

Large N_c gauge theory and Chiral Random Matrix Theory

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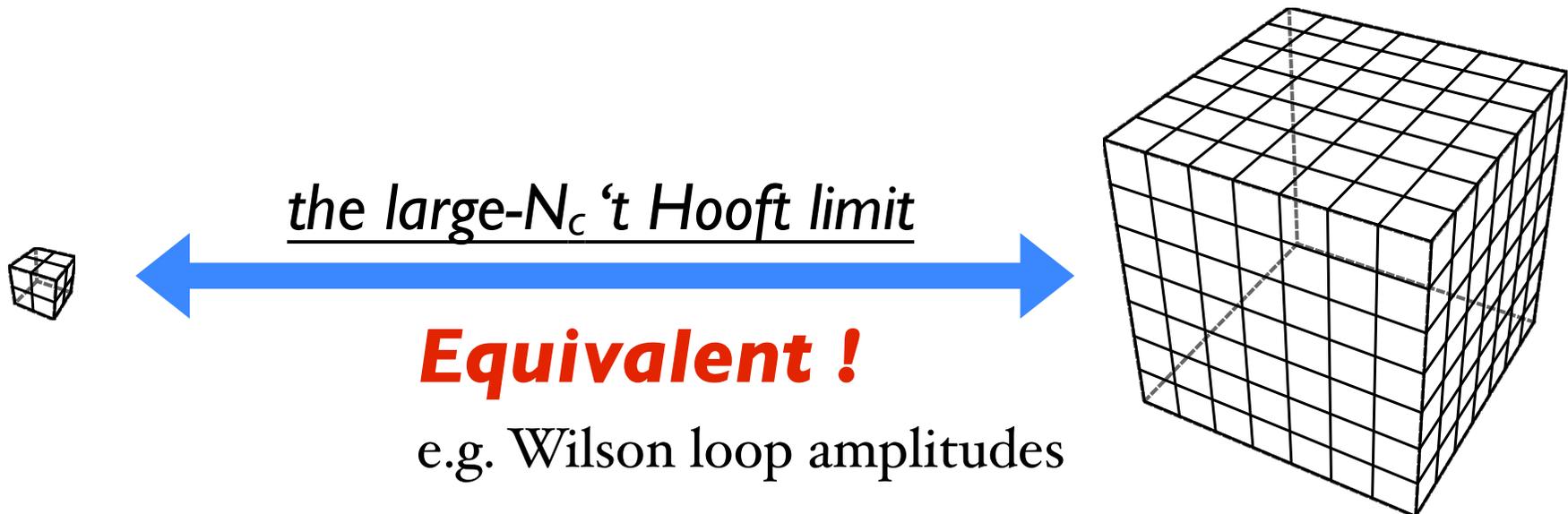
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Collaborators: Masanori Hanada, Norikazu Yamada (in progress)

In the large- N_c 't Hooft limit of QCD, the theory is greatly simplified (planarity) while sharing the same properties (asymptotic freedom, confinement, chiral symmetry breaking, ...) with QCD.

Consider $SU(N_c)$ gauge theory



This is so called the Large- N_c volume equivalence. (Eguchi-Kawai 82')
(Eguchi-Kawai equivalence)

Question: In $SU(N_c)$ gauge theory with $N_f=2$ adj. fermions at infinite V , the chiral symmetry is broken or not?

❖ $SU(2)$ with $N_f=2$ adjoint fermions - candidate of the Minimal Walking Technicolor model

Instead of considering the theory at infinite volume, in the 't Hooft large- N_c limit (thanks to Large- N_c volume equivalence) we study



$SU(N_c)$ gauge theory with $N_f=2$ adj. fermions at small V
($=2^4$ in our simulation)

The breakdown of the chiral symmetry can be detected by using the Chiral Random Matrix Theory (chRMT).

● Large N_c Volume Independence (80's)

In 1982, Eguchi and Kawai proposed that large volume $SU(N_c)$ gauge theory is identical to a single-site reduced matrix model in the large N_c limit if certain conditions are satisfied. The most crucial condition is the unbroken center symmetry.

However, single-site Eguchi-Kawai volume reduction Fails $\text{tr } U_\mu \neq 0$
Bhanot, Heller & Neuberger 82

Prescriptions in early 80's

Quenched EK

Bhanot, Heller & Neuberger 82

Fails Brigoltz & Sharpe 2008

Twisted EK

Gonzales-Arroyo & Okawa 83

Teper & Vairinhos 2007

Azeyanagi et al 2008

Bietenholtz et al 2007

Revival (2007 ~) - center symmetry restoration

Partial reduction ($aL > 1\text{fm}$)

Adjoint EK (AEK)

Deformed EK

Modified TEK

Kiskis, Narayanan, Neuberger (2003)

Kovtun, Unsal, & Yaffe (2007)

Unsal & Yaffe (2008)

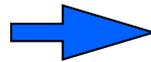
Gonzales-Arroyo & Okawa (2010)

● Adjoint Eguchi-Kawai model (AEK)

● Perturbative analysis

$N_f \geq 1$ fermions in adjoint representation with periodic boundary conditions

1-loop effective potential for Wilson lines is repulsive



stabilized center symmetry

Kovtun, Unsal & Yaffe 2007/2010

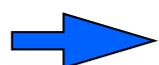
Azeyanagi, Hanada, Unsal, Yacoby. 2010

● Numerical simulations

single-site EK with $N_f=1,2$ adj. fermions

$(Z_{N_c})^4$ center symmetry is unbroken even at weak coupling & heavy fermions

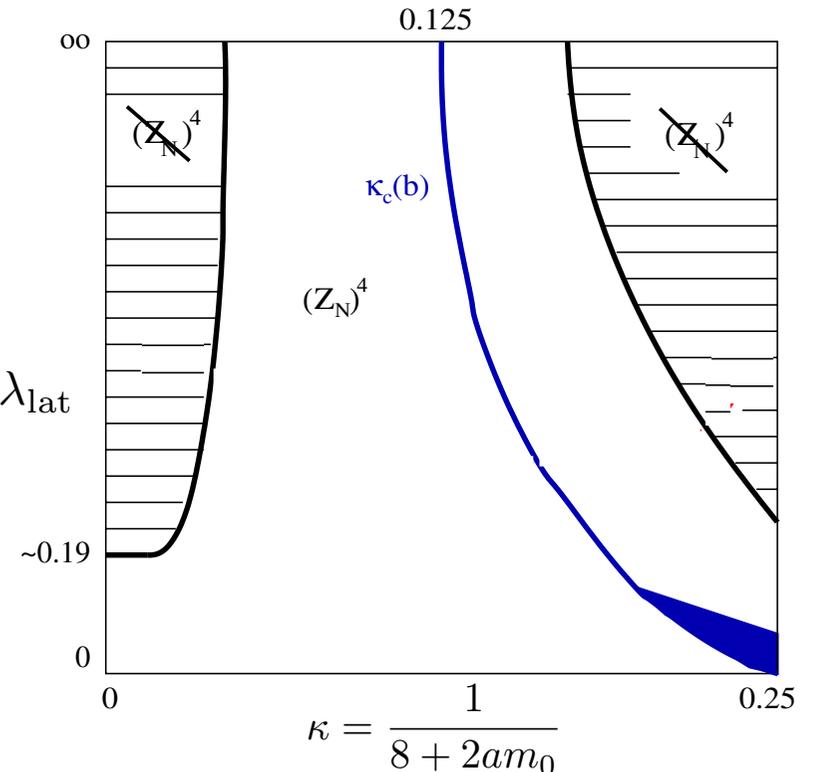
$$m_0 \sim \mathcal{O}(1/a) \quad b = 1/\lambda_{\text{lat}}$$



