

核力から出発した有効相互作用 の構築とその応用

角田 直文

Center for Nuclear Study, Univ. Tokyo

tsunoda@cns.s.u-tokyo.ac.jp

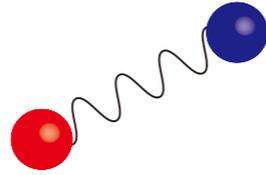
HPCI Field5: The origin of matter and the universe

QCD



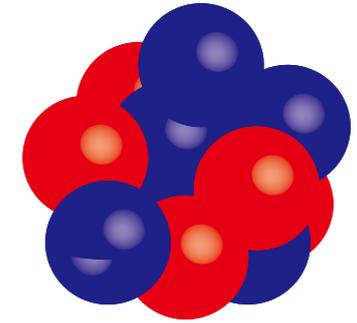
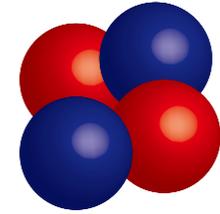
Lattice QCD
Effective Field Theory

Nuclear force

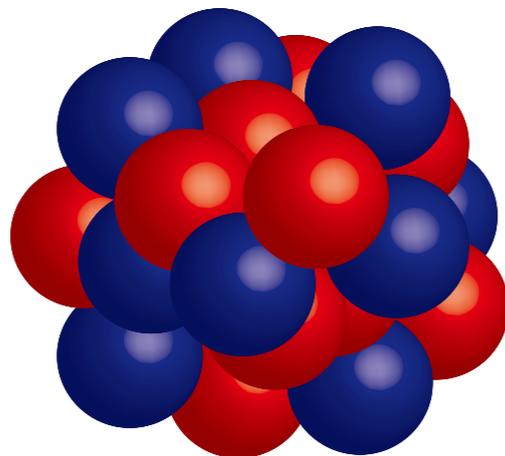


Few body techniques
No core shell model
and many others...

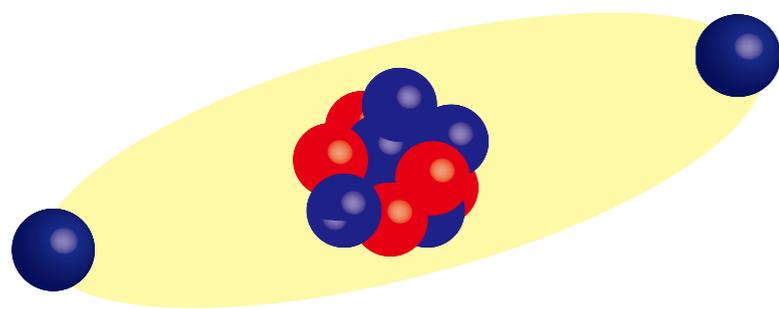
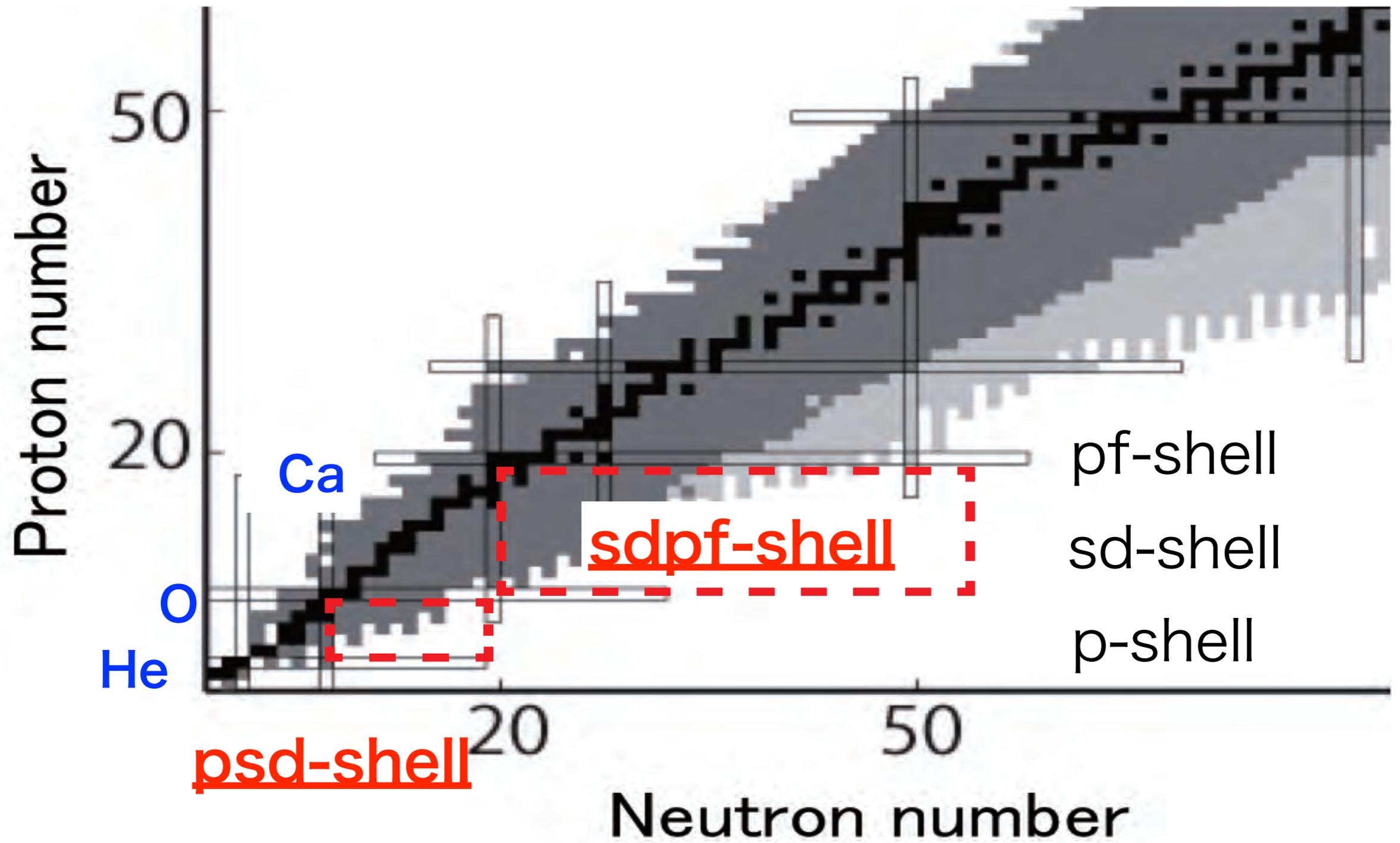
Light nuclei $\sim A \approx 10-20$



