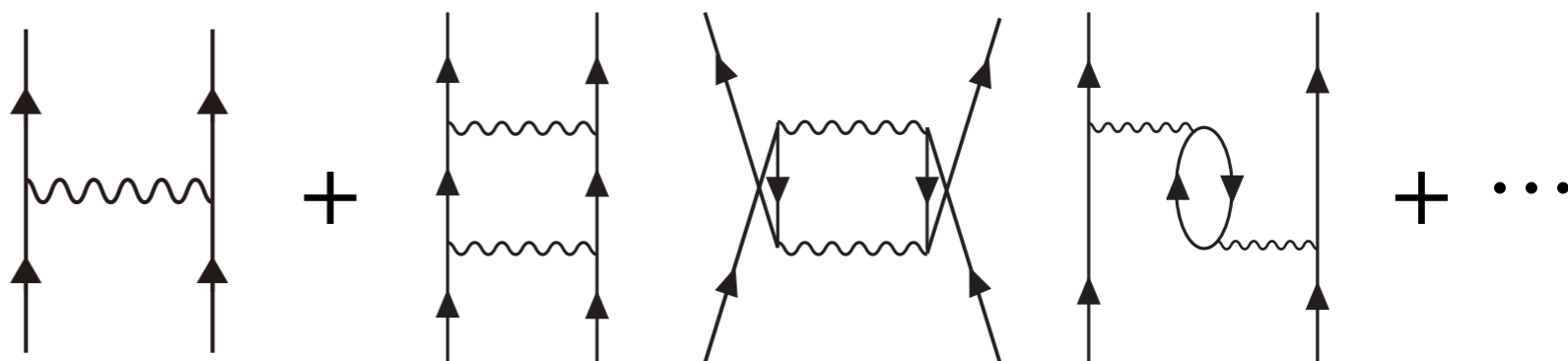
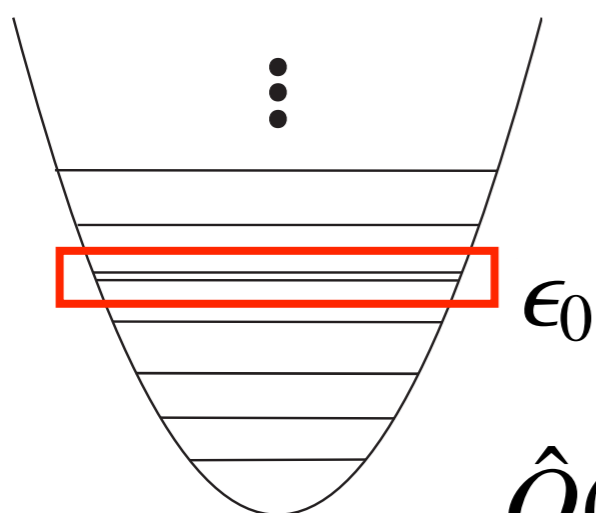
✱ Application to island of inversion (京でやったこと)

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# Many-body perturbation theory

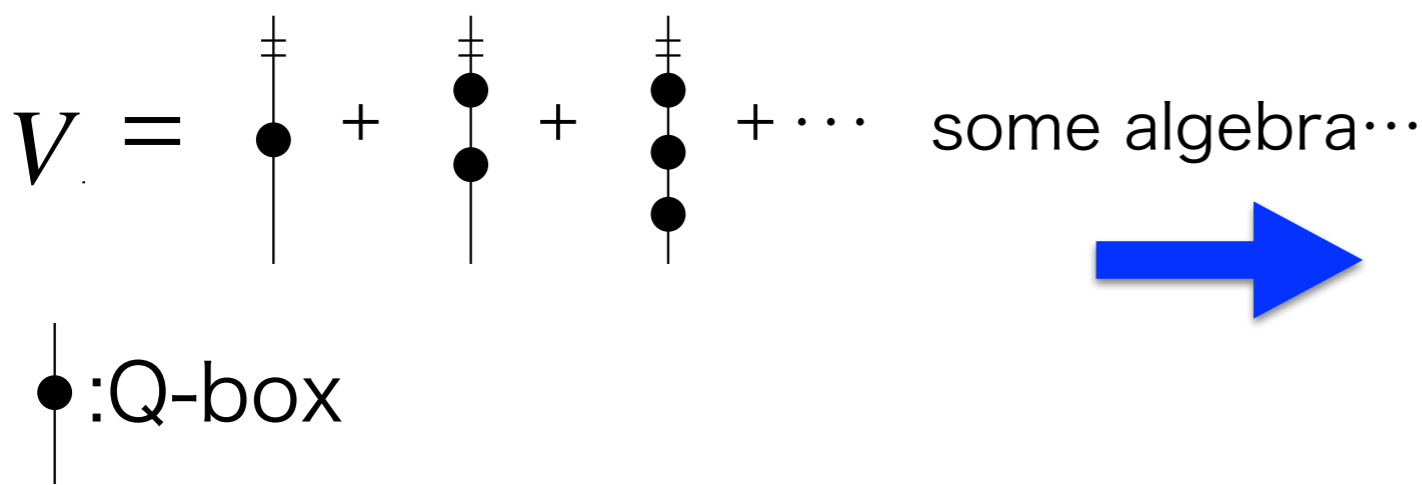
P: model space



$$\hat{Q}(E) = PVP + PVQ \frac{1}{E - QHQ} QVP$$

Q-box

As non-perturbative correction, we further include infinite repetition of Q-box



non-perturbative correction

$$V_{\text{eff}}^{(n)} = \hat{Q}(\epsilon_0) + \sum_{k=1}^{\infty} \hat{Q}_k(\epsilon_0) \{V_{\text{eff}}^{(n-1)}\}^k$$

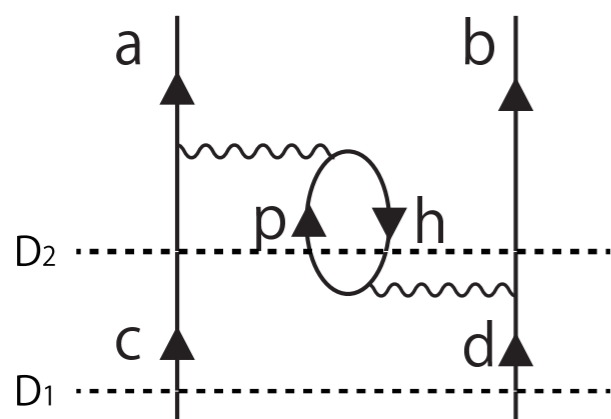
$$\hat{Q}_k(E) = \frac{1}{k!} \frac{d^k \hat{Q}(E)}{dE^k}$$

# Divergent problem of Q-box in non-degenerate model space

(A) KK method requires assumption that the model space is **degenerate**

(B) Naive perturbation theory leads a **divergence** in non-degenerate model space

Example



$$\frac{V_{ah,cp} V_{pb,hd}}{(\epsilon_c + \epsilon_d) - \epsilon_c - \epsilon_p + \epsilon_h - \epsilon_b}$$

→ Energy denominator is zero  
when  $\epsilon_d - \epsilon_b = \epsilon_p - \epsilon_h$

We need a theory which satisfies

(a) The assumption of degenerate model space is **removed**

(b) **Avoid** the divergence appearing in Q-box diagrams

→ **EKK method** as a re-summation scheme of KK method

# Extended KK method as a re-summation of the perturbative series

## EKK method

New parameter  $E$  (arbitrary parameter)

$$\begin{aligned} H &= H'_0 + V' \\ &= \begin{pmatrix} E & 0 \\ 0 & QH_0Q \end{pmatrix} + \begin{pmatrix} P\tilde{H}P & PVQ \\ QVP & QVQ \end{pmatrix}, \end{aligned}$$

$$H_{\text{BH}}(E) = PHP + PVQ \frac{1}{E - QH_0Q} QVP.$$

$$\tilde{H}_{\text{eff}}^{(n)} = \tilde{H}_{\text{BH}}(E) + \sum_{k=1}^{\infty} \hat{Q}_k(E) \{\tilde{H}_{\text{eff}}^{(n-1)}\}^k.$$

## KK method (conventional)

$$\begin{aligned} H &= H_0 + V \\ &= \begin{pmatrix} PH_0P & 0 \\ 0 & QH_0Q \end{pmatrix} + \begin{pmatrix} PVP & PVQ \\ QVP & QVQ \end{pmatrix} \end{aligned}$$

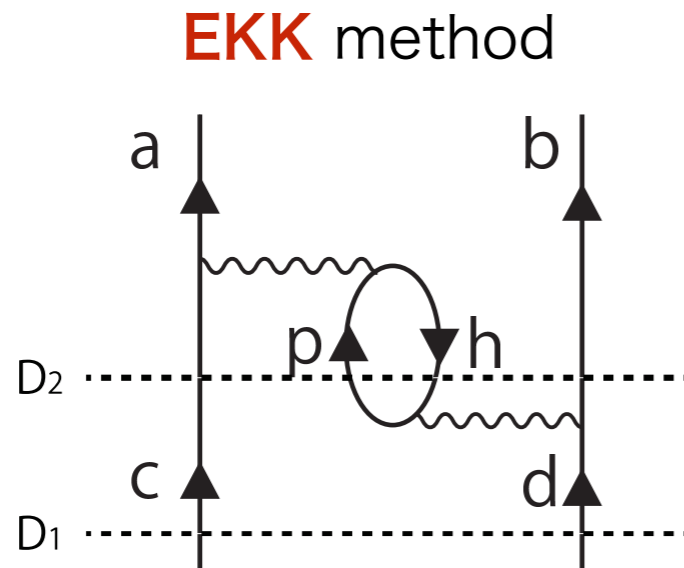
$$\hat{Q}(E) = PVP + PVQ \frac{1}{E - QH_0Q} QVP.$$

$$V_{\text{eff}}^{(n)} = \hat{Q}(\epsilon_0) + \sum_{k=1}^{\infty} \hat{Q}_k(\epsilon_0) \{V_{\text{eff}}^{(n-1)}\}^k.$$