non-perturbative study of pure $SU(N_c)$ gauge theory

Bringoltz, Sharpe, 2009

Bringoltz, Koren, Sharpe, 2011



● Chiral Random Matrix Theory (chRMT)

chRMT provides equivalent descriptions of the low-energy limit of QCD-like theories if the **chiral symmetry is spontaneously broken.**

$$Z = \int DW \prod_{f=1}^{N_f} \det(\mathcal{D} + m_f) e^{-\frac{N\beta}{4} \text{Tr}(W^\dagger W)}$$

$$\mathcal{D} = \begin{pmatrix} 0 & iW \\ iW^\dagger & 0 \end{pmatrix} \quad \begin{array}{l} \text{anti-} \\ \text{Hermitian} \end{array}$$

W is $N \times N$ complex matrix

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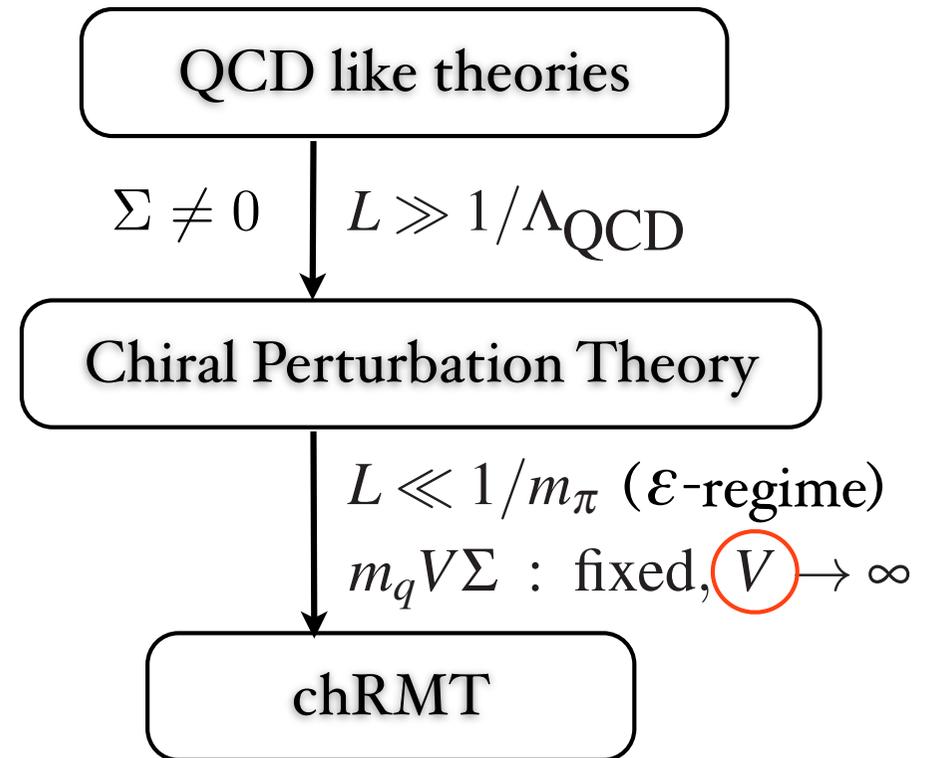
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spectral density of low-lying Dirac eigenvalues

$$\rho(z), \quad z = \lambda V \Sigma \sim \mathcal{O}(1)$$



- Order parameter of the spontaneous symmetry breaking

$$\Sigma = |\langle \bar{\psi} \psi \rangle| = \lim_{\epsilon \rightarrow 0} \lim_{m \rightarrow 0} \lim_{V \rightarrow \infty} \frac{\pi \rho(\epsilon)}{V} \quad \text{Banks-Casher relation}$$

● The chRMT & large- N_c gauge theories

In large- N_c gauge theories, colors are also relevant degrees of freedom. Thus,

$$V(N_c)^\alpha \longleftrightarrow N$$

To compare with chRMT

'chRMT limit': $m_q V(N_c)^\alpha \tilde{\Sigma}$: fixed, $N_c \rightarrow \infty$

spectral density: $\rho(z)$, $z = \lambda V(N_c)^\alpha \tilde{\Sigma} \sim \mathcal{O}(1)$

- **Order parameter of the spontaneous symmetry breaking**

$$\tilde{\Sigma} = |\langle \bar{\psi} \psi \rangle| = \lim_{\epsilon \rightarrow 0} \lim_{m \rightarrow 0} \lim_{N_c \rightarrow \infty} \frac{\pi \rho(\epsilon)}{V(N_c)^\alpha}$$

● The chRMT limit vs 't Hooft limit

'chRMT limit': $m_q V (N_c)^{\alpha \tilde{\Sigma}}$: fixed, $N_c \rightarrow \infty$

Different!

't Hooft limit': m_q, V : fixed, $N_c \rightarrow \infty$

The EK volume equivalence does not hold in the chRMT limit.