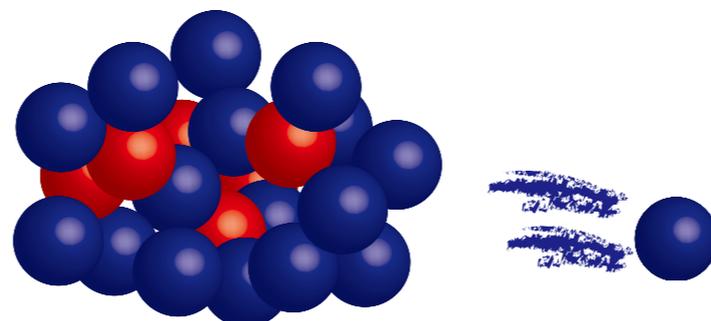
Medium mass nuclei $\sim A \approx 20-100$



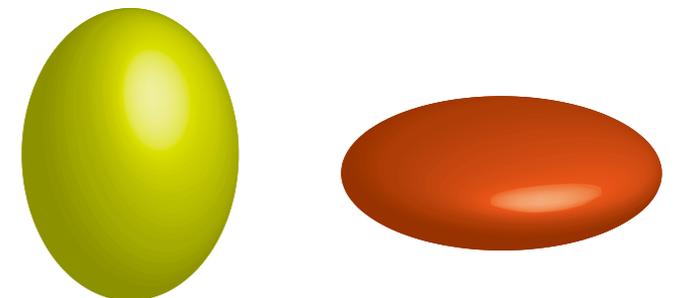
shell model with core
via the **effective interaction**
derived from **nuclear force**



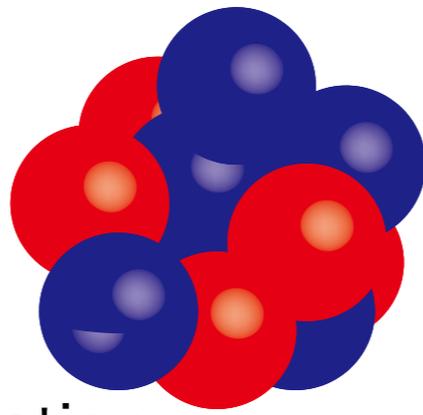
Neutron Halo



Dripline



Deformation



non-relativistic schrodinger equation

$$H|\Psi\rangle = E|\Psi\rangle \quad |\Psi\rangle : \text{many-body wave function}$$

Solve the Schrodinger equation with :

1. Nuclear force
2. Degrees of freedom which is appropriate for the problem we like to attack

What we have to overcome are:

- A. Nuclear force has strong short range repulsion (singularity) \rightarrow Vlowk interaction (not today)
- B. Unless we take whole Hilbert space, we need to fix **effective interaction** for selected model space

Nuclear force and Nuclear shell model

Shell model Hamiltonian

$$H = \sum_i \epsilon_i a_i^\dagger a_i + \sum_{ijkl} V_{ij,kl} a_i^\dagger a_j^\dagger a_l a_k.$$

input

ϵ_i : single particle energies

$V_{ij,kl}$: two-body matrix elements

diagonalization

output

Nuclear properties

Binding energy, energy spectrum, transition probability , etc...