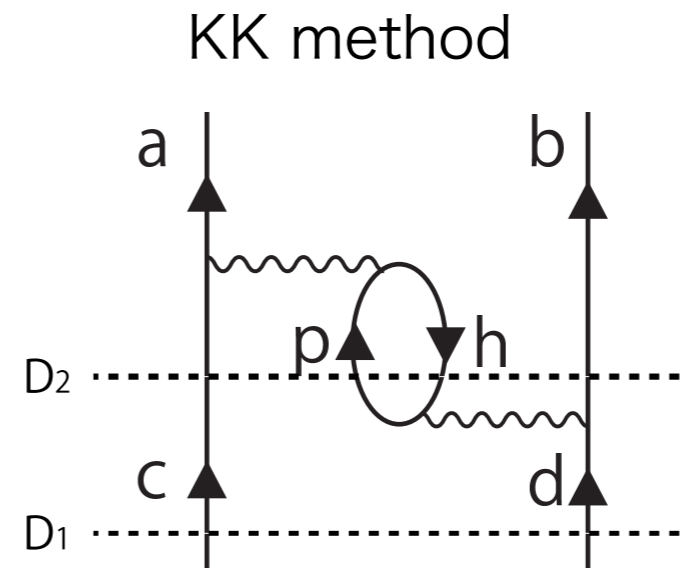
- EKK method can be interpreted as a re-summation of KK method
- All the arguments are kept unchanged with the new division of the Hamiltonian

N. Tsunoda, K. Takayanagi, M. Hjorth-Jensen, and T. Otsuka, Phys. Rev. C 89, 024313 (2014).

# Example: EKK method avoids the divergences



$$E - \frac{V_{ah,cp} V_{pb,hd}}{\epsilon_c - \epsilon_b - \epsilon_p + \epsilon_h}$$

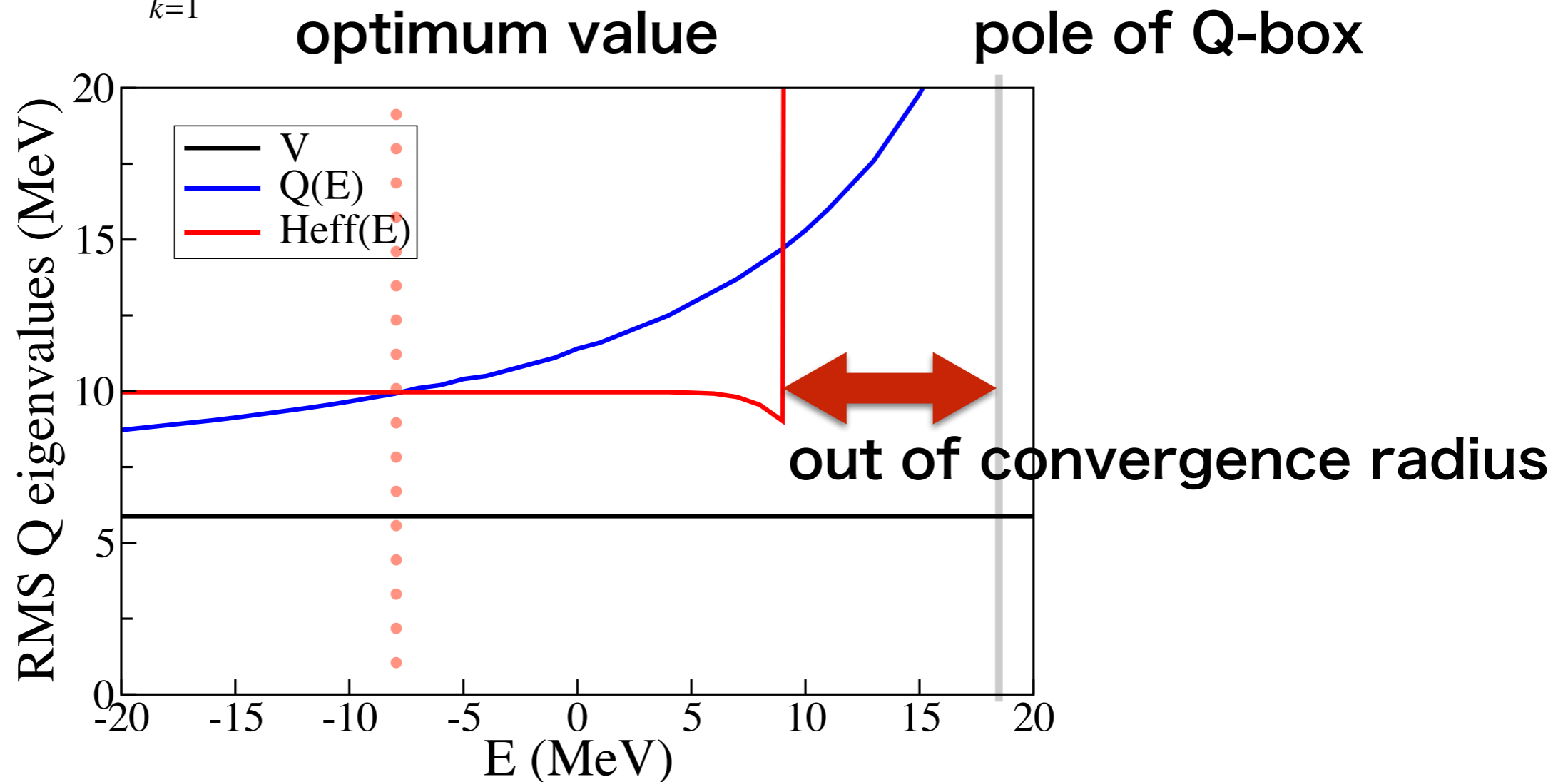


$$\frac{V_{ah,cp} V_{pb,hd}}{(\epsilon_c + \epsilon_d) - \epsilon_c - \epsilon_p + \epsilon_h - \epsilon_b}$$

- We can choose E to avoid divergence !
- Note that the choice of E is arbitrary and should give the same result if the Q-box is calculated without any approximation.
- Inversely, E-dependence is a **measure of error** coming from the approximation

# E-dependence w and w/o non-perturbative correction

$$\tilde{H}_{\text{eff}}^{(n)} = \tilde{H}_{\text{BH}}(E) + \sum_{k=1}^{\infty} \hat{Q}_k(E) \{\tilde{H}_{\text{eff}}^{(n-1)}\}^k.$$



- ★ Non perturbative correction vanishes the E-dependence
- ★ Optimum value of E

# EKK code algorithm

**SPEs**      **TBMEs**      所要時間目安

$$H = \sum_i \epsilon_i a_i^\dagger a_i + \sum_{ijkl} V_{ij,kl} a_i^\dagger a_j^\dagger a_l a_k \cdot \text{(single node 換算)}$$

read bare interaction from file

~10 s

calculate SPEs ~ o(10) elements

~5 s

calculate  $\epsilon_i$

calculate TBMEs  
~ o(10<sup>2</sup>~10<sup>4</sup>) elements

**10 h~100 days** calculate  $\hat{Q}_k(E)$

calculate folded diagrams

~5 s

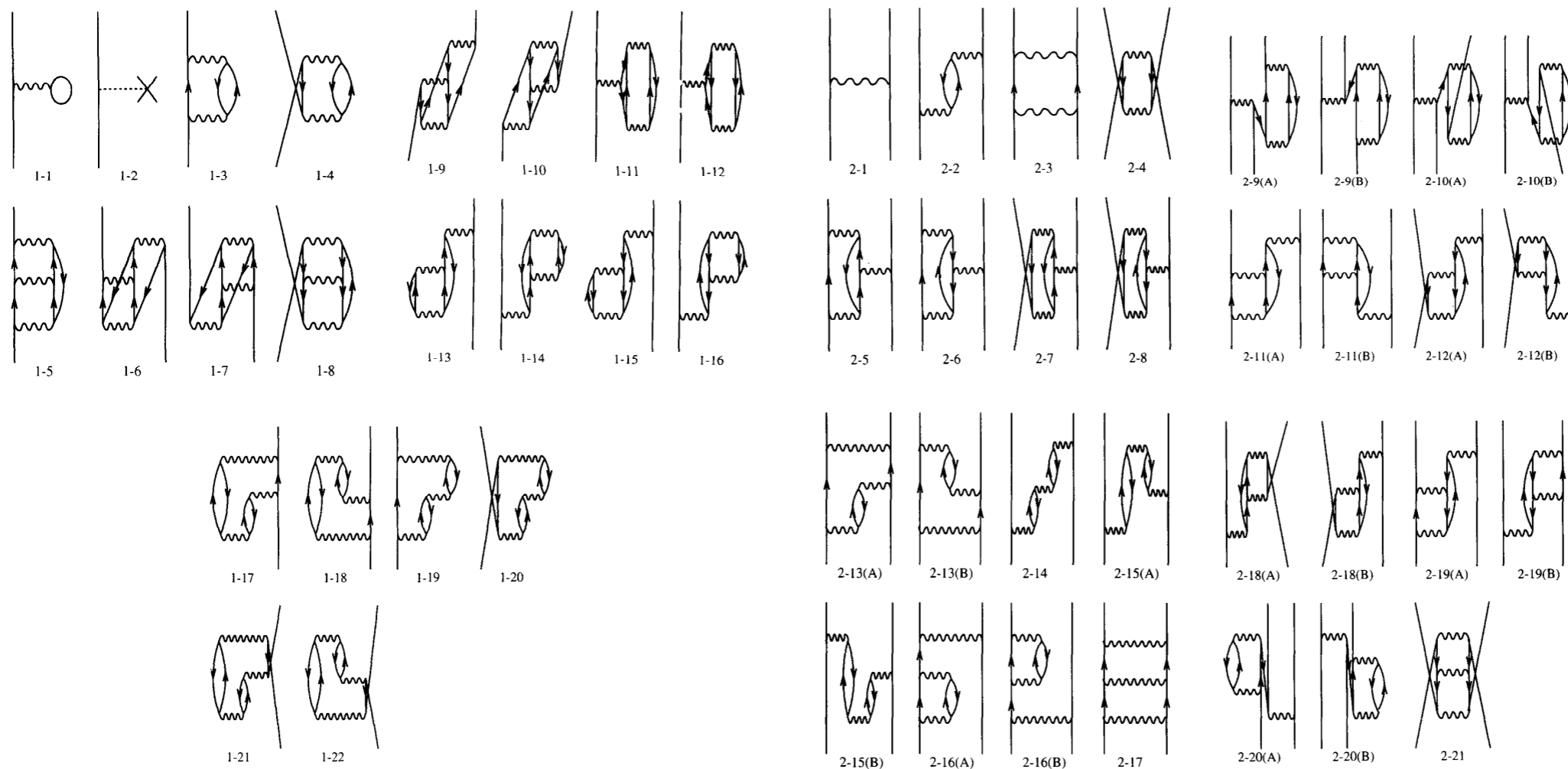
iteration

$$\tilde{H}_{\text{eff}}^{(n)} = \tilde{H}_{\text{BH}}(E) + \sum_{k=1}^{\infty} \hat{Q}_k(E) \{ \tilde{H}_{\text{eff}}^{(n-1)} \}^k$$