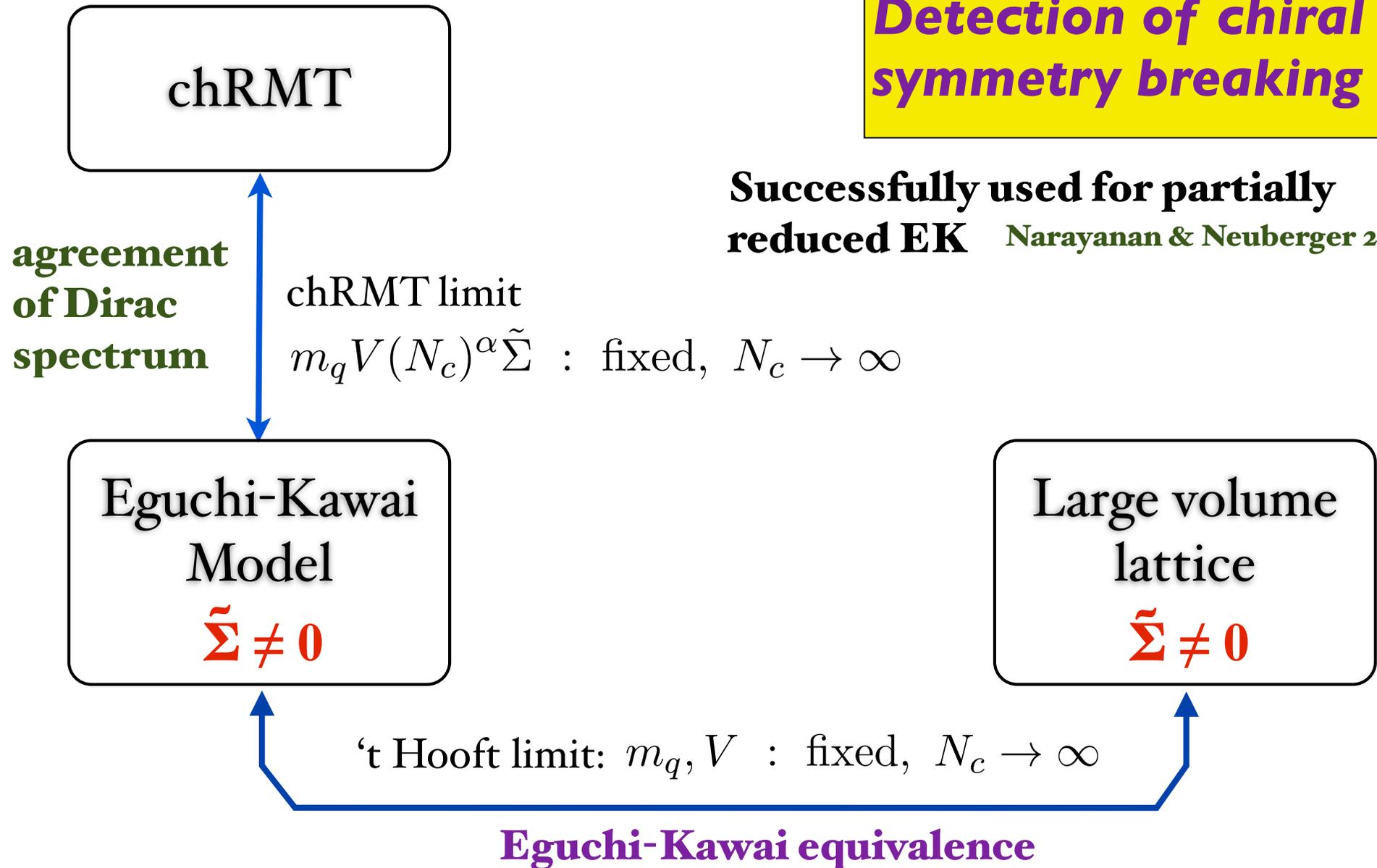
Can we still study the chiral symmetry breaking using the EK volume equivalence and chRMT in large- N_c gauge theories?

Yes!

● **EK model, large V LGT and chRMT**

Detection of chiral symmetry breaking

Successfully used for partially reduced EK Narayanan & Neuberger 2006



To establish the method, we numerically study *the quenched approximation*, where *the answer is known*.

In the large- N_c 't Hooft limit,

SU(N_c) EK model with $N_f=2$ heavy adjoint fermions
($m_0 \sim 1/a$) on a 2^4 lattice

**Center symmetry is unbroken and
thus the EK equivalence holds.**

Adjoint fermions are not dynamical.

pure SU(N_c) lattice gauge theory

Chiral symmetry is spontaneously broken?

● Action and Symmetries

● Action for reduced model on a 2^4 lattice

$$Z_{reduced} = \int DU D\bar{\psi} D\psi \exp(S_g + S_f), \quad n \in 2^4 \quad \text{Periodic B.C.}$$

$$S_g = 2N^2 b \sum_n \sum_{\mu < \nu} \left(1 - \frac{1}{N_c} \text{ReTr} P_{\mu\nu}(n) \right), \quad b = \frac{1}{g^2 N_c}$$

$$S_f = \sum_{j=1}^2 \sum_n \bar{\psi}_{n,j} \left(1 - \kappa \left[\sum_{\mu=1}^4 (1 - \gamma_\mu) U_{n,\mu}^{adj} \psi_{n+\mu,j} + (1 + \gamma_\mu) U_{n-\mu,\mu}^{\dagger,adj} \psi_{n-\mu,j} \right] \right).$$

$$\kappa = \frac{1}{2(m_0 a + 4)}, \quad m_0 \sim \mathcal{O}\left(\frac{1}{a}\right)$$

● Symmetries

$SU(N_c)$ **gauge symmetry:** $U_{n,\mu} \rightarrow \Omega_n U_{n,\mu} \Omega_{n+\mu}^\dagger, \quad \Omega_n \in SU(N_c)$

$Z_{N_c}^4$ **center symmetry:** $U_{n,\mu} \rightarrow e^{2\pi i n_\mu / N_c} U_{n,\mu}, \quad n_\mu \in Z_{N_c}$

● Simulation details

- **fermions in action**

quenched simulation ($\kappa=0$)

heavy Wilson fermions in adjoint rep. ($\kappa=0.09$)

↔ **unbroken center symmetry**

- **probe fermions**

massless overlap fermions in adjoint rep.

↔ **closer to RMT limit**

↔ **good chiral symmetry**

- **number of colors up to $N_c=16$**

- **bare coupling $b = 0.1, 0.2, 0.3, 0.4, 0.5$**
strong **weak**

- **Hybrid Monte Carlo (HMC) algorithm**

500 configs (10 trajectories) for each ensemble

200 trajectories for thermalization

● Tests of the Center Symmetry

If the $(\mathbf{Z}_{N_c})^4$ center symmetry is **unbroken**,

(1) Polyakov loop scatters radially in the vicinity of the origin.

$$P_\mu = \frac{1}{N_c} \text{tr} U_{n,\mu} U_{n+\mu,\mu}, \quad n \in 2^4$$

(2) Magnitude of the Polyakov loop goes to zero in the large N_c limit.