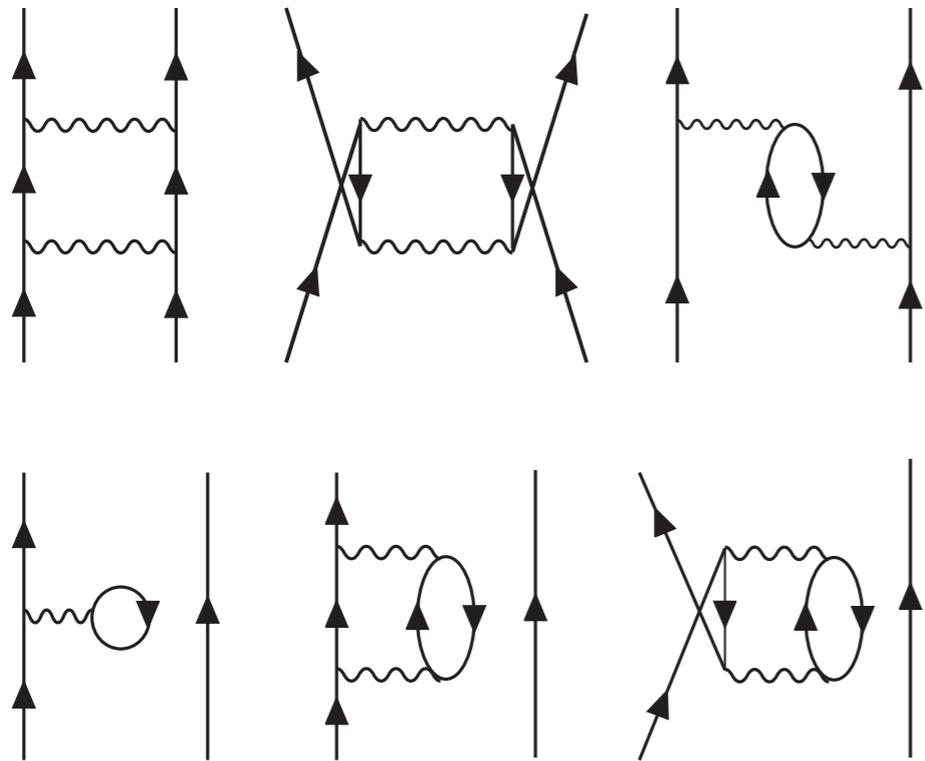
Derive Shell model Hamiltonian based on Nuclear force and many-body theories

Q-box expansion

Q-box is the ingredient of effective interaction and approximated by perturbation theory

$$\hat{Q}(E) = PVP + PVQ \frac{1}{E - QHQ} QVP$$

$$= PVP + PVQ \frac{1}{E - QH_0Q} QVP + PVQ \frac{1}{E - QH_0Q} QVQ \frac{1}{E - QH_0Q} QVP + \dots$$



Diagrams appearing in 2nd order Q-box

beyond perturbation contribution

$$V_{\text{eff}}^{(n)} = \hat{Q}(\epsilon_0) + \sum_{k=1}^{\infty} \hat{Q}_k(\epsilon_0) \{V_{\text{eff}}^{(n-1)}\}^k$$

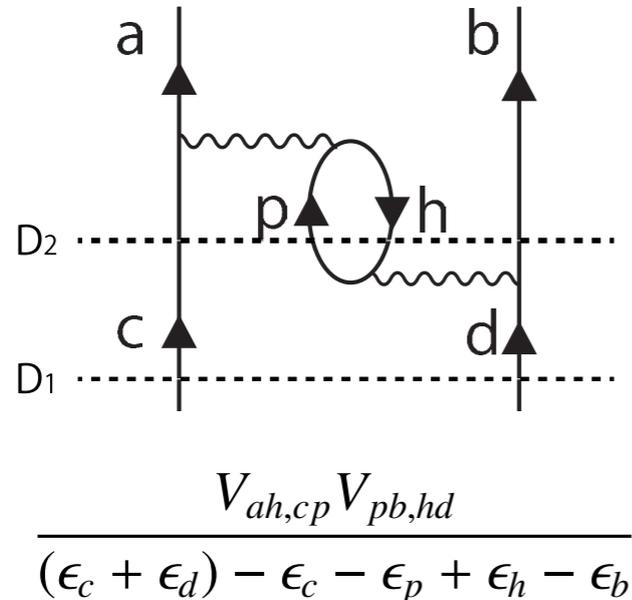
$$\hat{Q}_k(E) = \frac{1}{k!} \frac{d^k \hat{Q}(E)}{dE^k}$$

Divergent problem of Q-box in non-degenerate model space

(A) Folded diagram theory requires assumption that the model space is **degenerate**

(B) Naive perturbation theory leads a **divergence** in non-degenerate model space