# EKK code algorithm

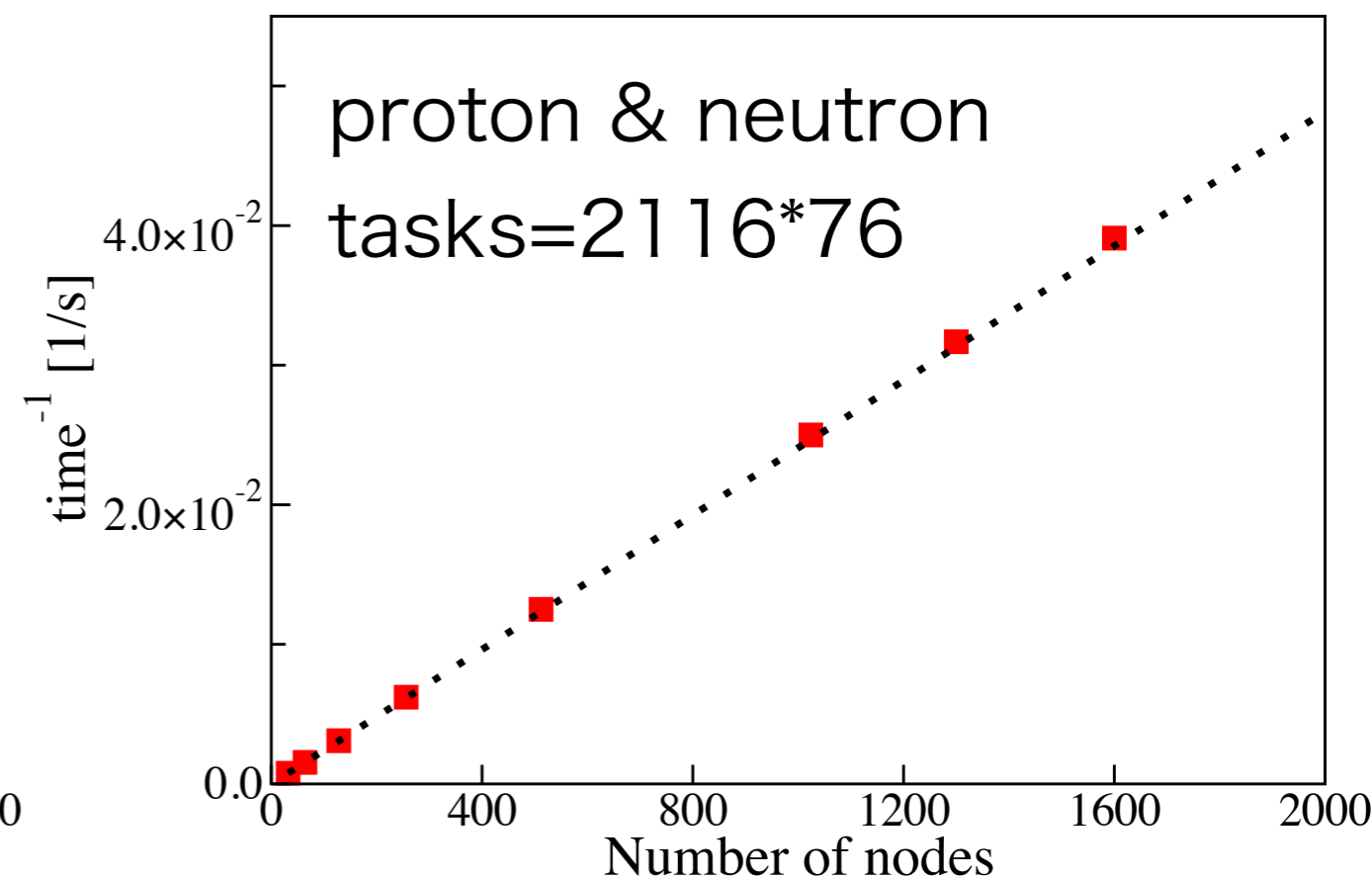
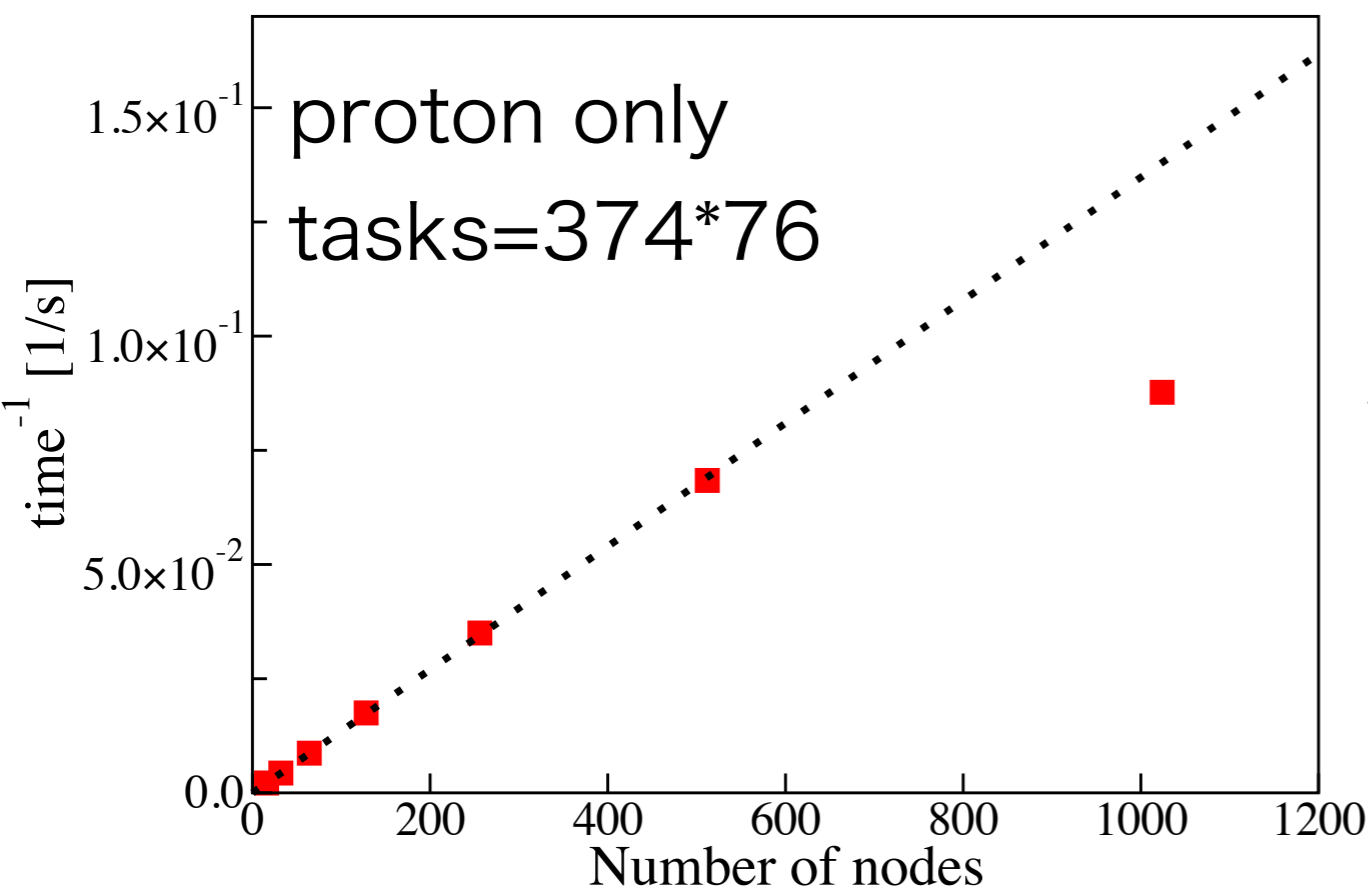
About 50 diagrams (76 tasks) to be calculated for **each** TBMEs (2116 for sd<sub>pf</sub>-shell, 6380 for pf<sub>sdg</sub>-shell) → MPI parallel calculation (master-slave scheme)