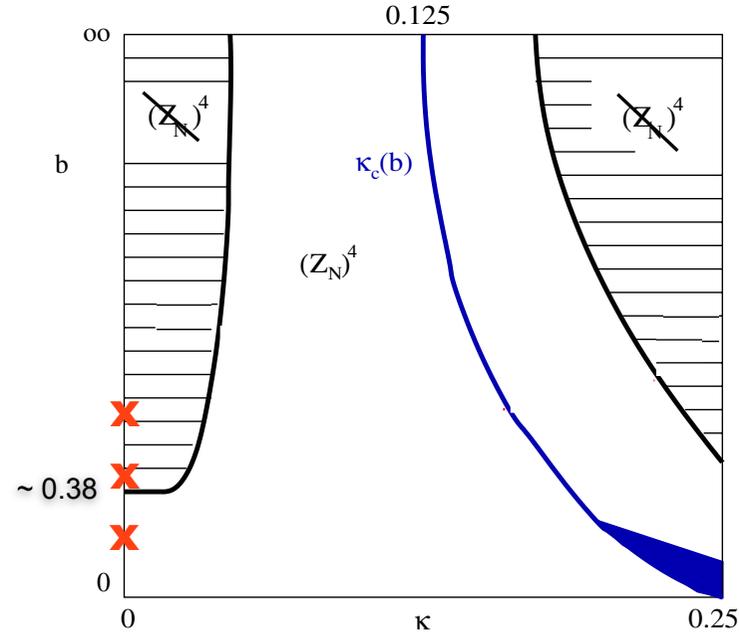
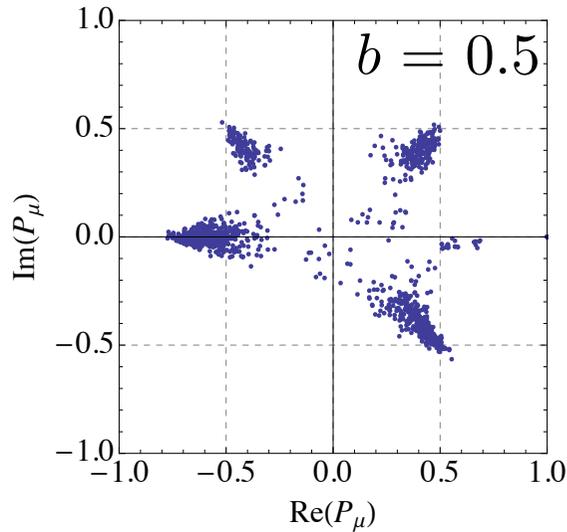
$$|P_\mu| \longrightarrow 0 \quad \text{as} \quad N_c \longrightarrow \infty$$

(3) Average plaquette value agrees with that of the large volume theory in the large N_c limit.

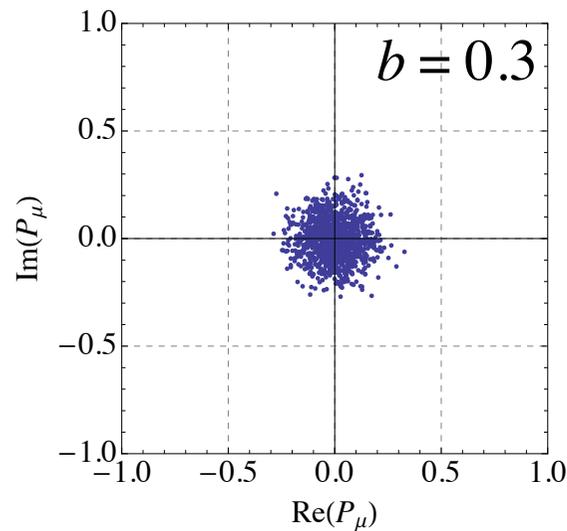
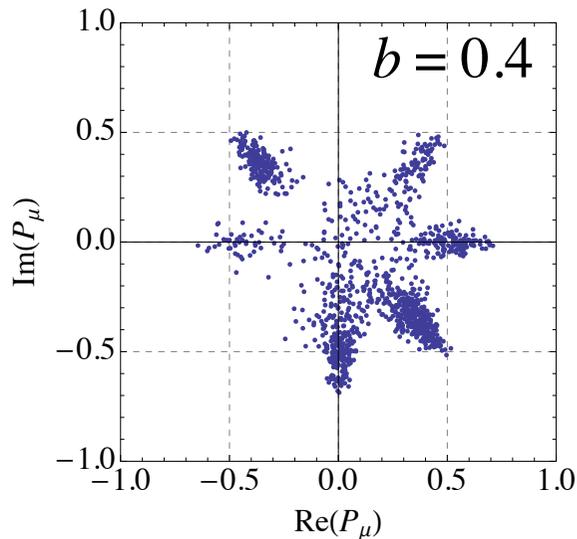
$$\langle \text{plaquette} \rangle = \frac{1}{N_c} \left\langle \sum_{\mu < \nu} \text{Tr} U_{n,\mu} U_{n+\mu,\nu} U_{n+\nu,\mu}^\dagger U_{n,\nu}^\dagger \right\rangle_n, \quad n \in 2^4$$

● Polyakov Loops - $N_f = 0$ (Quenched)

For $N_c = 8$,



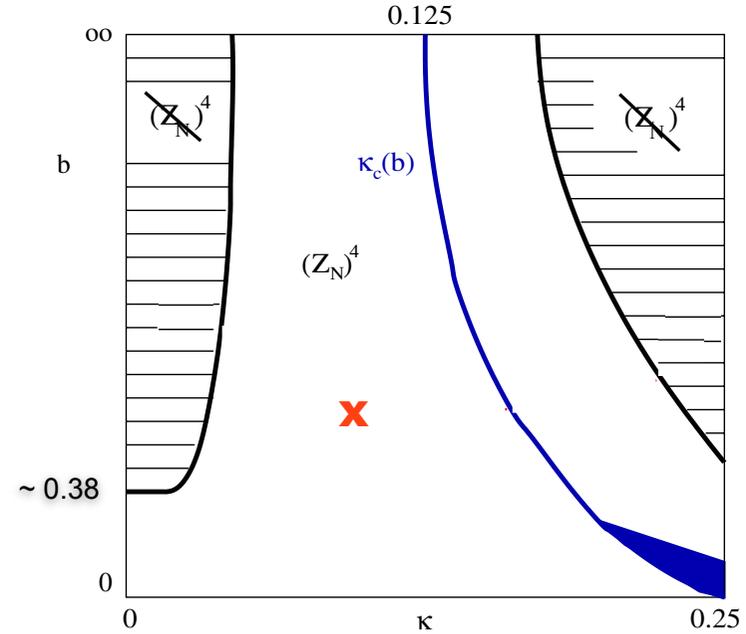
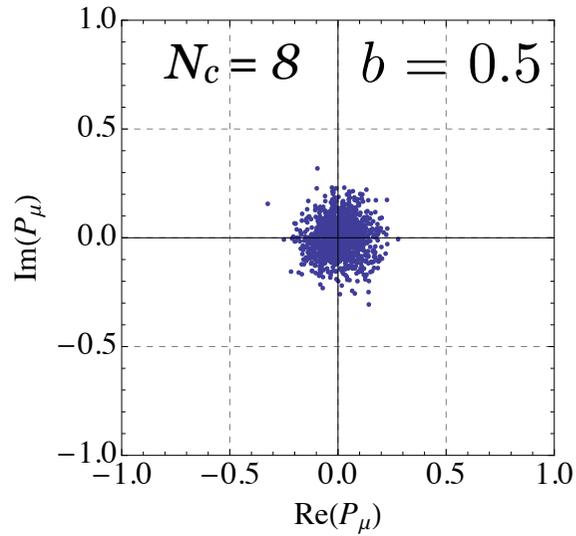
Center Symmetry is **broken** at weak coupling limit.