Example



→ Energy denominator is zero
when $\epsilon_d - \epsilon_b = \epsilon_p - \epsilon_h$

We need a theory which satisfies

(a) The assumption of degenerate model space is **removed**

(b) **Avoid** the divergence appearing in Q-box diagrams

→ **EKK method as a re-summation scheme of KK method**

Extended KK method as a re-summation of the perturbative series

EKK method is derived with the following re-interpretation of the Hamiltonian

$$H = H'_0 + V'$$

$$= \begin{pmatrix} E & 0 \\ 0 & QH_0Q \end{pmatrix} + \begin{pmatrix} P\tilde{H}P & PVQ \\ QVP & QVQ \end{pmatrix},$$

New parameter E (arbitrary parameter)

Change PH₀P part of the unperturbed Hamiltonian

KK method

$$\hat{Q}(E) = PVP + PVQ \frac{1}{E - QHQ} QVP$$

$$V_{\text{eff}}^{(n)} = \hat{Q}(\epsilon_0) + \sum_{k=1}^{\infty} \hat{Q}_k(\epsilon_0) \{V_{\text{eff}}^{(n-1)}\}^k.$$

EKK method

$$H_{\text{BH}}(E) = PHP + PVQ \frac{1}{E - QHQ} QVP.$$

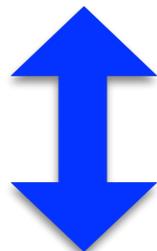
$$\tilde{H}_{\text{eff}}^{(n)} = \tilde{H}_{\text{BH}}(E) + \sum_{k=1}^{\infty} \hat{Q}_k(E) \{\tilde{H}_{\text{eff}}^{(n-1)}\}^k.$$

- One can take E so as to avoid the divergence !
- Final result does not depends on E.

Extended KK method as an analogy of Taylor series

KK method

$$V_{\text{eff}}^{(n)} = \hat{Q}(\epsilon_0) + \sum_{k=1}^{\infty} \hat{Q}_k(\epsilon_0) \{V_{\text{eff}}^{(n-1)}\}^k.$$



$$e^x = 1 + \sum_{k=1}^{\infty} \frac{1}{k!} x^k$$

Taylor expansion
around x=0

EKK method

$$\tilde{H}_{\text{eff}}^{(n)} = \tilde{H}_{\text{BH}}(E) + \sum_{k=1}^{\infty} \hat{Q}_k(E) \{\tilde{H}_{\text{eff}}^{(n-1)}\}^k.$$



$$e^x = e^E + \sum_{k=1}^{\infty} \frac{e^E}{k!} (x - E)^k$$

Taylor expansion
around x=E

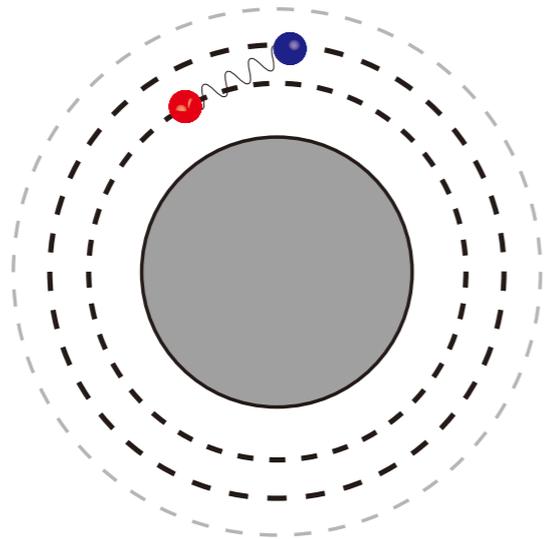
→ Result does **not** depend on E

Nuclear force and Nuclear shell model

Shell model Hamiltonian

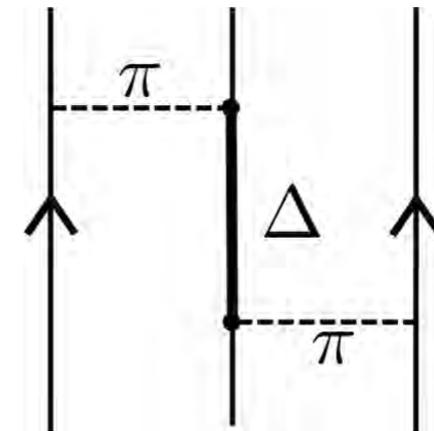
$$H = \sum_i \epsilon_i a_i^\dagger a_i + \sum_{ijkl} V_{ij,kl} a_i^\dagger a_j^\dagger a_l a_k.$$