Each diagram → openMP parallel calculation for particle loop

# MPI+openMP scaling

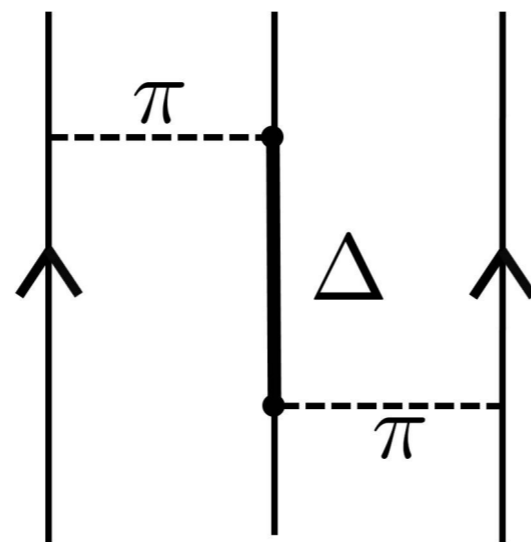
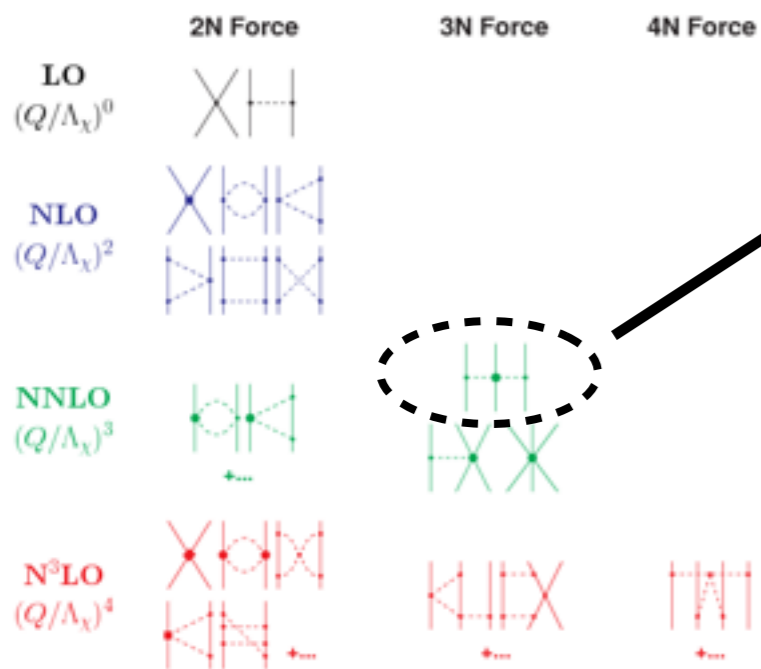
Test case: sdpf-shell  
13major shells  
oakforest (64 core, max 2048 nodes)



# of nodes ~ # of TBMEs までは良い効率で並列化可能  
(グラフでは)

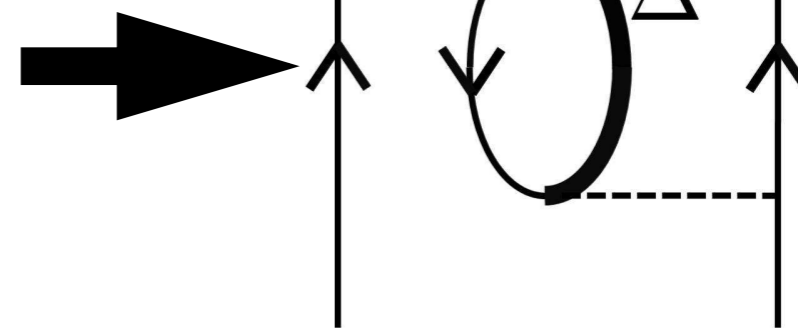
# 3N interaction ( $\Delta$ -hole interaction)

main contribution of 3N



Fujita-Miyazawa type  
**3N** interaction

summation with hole state



Effective  
2N interaction

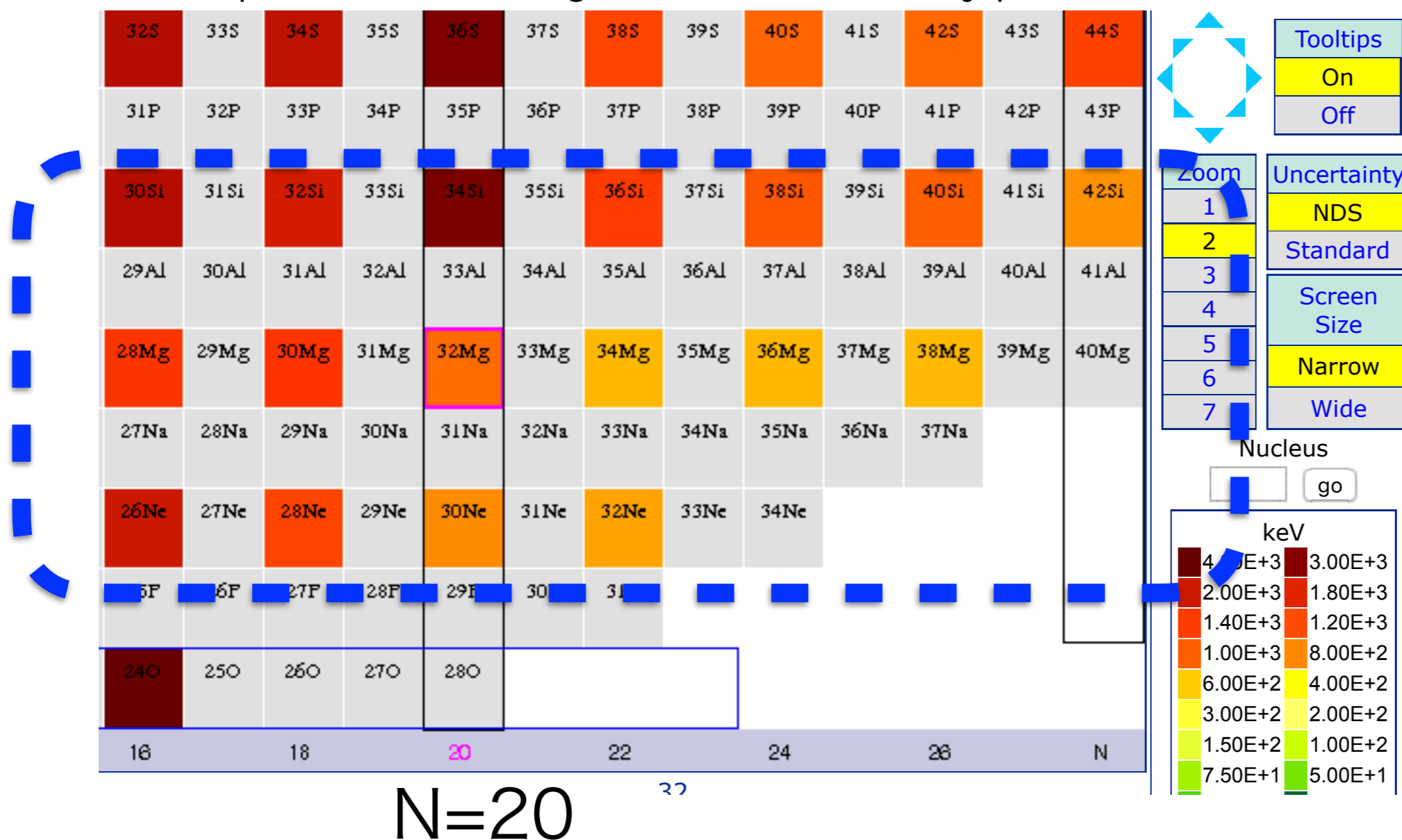
- Adding up effective 2N interaction derived from 3N interaction to EKK 2N effective interaction [1]
- This is one of the lowest order interaction from 3N force and for higher order we are working on...

[1] T. Otsuka, T. Suzuki, J. D. Holt, A. Schwenk, and Y. Akaishi, Phys. Rev. Lett. 105, 032501 (2010).

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- ✱ Conclusion

# Island of inversion

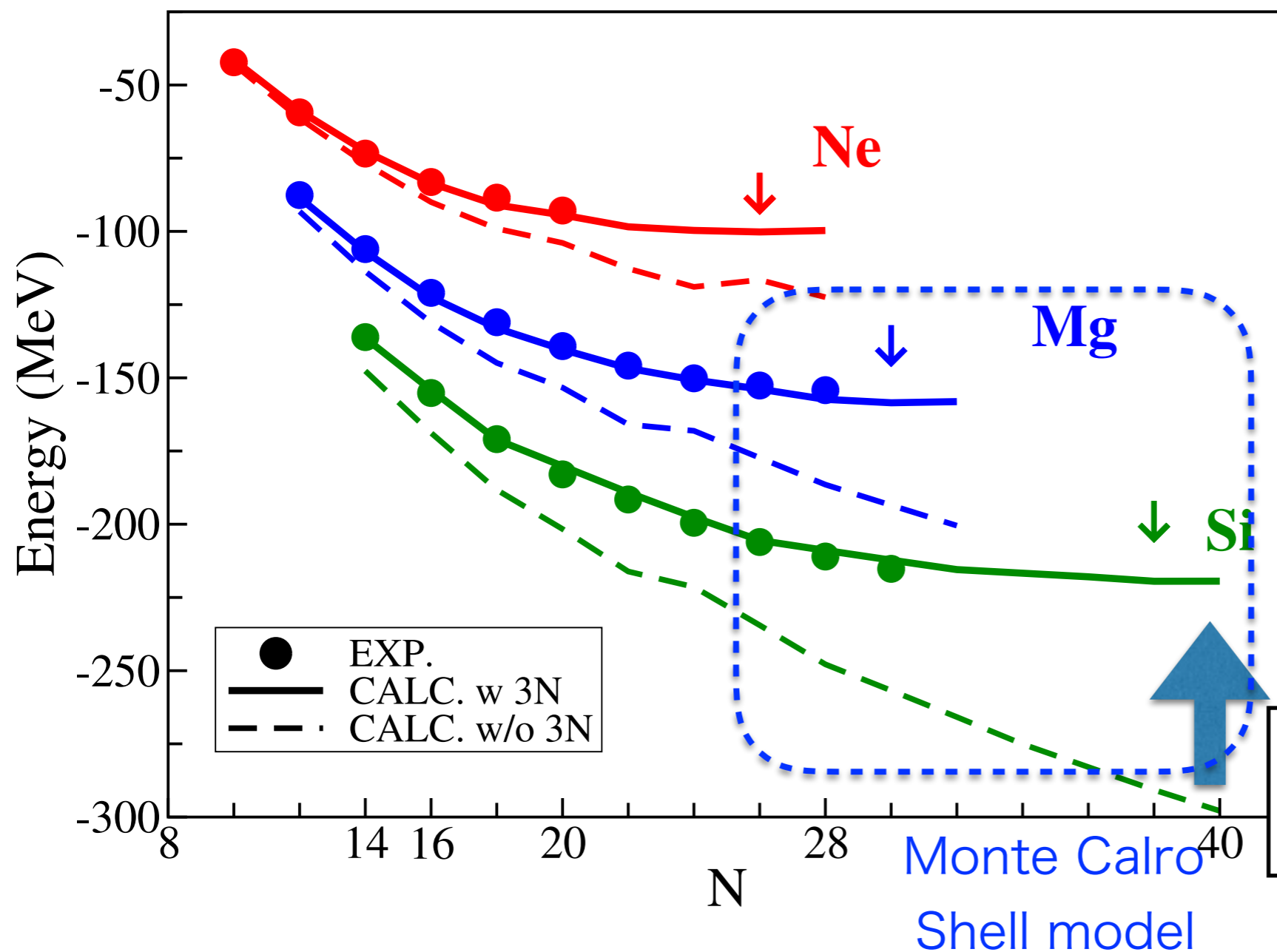
<http://www.nndc.bnl.gov/nudat2/reCenter.jsp?z=12&n=20>



N=20

- Around Ne and Mg region N=20 major gap disappears. (small 2+ energy for even-even nuclei, large deformation, etc...)
- Ground state is consist of “**inverse**” configuration, i.e. intruder configuration
- Can microscopic theory describe this disappearance of major magic number?