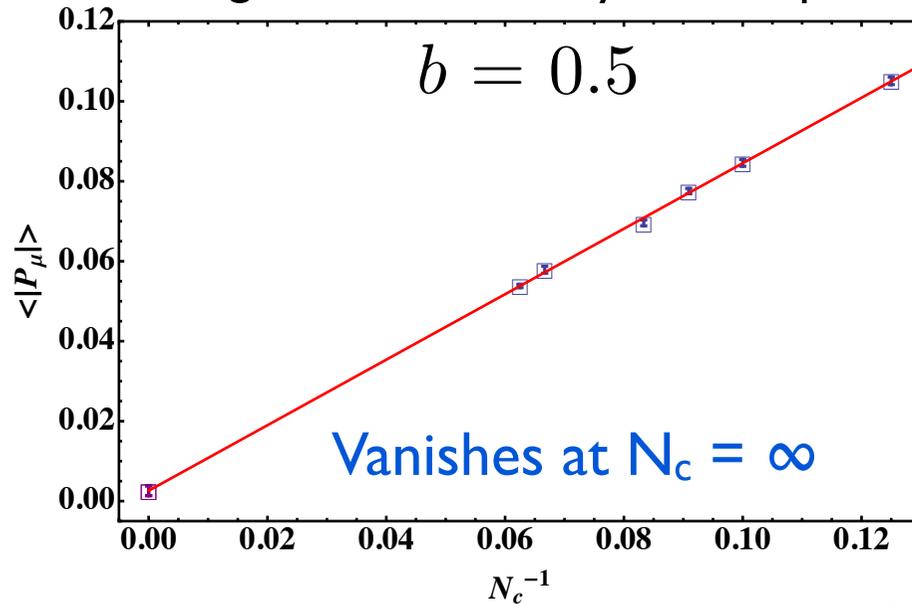
$$b = 1/(g^2 N_c)$$

Center Symmetry is **restored** in the strong coupling limit.

● Polyakov Loops - $N_f=2$ Heavy Adjoints Fermions



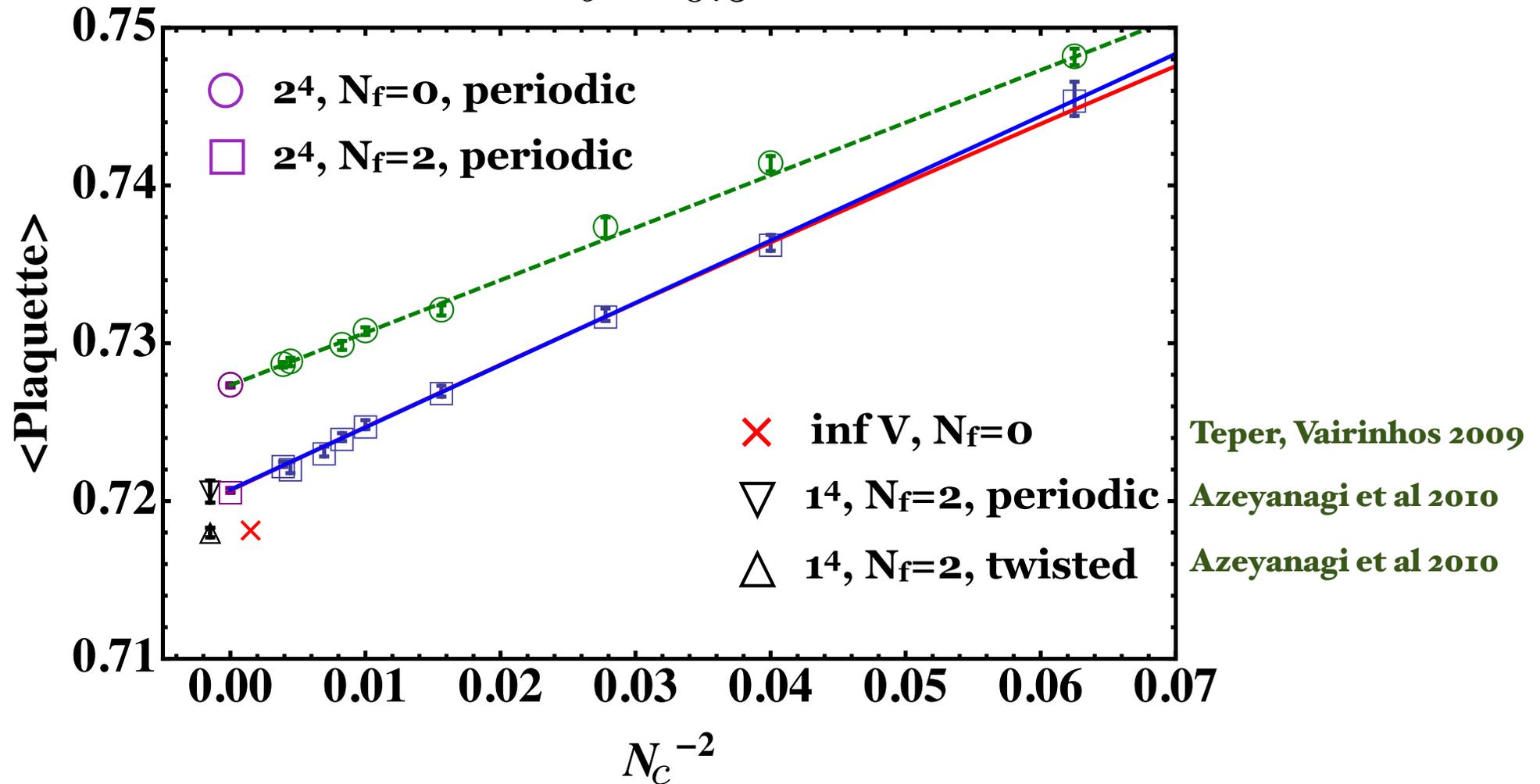
Magnitude of the Polyakov loops



Center Symmetry is **unbroken**
even in the weak coupling limit !

● Plaquette values (large N_c limit)

$$b = 0.5$$



Plaquette value for $N_f=2$ is consistent with that from large volume calculation, while $N_f=0$ value is far above all others.