Effective interaction in two-body space



+

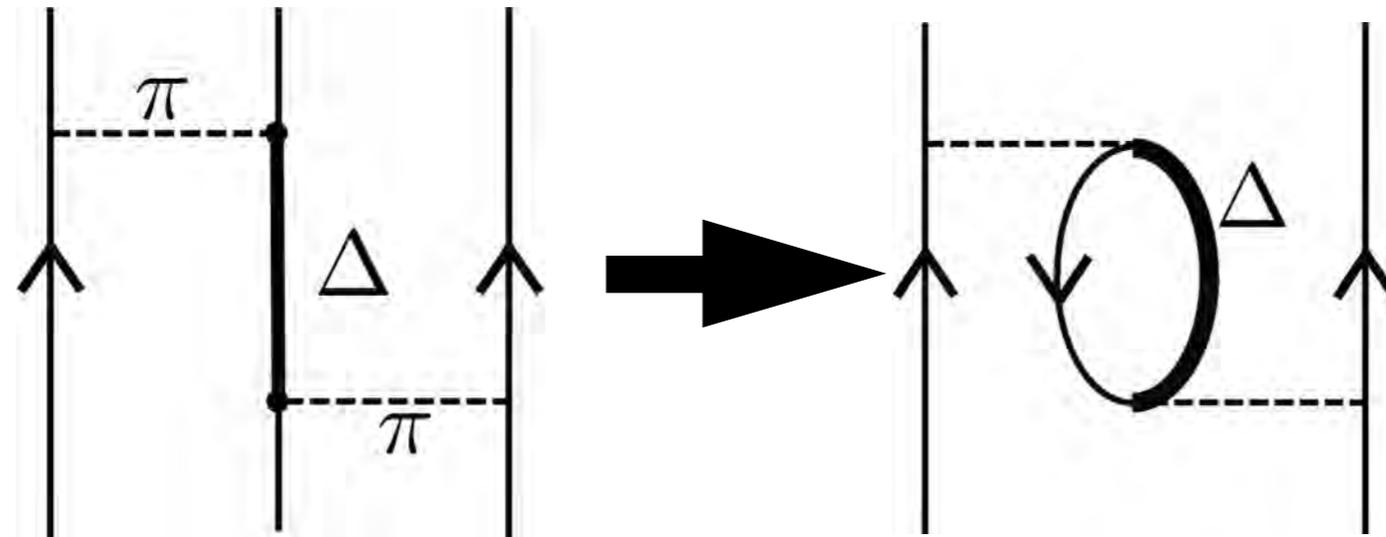
Three-body force



Reduce interaction to the model space
perturbatively
Many-body perturbation theory