# Ground state energies and dripline



Many body perturbation  
up to 3rd order  
starting from N3LO  
P+Q: 17hw (converged)  
SPE fitted

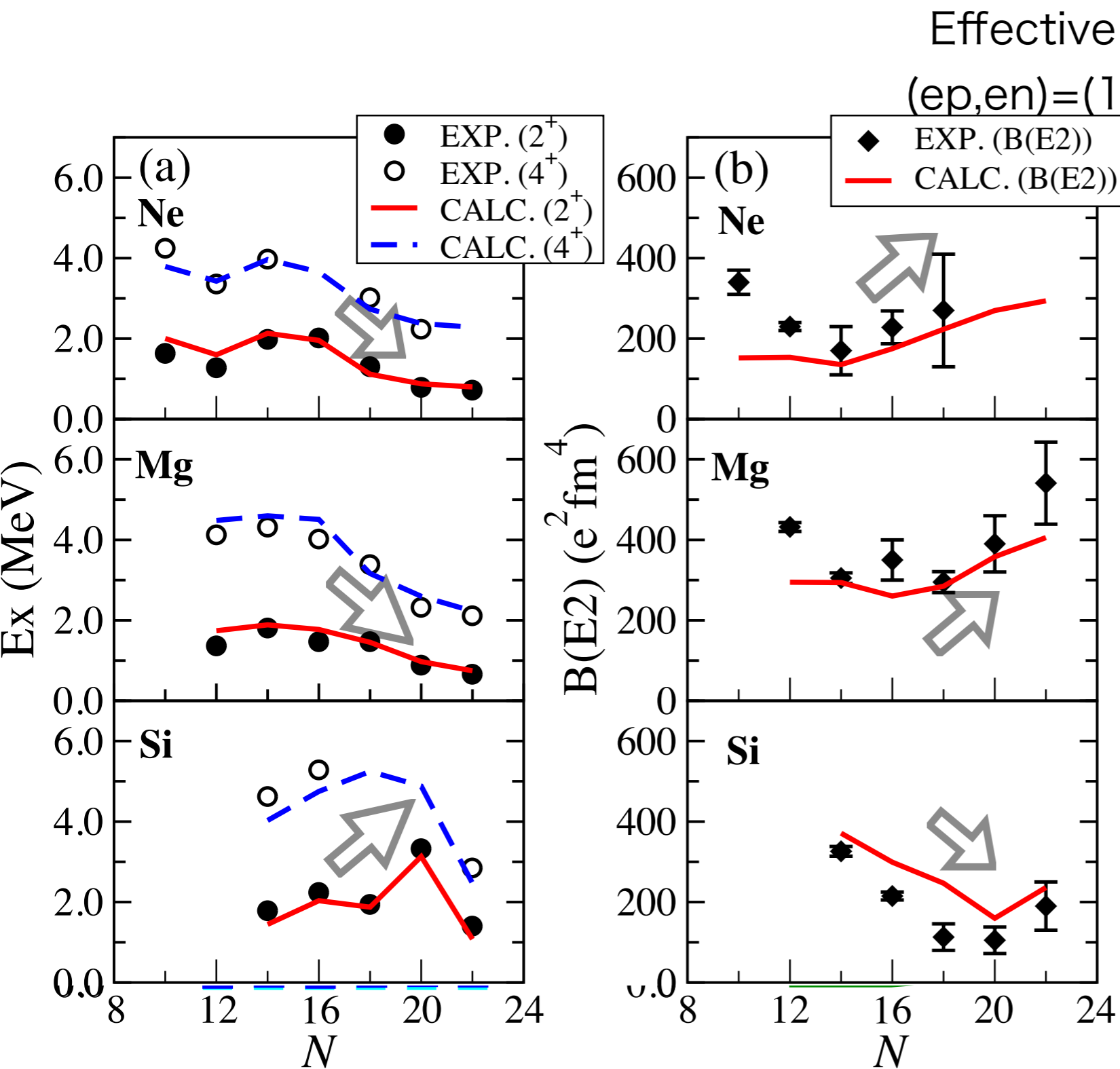
## Resource:

$\sim 10^6$  node\*hour  
in K-computer for  
around **100 states**

Contribution  
from 3N force

- Contribution of 3N force is significant in neutron-rich nuclei
- Predictions of dripline
- Combination of **Microscopic theory** and **Large scale calc.**

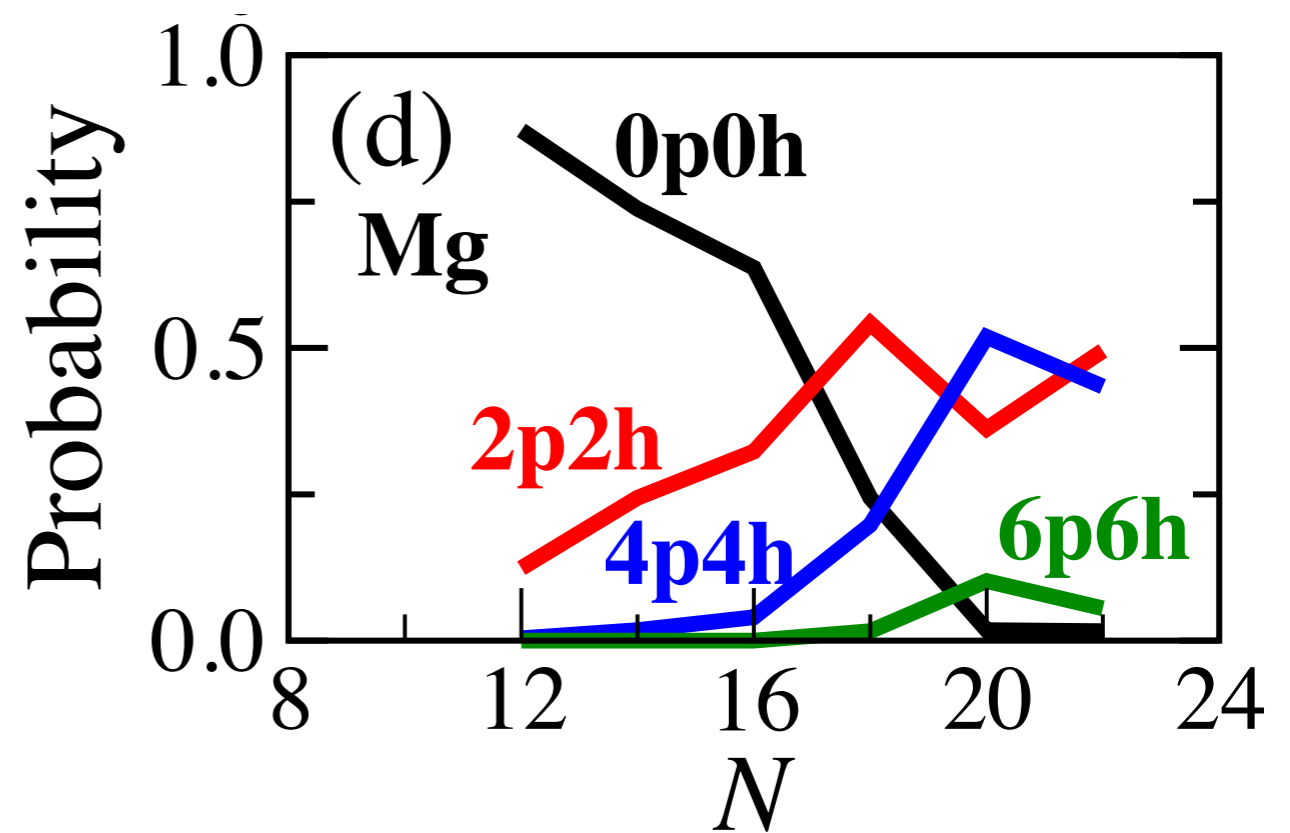
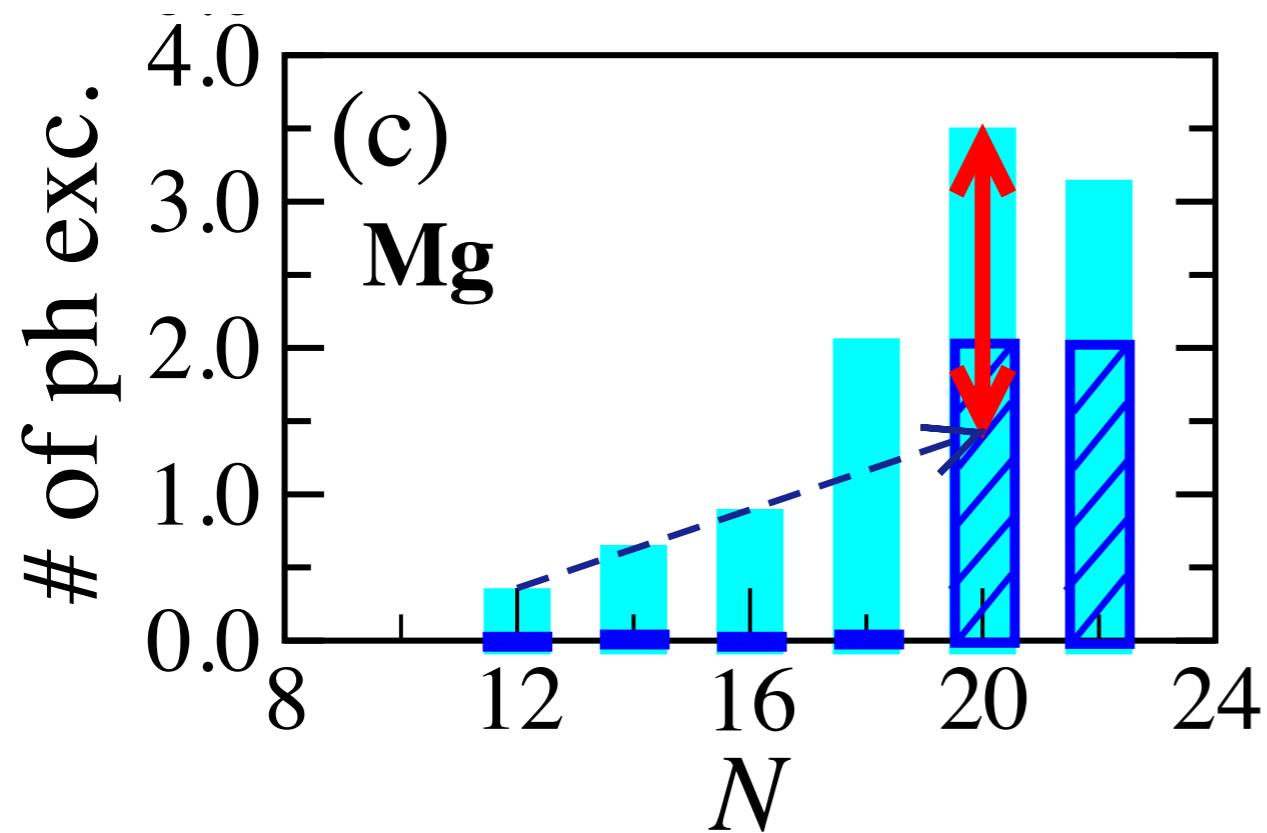
# Shell structure in “island of inversion”