● Spectrum of the Dirac operator

overlap-Dirac operator of fermions in adj. representation

$$H = \gamma_5 D = \gamma_5 \frac{1}{2} \left[1 + \gamma_5 \operatorname{sgn}[H_w(-m_0)] \right] \quad \gamma_5 H_w(-m_0)$$

Wilson-Dirac op.

fund. fermions \mathbf{N} \longrightarrow adj. fermions \mathbf{N}^2

calculate the low lying eigenvalues λ_i $z_i = \lambda_i V(N_c)^\alpha \tilde{\Sigma}$

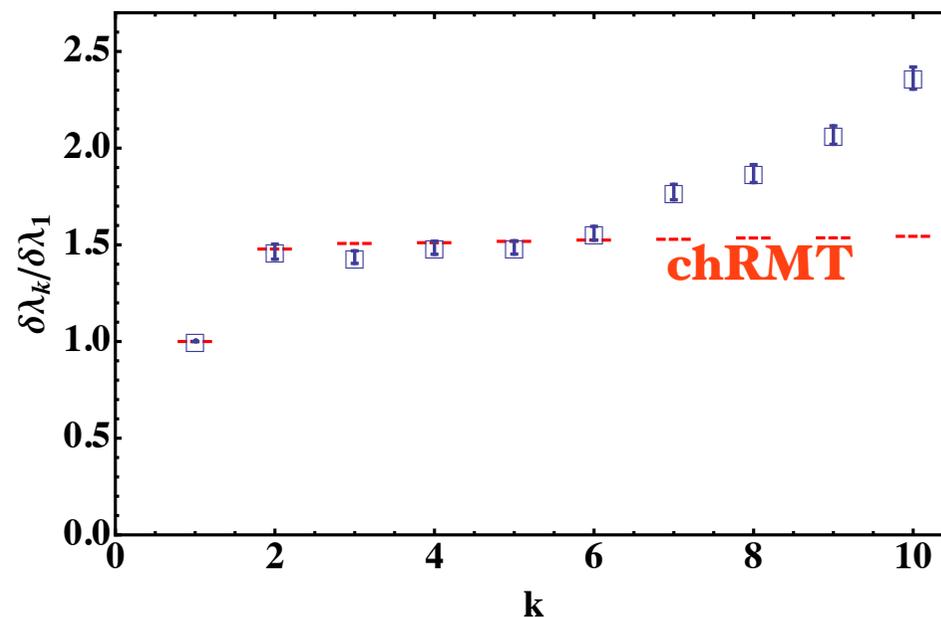
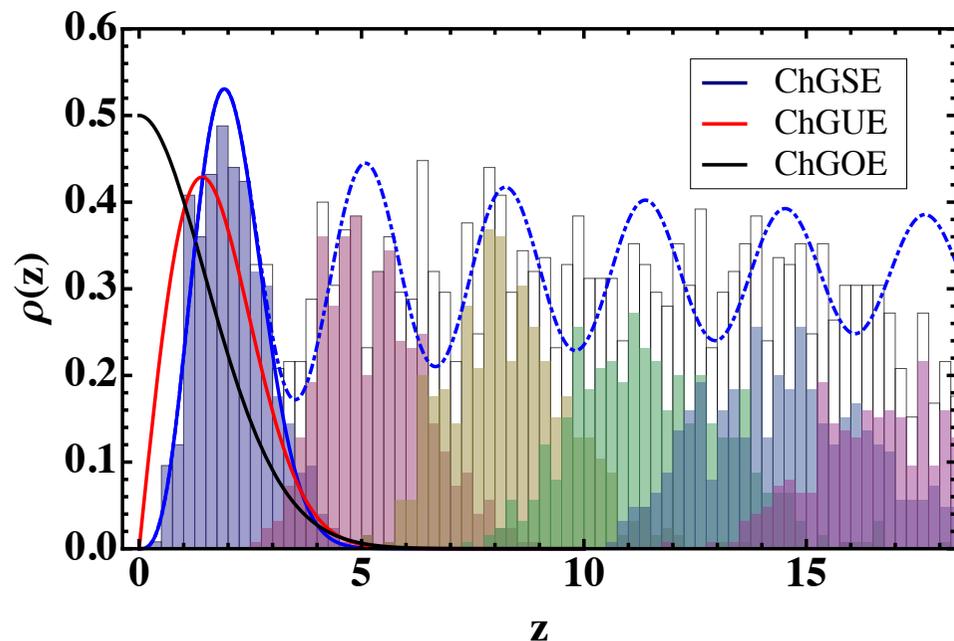
chRMT predictions:

distribution of the smallest eigenvalue for quenched theory

$$P_{\min}(z) = \begin{cases} \frac{2+z}{4} e^{-(z/2)-(z^2/8)} & \text{ChGOE} \\ \frac{z}{2} e^{-z^2/4} & \text{ChGUE} \\ \sqrt{\frac{\pi}{2}} z^{3/2} I_{3/2}(z) e^{-z^2/2} & \text{ChGSE} \\ & \text{QCD (adj)} \end{cases}$$

● Comparison with chRMT prediction

$$N_c = 16, b = 0.5, N_f = 2, V = 24$$



(reduced model)

chRMT prediction *perfectly* agrees with the numerical data -
chiral symmetry is broken.

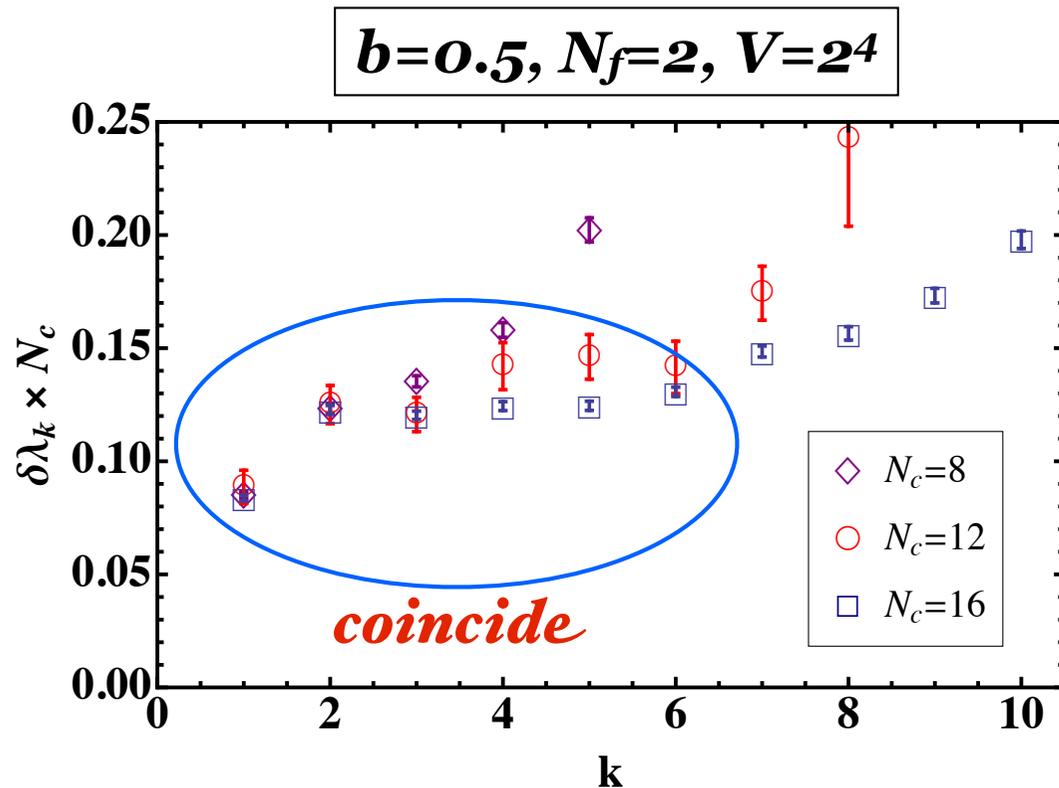


chiral symmetry is broken in the large- N_c quenched QCD

EK volume equivalence

● Scaling behavior of the Dirac eigenvalues

chRMT limit: mVN_c^α is fixed, while N_c (and/or V) $\rightarrow \infty$



λVN_c^α is fixed in the same way.