Fujita-Miyazawa interaction

3N interaction



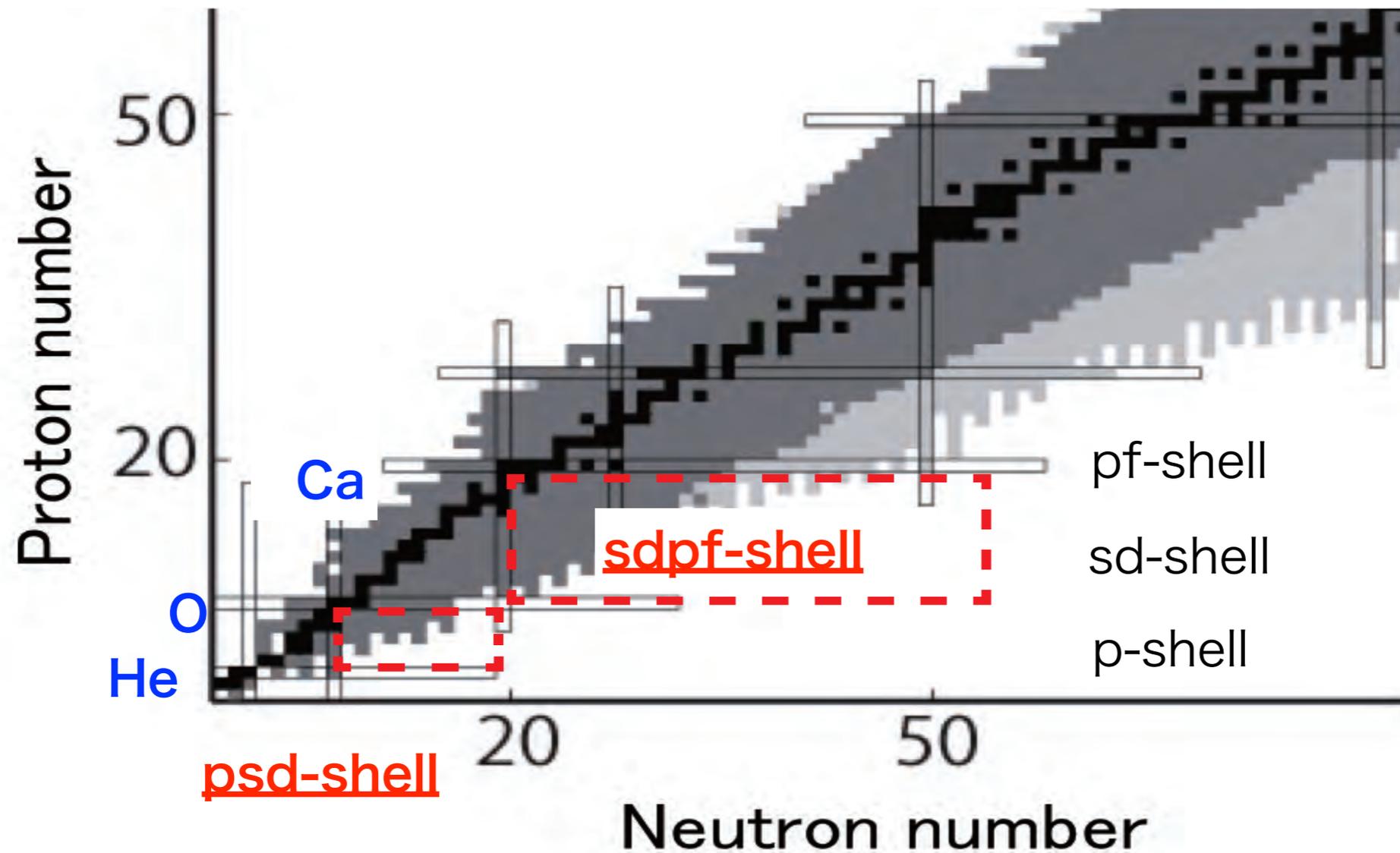
summation with hole

Fujita-Miyazawa type
3N interaction

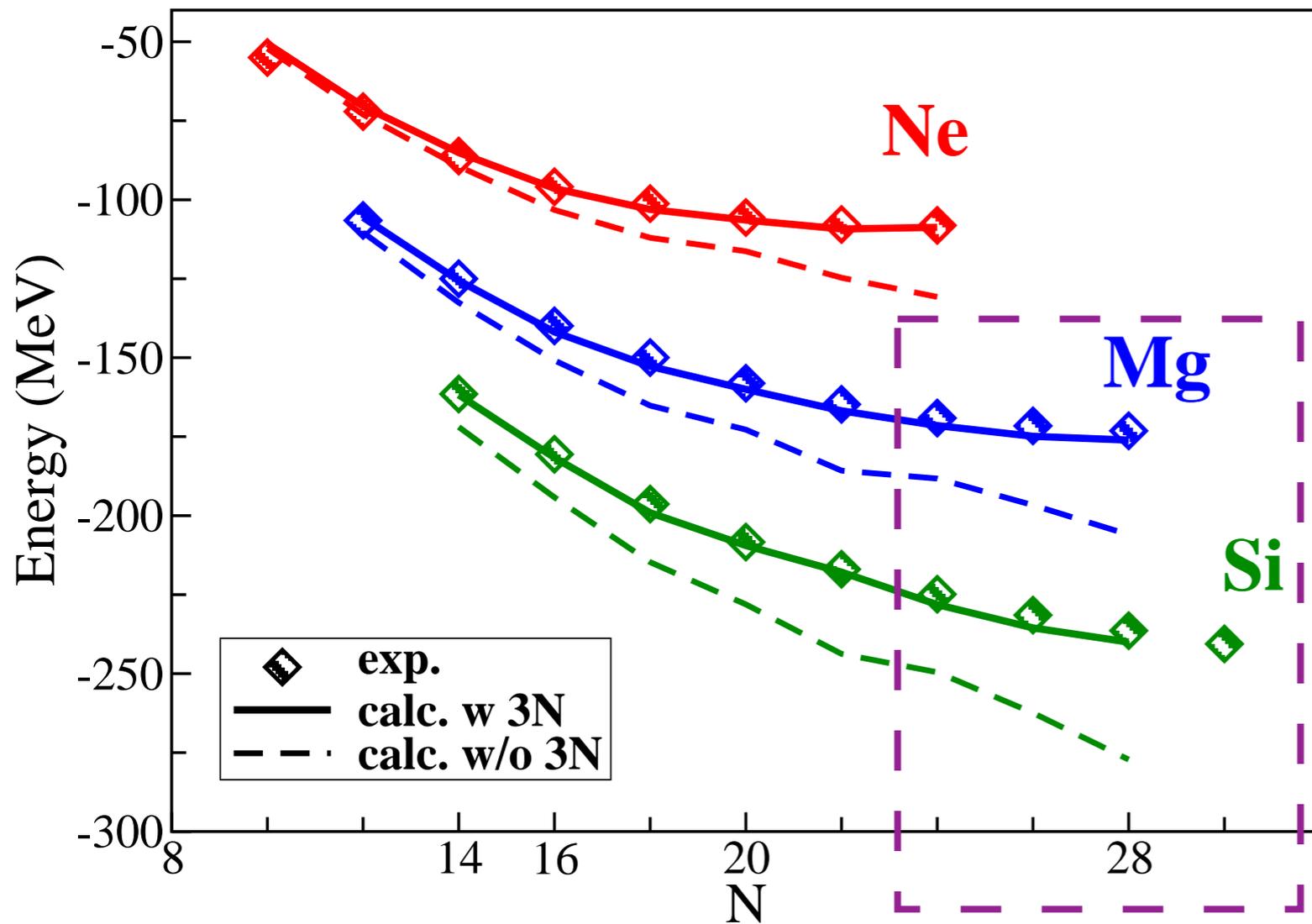
Effective
2N interaction

- Adding up effective 2N interaction derived from 3N interaction to EKK 2N effective interaction
- This is one of the lowest order interaction from 3N force and for higher order we are working on...