Clear indication of breaking of  $N=20$  gap for Ne and Mg.

$N=20$  gap remains in Si case.

# Wave function of Mg isotopes



“modest” and “abrupt” excitation

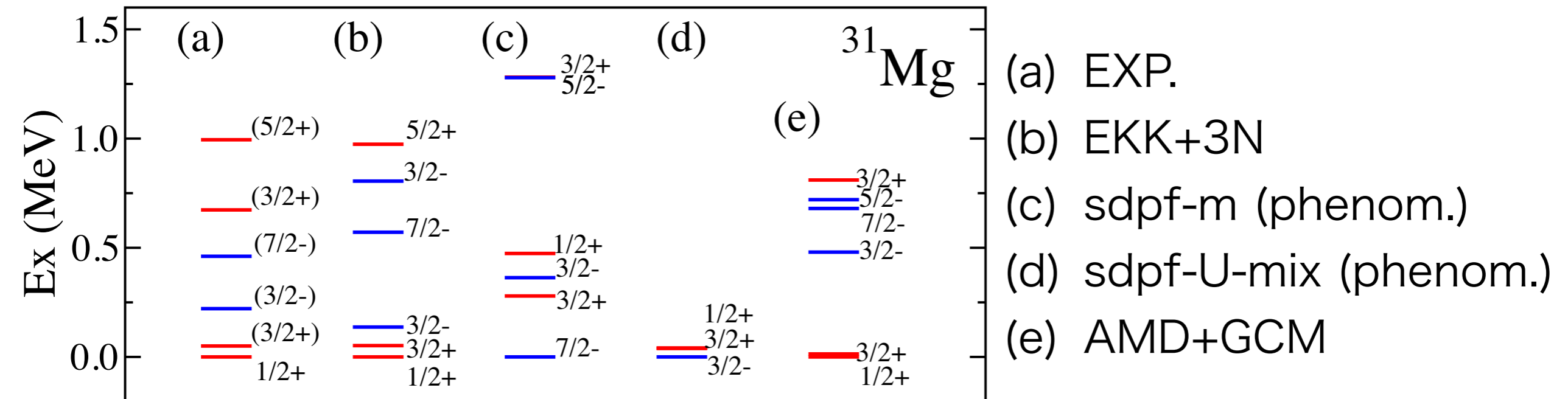
modest: shifting between two shells (e.g. pairing)

abrupt : strong deformation

Abrupt excitation roughly corresponds to conventional 2p2h excitation model

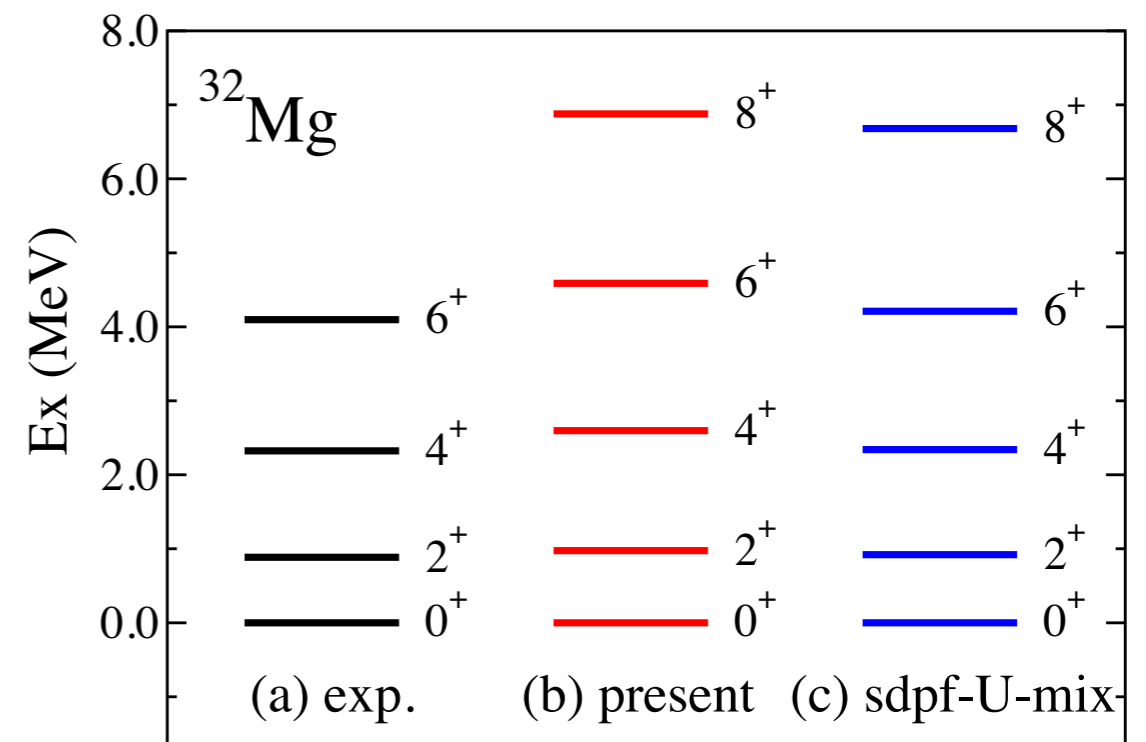


# 31Mg



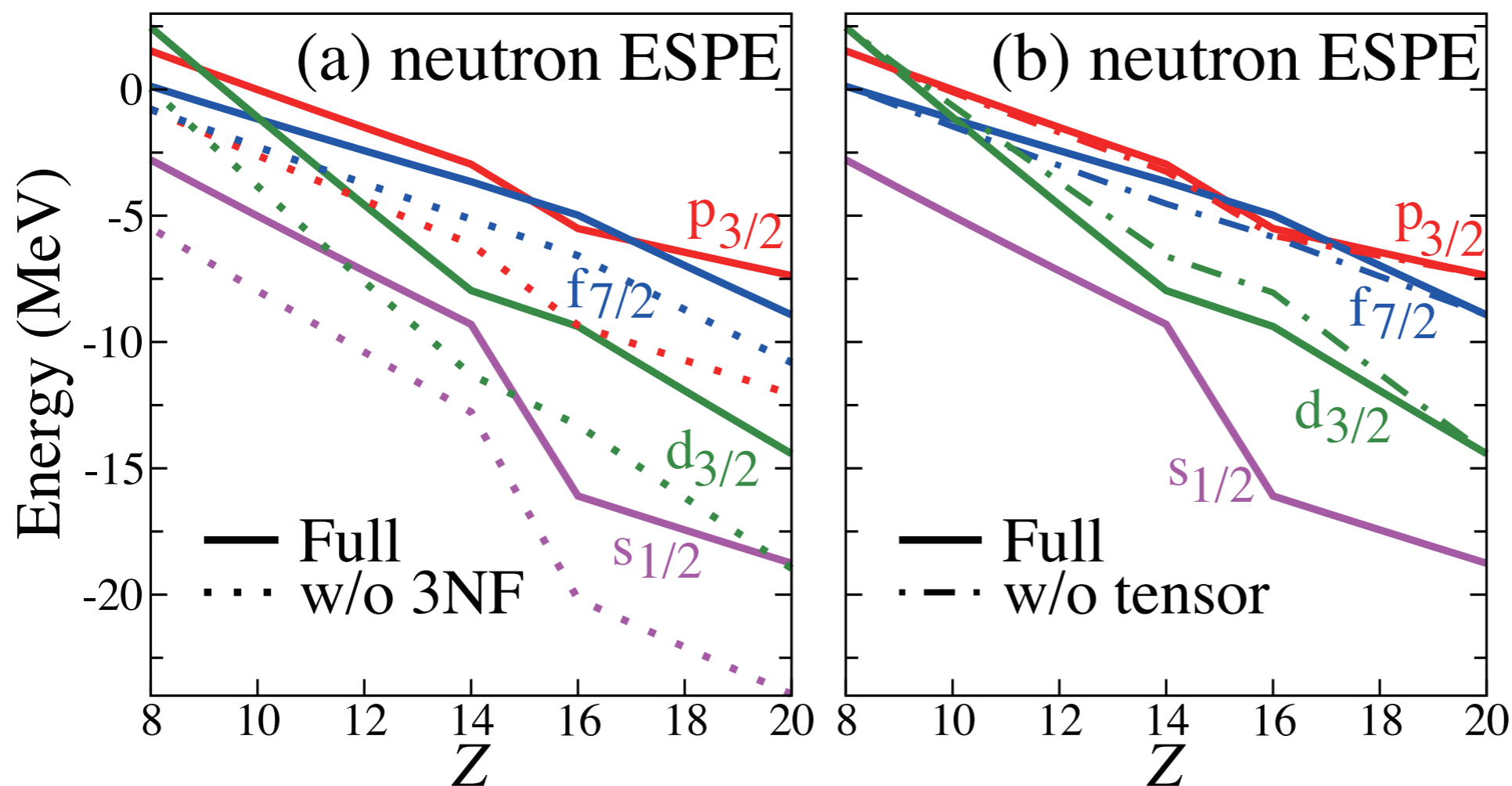
- onset of island of inversion
- ordering of levels reproduced
- positive=2hw dominated
- negative=1hw dominated