$\alpha = 1$

$1/N_c$ scaling behavior has also been found in a single-site simulations for two heavy adjoint fermions.

Narayanan & Hietanen 2012

● Conclusion and Future work

- Numerical study of QCD-like theories on a small lattice is possible by using EK volume equivalence and chiral RMT in the large N_c limit.

Quenched QCD - chiral symmetry is spontaneously broken

- Applying this method to

SU(N_c) gauge theory with $N_f = 2$ dynamical adjoint fermions

SU(2) two adjoint fermions - candidate of the Minimal Walking Technicolor model

- EK model is economical? Not really, because $V_{eff} \sim N_c$.

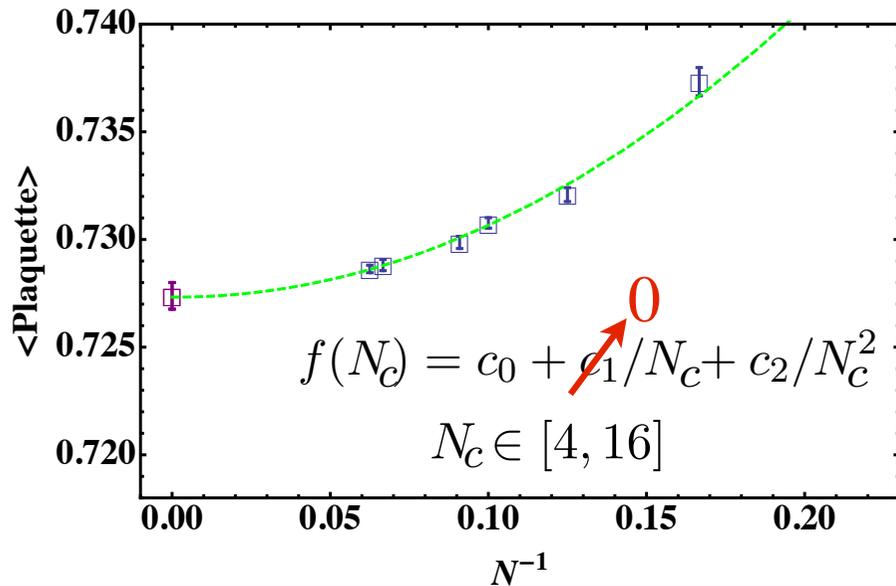
e.g. twisted boundary condition

$$1/N_c \longrightarrow 1/N_c^2$$

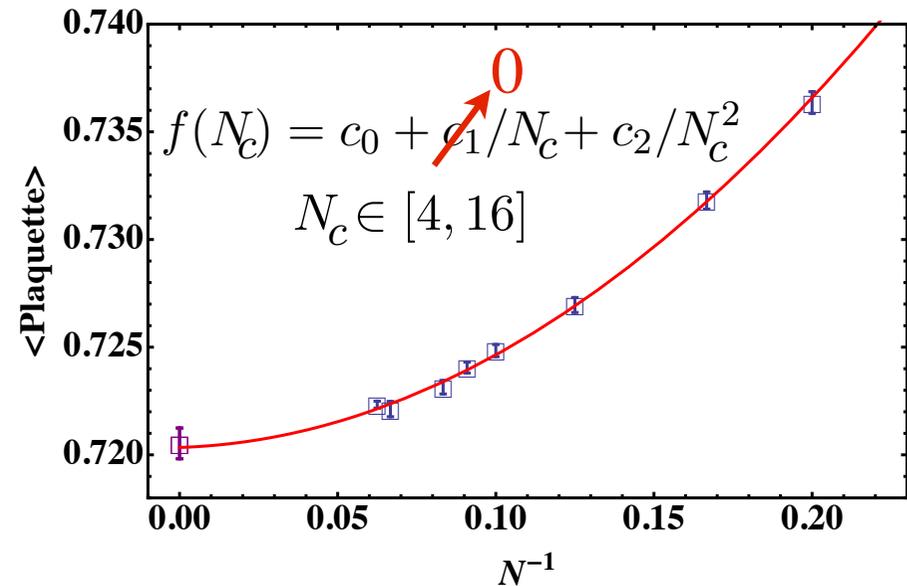
Thank you !

Backup slide 1 - Correction to plaquette values

$N_f = 0, b = 0.5$



$N_f = 2, b = 0.5$



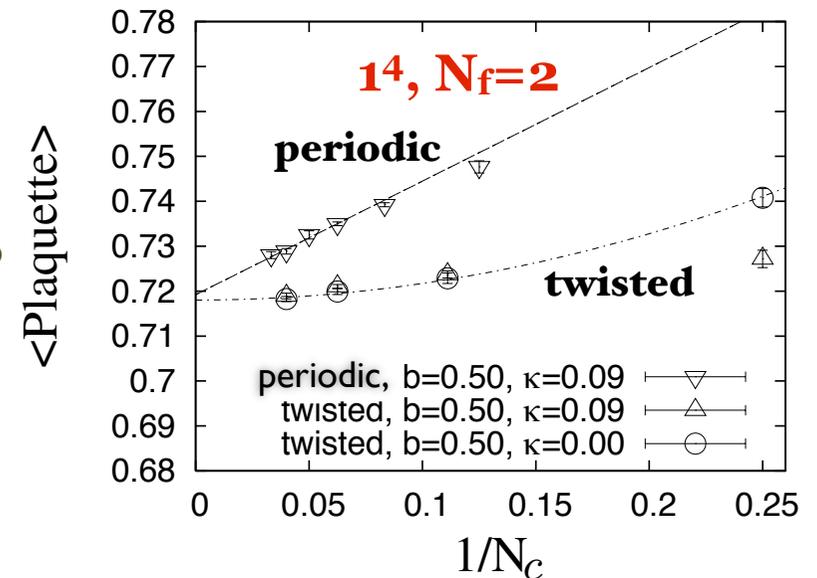
plaquette value scales as $1/N_c^2$.

cf) Plaquette value scales as $1/N_c$ in single-site simulation with two heavy adjoint fermions.