Application to sdpf-shell



Ground state energies



Typical dimension to diagonalize (4hw)

$${}^{32}\text{Ne} = 1.78 \times 10^8$$

$${}^{34}\text{Mg} = 2.26 \times 10^9$$

$${}^{36}\text{Si} = 8.10 \times 10^9$$

$${}^{38}\text{Si} = 4.53 \times 10^{11}$$

Monte Carlo Shell model

PC

workstation

super comp.

k-comp.

MCSM

10^6

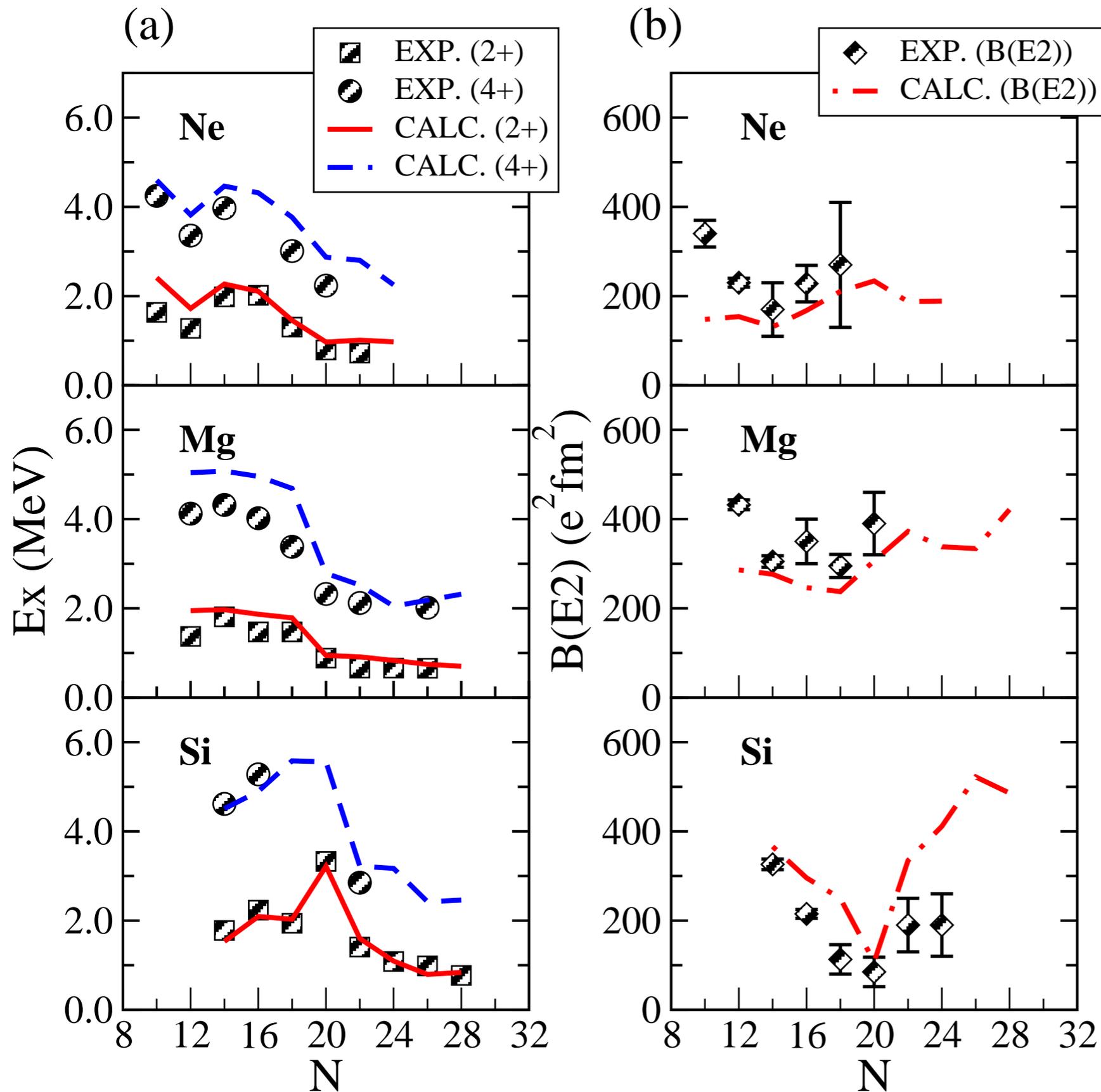
10^8

10^9

10^{11}

dim.