## Deformation of $^{32}\text{Mg}$



# Evolution of single particle states

Effective single particle energies at N=20 isotones



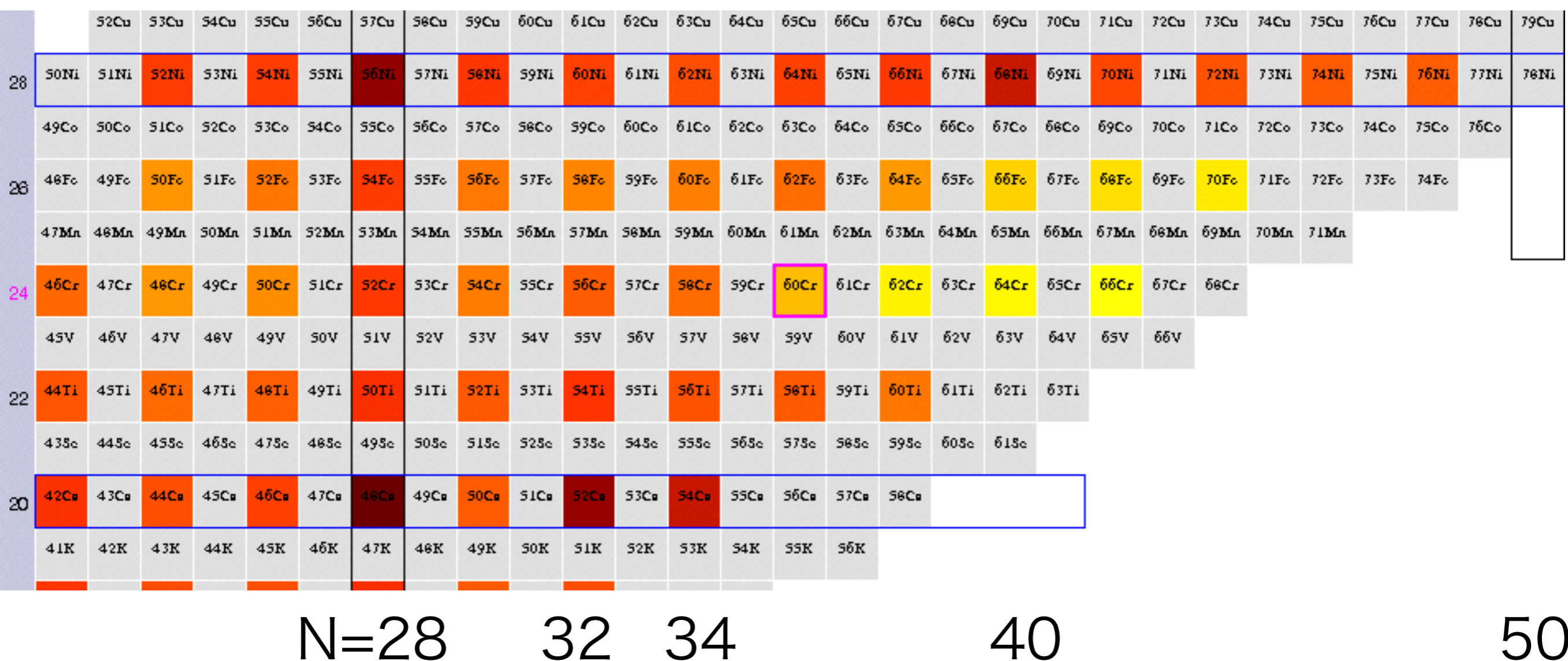
3NF: general shift

Tensor force: drive sd to pf gap

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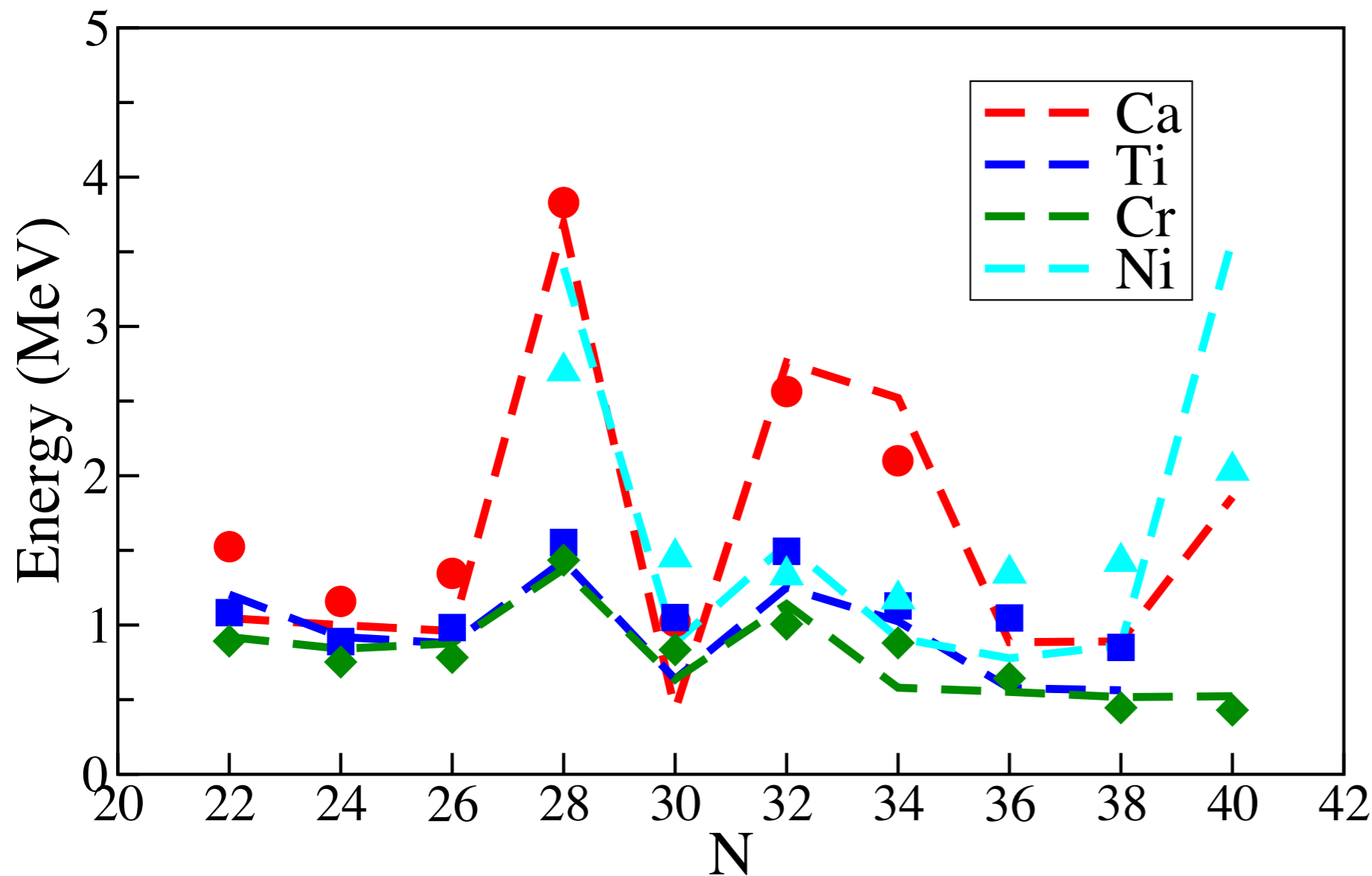
# Ca~Ni isotopes, several magic numbers

## $E^{2+}$ of even-even isotopes



- “Second” island of inversion
- Wider area, larger calculations (ポストト京に向けて)

# $E^{2+}$ of Ca, Ti, Cr and Ni isotopes (preliminary)



## Interaction

pf+sdg-shell

P+Q: 13 major shell

vlowk + 3NF(FM)