Bringoltz, Koren, Sharpe, 2011 Azeyanagi et al, 2010

1/N_c correction is expected to be suppressed by $V=1/2^4$

$$\langle \square \rangle = \langle \square_\infty \rangle + C_1 \left(\frac{N_c V}{V_{eff}} \right)^{-1} + C_2 N_c^{-2} + \dots$$



● Backup slide 2 - $1/N_c$ scaling, zero modes and V_{eff}

● Spacing of the low-lying Dirac eigenvalues

In the 't Hooft limit, $\langle \bar{\psi}\psi \rangle \sim N_c^2$ then $\Delta\lambda \sim \frac{1}{N_c^2}$ **Leutwyler & Smilga 1992**

Our finding: $\Delta\lambda \sim \frac{1}{N_c}$

In compact space, momentum-zero modes are not gauged away and thus the low-lying Dirac spectrum can be determined by the zero modes, where their number scales as N_c . **Azeyanagi, Hanada, Unsal, Yacoby. 2010**

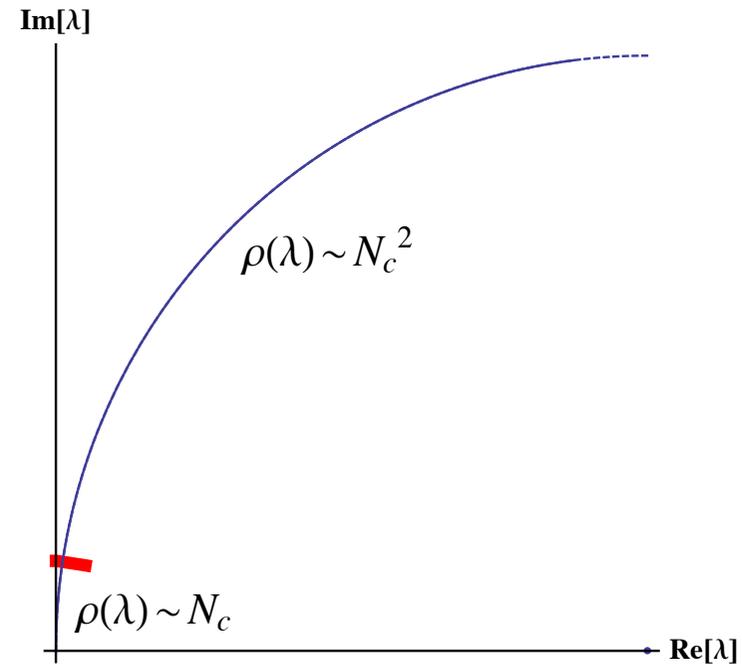
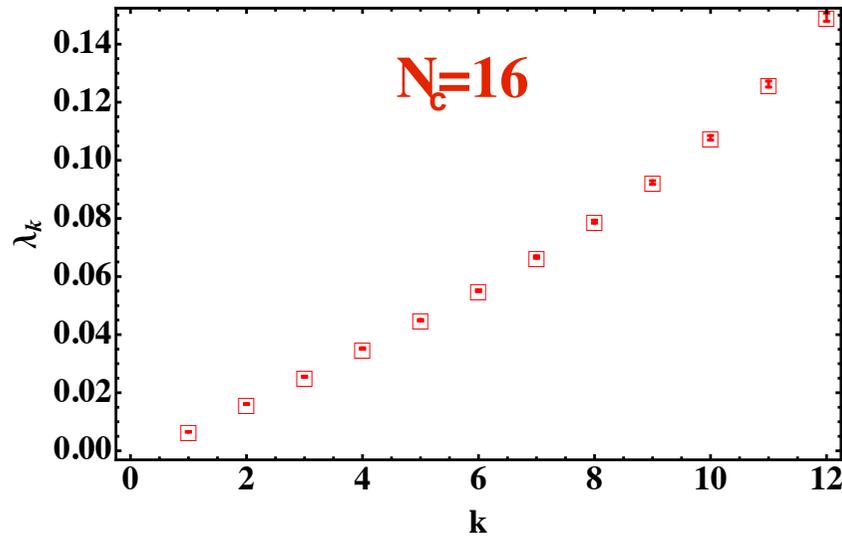
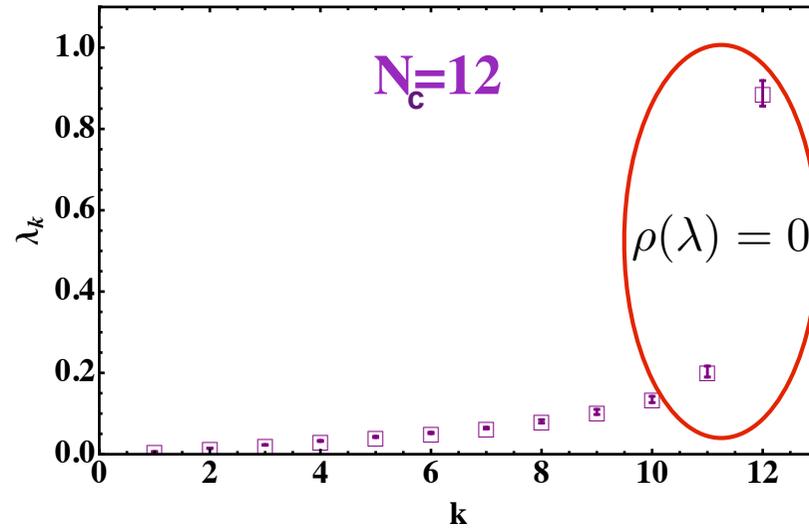
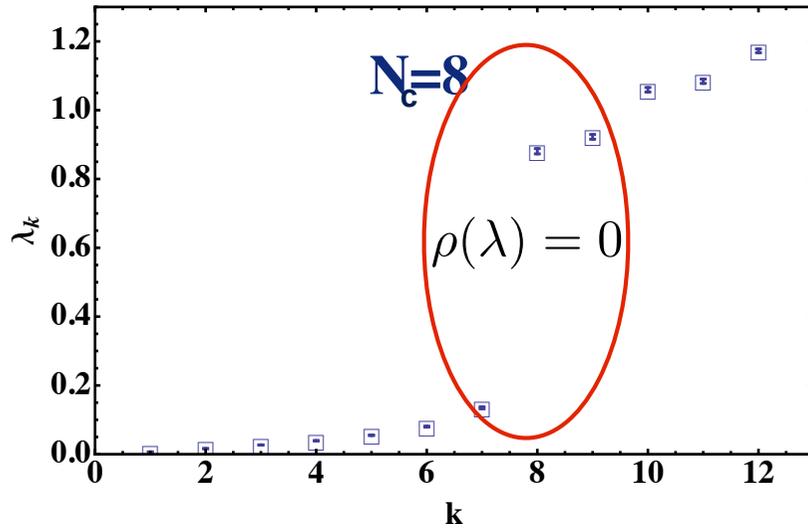
● Effective volumes V_{eff}

Finite N_c correction to the plaquette values and N_c dependence of the low-lying Dirac spectrum presumably indicate that the effective volume scales as N_c for AEK model with periodic b.c.

Volume expansion as an orbifold projection for adjoint QCD.

Kovtun, Unsal & Yaffe 2007

Backup slide 3 - *Bump* in the Dirac eigenvalues



of zero modes of the link matrix = $N_c - 1$