shell structure ~ island of inversion

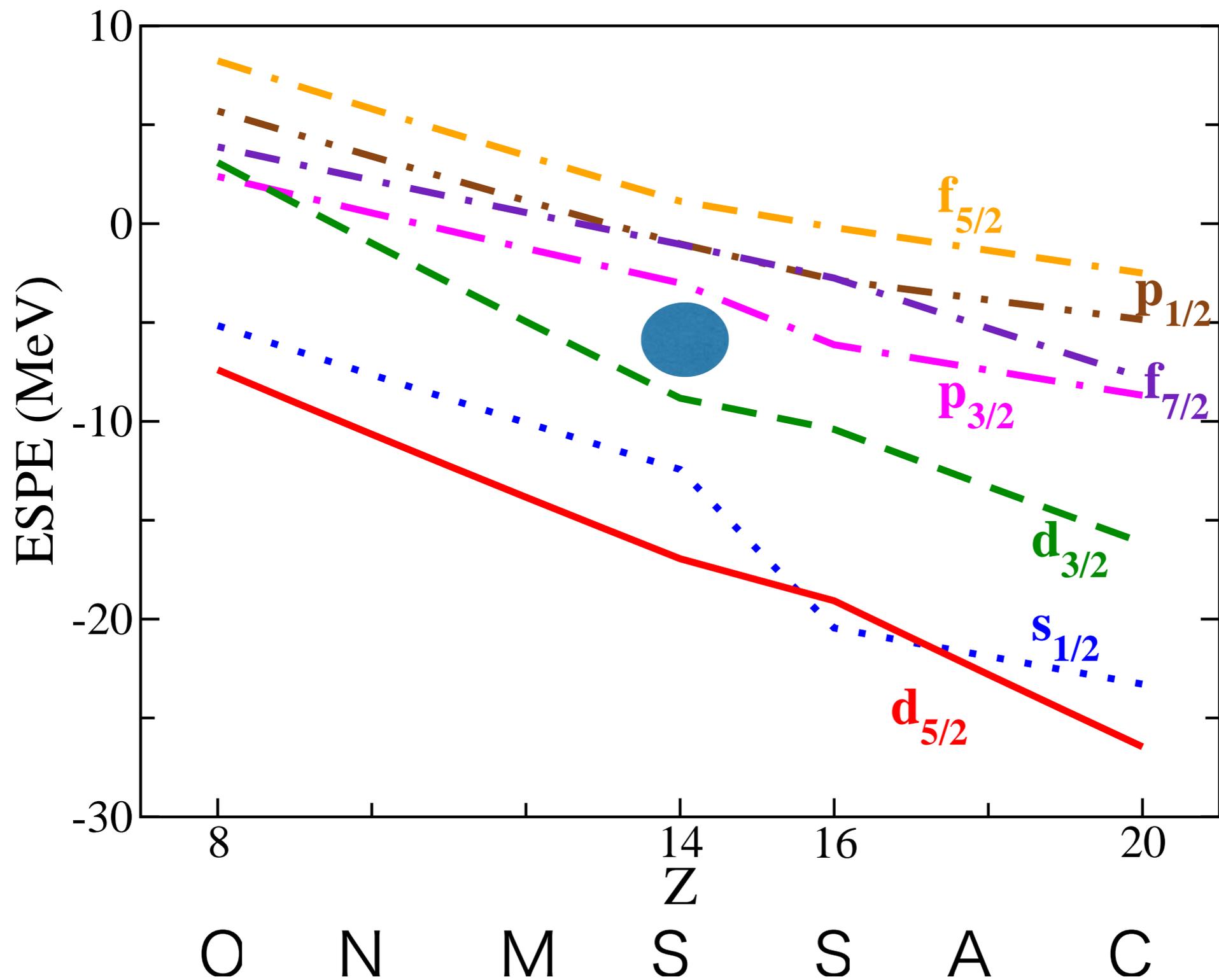


Ne, Mg: $N=20$
gap Vanishes

vs

Si: $N=20$ gap
appears

Effective single particle energy



Summary and conclusion

- To describe medium mass nuclei starting from nuclear force, we need some method to derive effective interaction.
- EKK method is developed to derive effective interaction for multi-shell, which has an energy parameter E that we can estimate the accuracy of the approximation via E -dependence of the final results.
- As the application of EKK method Ne, Mg, Si isotopes are discussed.
 - Physics in island of inversion is well described with EKK+3N framework

Collaborators

- Takaharu Otsuka (Univ. Tokyo)
- Noritaka Shimizu (CNS)
- Kazuo Takayanagi (Sofia Univ.)
- Toshio Suzuki (Nihon Univ.)
- Morten Hjorth-Jensen (Oslo Univ.)

e-mail: tsunoda@cns.s.u-tokyo.ac.jp