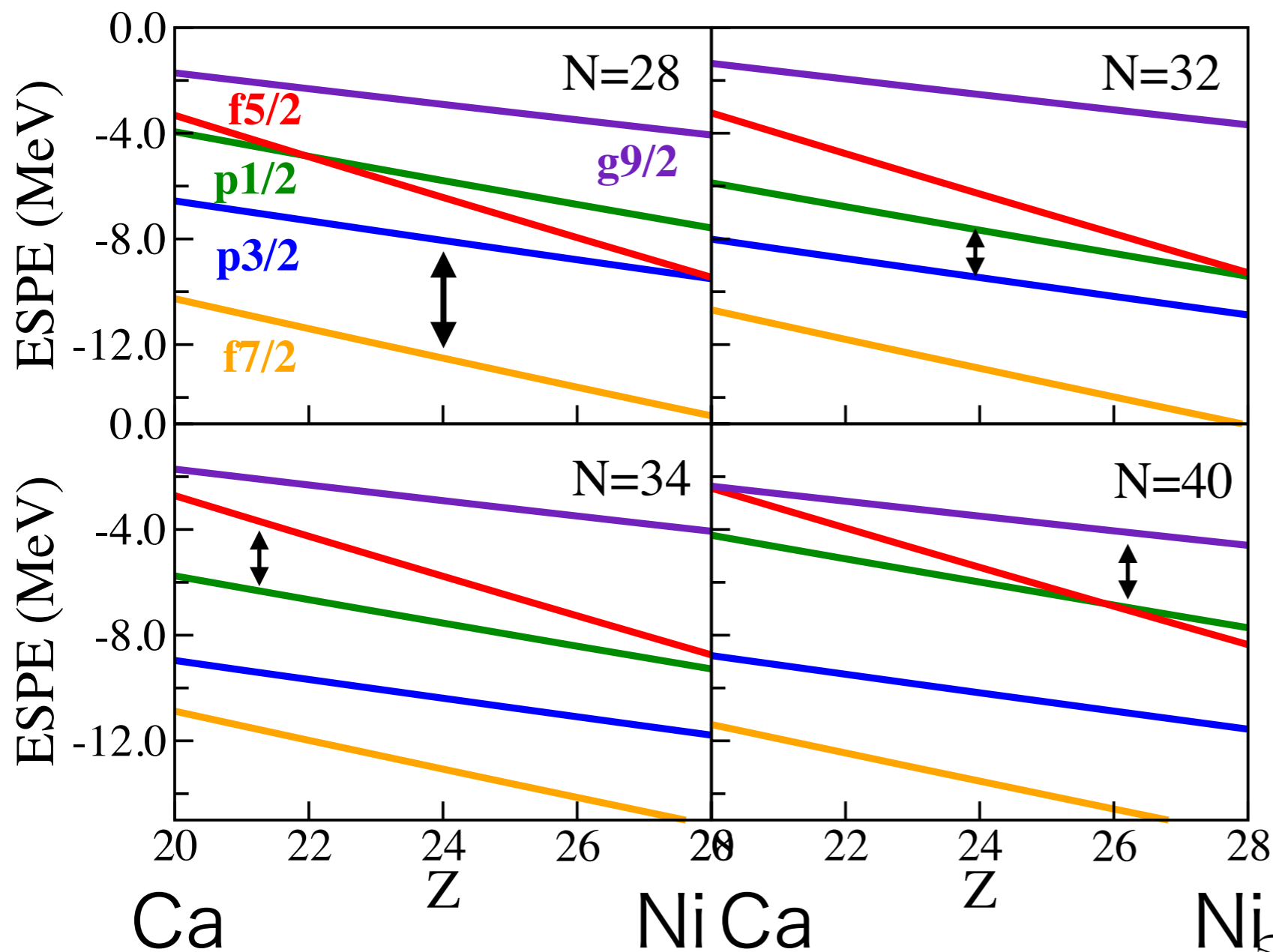
3rd order perturbation

6380 TBMEs

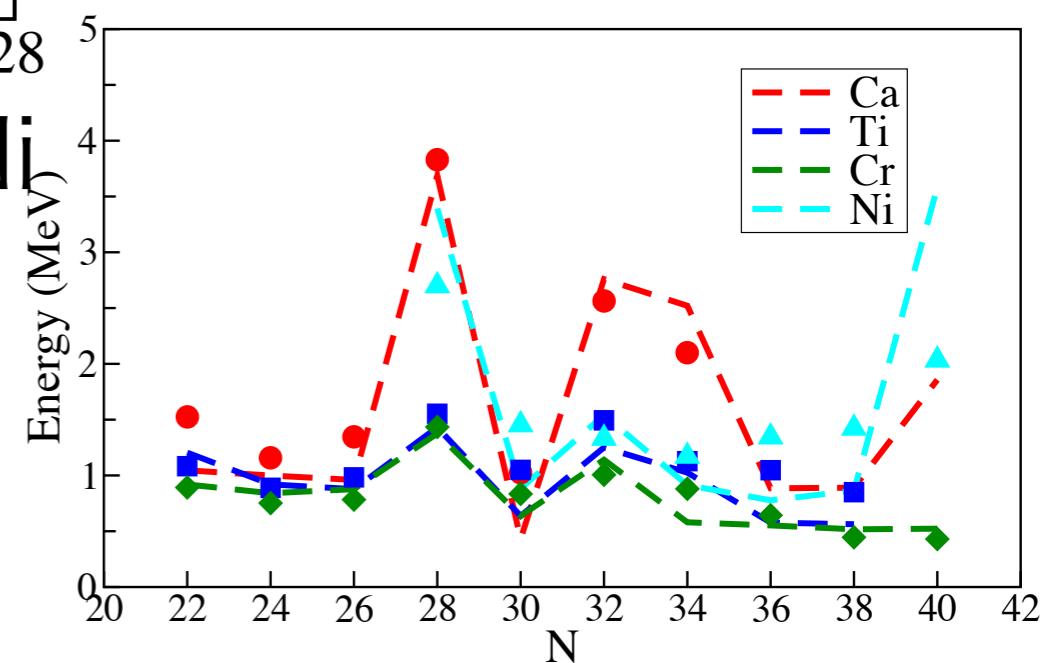
500 node\*h in K

Appearance and disappearance of  
 $N=28, 32, 34, 40$  magic numbers

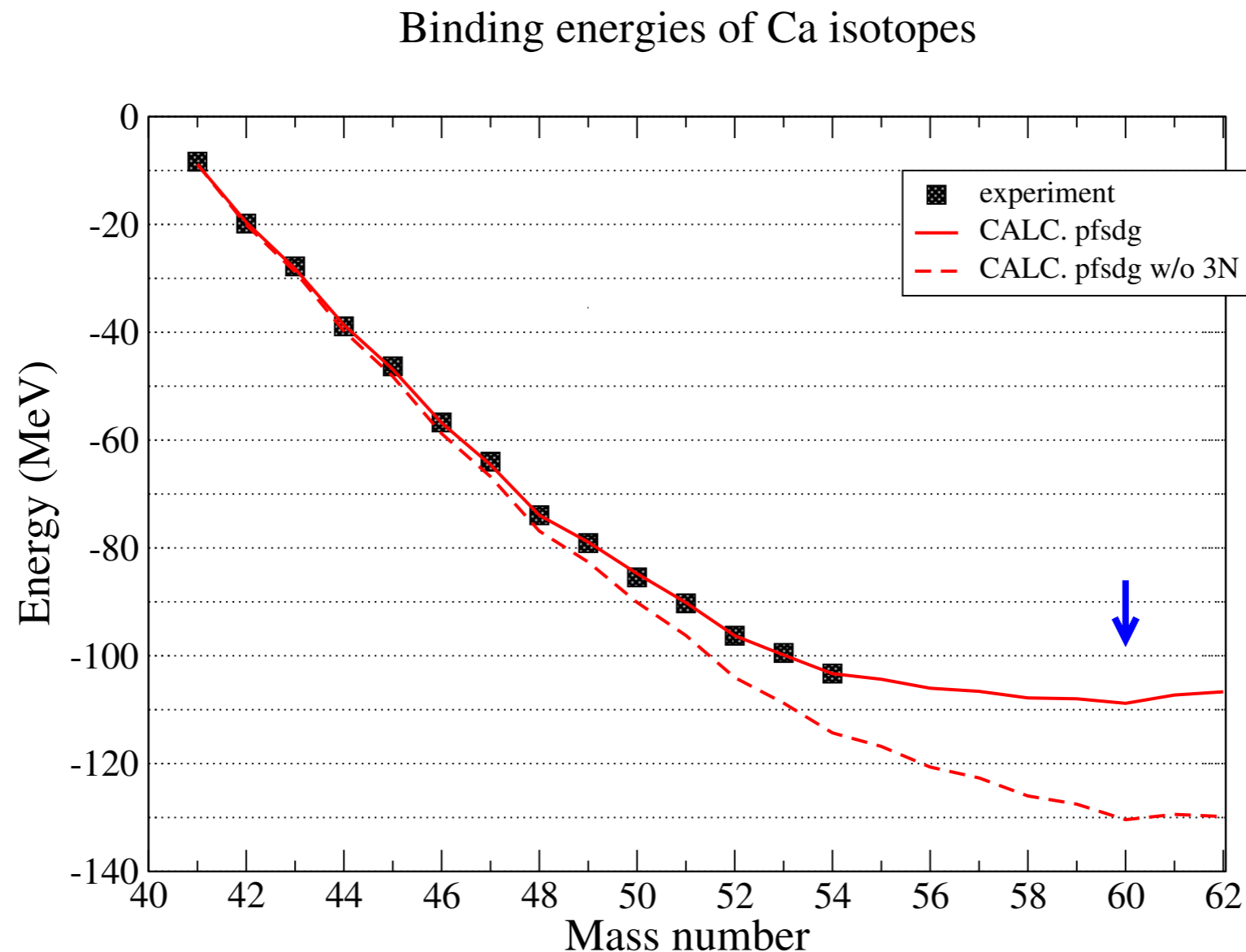
# Evolution of single particle states



N=28, 32 gap looks stable  
comparing to N=34,40 gap



# Binding energy and 3N force (preliminary)



3NF is important for binding energy, in the same way as in the case of sd<sub>2</sub>pf-shell

pf+sdg-shell の計算は、ある程度小規模の計算により、うまくいきそうなことが分かった段階。

## 今後と課題

### 1. 三体力の扱い

- ・ Fujita-Miyazawa type だけでなく、カイラル摂動論に基づいた他の方法
- ・ 摂動論の基底の変換(Hartree-Fock basis)

### 2. EKK 計算を精密化

- ・ テストケースでは、13 major-shell だったが、収束にはもう少し必要 → ポスト京に向けて、コードをさらに高速化



# Summary and conclusion

- MBPT is the theory to construct the effective Hamiltonian starting from nuclear force.
- **EKK method** is introduced to derive the effective interaction for the shell model which is applicable to **multi-shell** system.
- As an application of EKK method, the physics in the “**island of inversion**” is discussed in K-computer.
- As a future project, physics in pf+sdg-shell is discussed with preliminary results.
- **EKK** and **3N** combination is the powerful tool to explore the wide area of the nuclear chart

# Collaborators

- Takaharu Otsuka (Univ. Tokyo)
- Noritaka Shimizu (CNS)
- Sota Yoshida (Univ. Tokyo)
- Takayuki Miyagi (Univ. Tokyo)
- Kazuo Takayanagi (Sofia Univ.)
- Toshio Suzuki (Nihon Univ.)
- Morten Hjorth-Jensen (Oslo Univ.)